

## A reformulation of the generalized $q$ -Painlevé VI system with $W(A_{2n+1}^{(1)})$ symmetry

TAKAO SUZUKI<sup>†</sup>

Department of Mathematics, Kindai University, 3-4-1, Kowakae, Higashi-Osaka,  
Osaka 577-8502, Japan

<sup>†</sup>Corresponding author. Email: suzuki@math.kindai.ac.jp

Communicated by: Masatoshi Noumi

[Received on 28 September 2016; editorial decision on 21 December 2016; accepted on 24 December 2016]

In the previous work, we introduced the higher order  $q$ -Painlevé system  $q$ - $P_{(n+1,n+1)}$  as a generalization of the Jimbo–Sakai’s  $q$ -Painlevé VI equation. It is derived from a  $q$ -analogue of the Drinfeld–Sokolov hierarchy of type  $A_{2n+1}^{(1)}$  and admits a particular solution in terms of the Heine’s  $q$ -hypergeometric function  ${}_{n+1}\phi_n$ . However, the obtained system is insufficient as a generalization of  $q$ - $P_{VI}$  due to some reasons. In this article, we rewrite the system  $q$ - $P_{(n+1,n+1)}$  to a more suitable one.

*Keywords:* discrete Painlevé equations; affine Weyl groups; basic hypergeometric functions.

### 1. Introduction

Several generalizations of the Painlevé VI equation ( $P_{VI}$ ) have been proposed ([1, 3, 5, 11, 12, 14, 15, 19]). We focus on the higher order Painlevé system  $P_{(n+1,n+1)}$  given in [1, 14], or equivalently the Schlesinger system  $\mathcal{H}_{n+1,1}$  given in [19], among them. It can be regarded as a generalization from a viewpoint of a particular solution in terms of the hypergeometric function  ${}_{n+1}F_n$  ([13, 20]). The aim of this article is to introduce its  $q$ -analogue. This  $q$ -difference equation becomes a generalization of the  $q$ -Painlevé VI equation ( $q$ - $P_{VI}$ ) given in [4].

The investigation of generalizations of  $q$ - $P_{VI}$  has been developed in recent years ([8–10, 16, 18]). In the previous work [16], we proposed the higher order  $q$ -Painlevé system  $q$ - $P_{(n+1,n+1)}$ , whose explicit formula will be given in Section 2, as a  $q$ -analogue of  $P_{(n+1,n+1)}$ . It is derived from the  $q$ -Drinfeld–Sokolov hierarchy of type  $A_{2n+1}^{(1)}$ , contains  $q$ - $P_{VI}$  in the case of  $n = 1$  and admits a particular solution in terms of the  $q$ -hypergeometric function  ${}_{n+1}\phi_n$ . However, this system is insufficient as a generalization of  $q$ - $P_{VI}$  due to the following two reasons.

- (1) The system  $q$ - $P_{(n+1,n+1)}$  is probably reducible and reduces to a system of  $2n$ -th order.
- (2) We do not express the backward  $q$ -shifts  $(\underline{x}_i, \underline{y}_i)$  as functions in  $(x_i, y_i)$ .

The aim of this article is to solve those two problems. We reduce the system  $q$ - $P_{(n+1,n+1)}$  to a more suitable one as a generalization of  $q$ - $P_{VI}$ .

This article is organized as follows. In Section 2, we recall the definition of  $q$ - $P_{(n+1,n+1)}$  and its properties, namely a Lax pair, an affine Weyl group symmetry, a relationship with  $q$ - $P_{VI}$  and a particular

solution in terms of  ${}_{n+1}\phi_n$ . In Section 3, we formulate a system of  $q$ -difference equations of  $2n$ -th order which is equivalent to  $q$ - $P_{VI}$  in the case of  $n = 1$ . It is the main result of this article. In Section 4, we describe an action of the affine Weyl group on the  $2n$ -th order system given in the previous section.

## 2. Review: higher order $q$ -Painlevé system $q$ - $P_{(n+1,n+1)}$

In the following, we use notations

$$\overline{x(t)} = x(qt), \quad \underline{x(t)} = x(q^{-1}t),$$

where  $t, q \in \mathbb{C}$  and  $|q| < 1$ .

The system  $q$ - $P_{(n+1,n+1)}$  given in [16] is expressed as a system of  $q$ -difference equations

$$\begin{cases} x_{i-1} - x_i = \frac{b_{i-1}\underline{x}_{i-1}}{1 + \underline{x}_{i-1}y_{i-1}} - \frac{a_i x_i}{1 + \underline{x}_i y_{i-1}} \\ \underline{y}_{i-1} - \underline{y}_i = \frac{a_i y_{i-1}}{1 + \underline{x}_i y_{i-1}} - \frac{b_i y_i}{1 + \underline{x}_i y_i} \end{cases} \quad (i = 1, \dots, n+1),$$

with a constraint

$$\prod_{i=1}^{n+1} a_i \frac{1 + \underline{x}_i y_i}{1 + \underline{x}_i y_{i-1}} = q^{-n/2}, \quad (2.1)$$

where

$$b_0 = qb_{n+1}, \quad x_0 = tx_{n+1}, \quad y_0 = \frac{q}{t}y_{n+1}.$$

**REMARK 2.1** In the previous work [18], a higher order generalizations of  $q$ - $P_{VI}$  were presented by Tsuda. Since his  $q$ -Painlevé system can be regarded as a  $q$ -analogue of the Schlesinger system  $\mathcal{H}_{n+1,1}$ , we conjecture that his system coincides with  $q$ - $P_{(n+1,n+1)}$ . However, a relationship between both  $q$ -Painlevé systems has not been clarified yet.

We derived the system  $q$ - $P_{(n+1,n+1)}$  by a similarity reduction from the  $q$ -Drinfeld–Sokolov hierarchy of type  $A_{2n+1}^{(1)}$ . Hence the following theorem is obtained naturally via the construction of the system.

**THEOREM 2.2 ([16])** The system  $q$ - $P_{(n+1,n+1)}$  is given as the compatibility condition of a system of linear  $q$ -difference equations

$$\Psi(q^{-1}z, t) = M(z, t)\Psi(z, t), \quad \Psi(z, q^{-1}t) = B(z, t)\Psi(z, t). \quad (2.2)$$



for  $j = 1, \dots, n + 1$ . Also let  $\pi$  be a birational transformation defined by

$$\begin{aligned}\pi(a_i) &= q^{-\rho_1} b_i, & \pi(b_i) &= q^{-\rho_1} a_{i+1} \quad (i = 1, \dots, n), \\ \pi(a_{n+1}) &= q^{-\rho_1} b_{n+1}, & \pi(b_{n+1}) &= q^{-\rho_1-1} a_1, & \pi(\rho_1) &= -\rho_1 - \frac{1}{n+1}, \\ \pi(x_i) &= q^{-2\rho_1} t^{\rho_1} \underline{y}_i, & \pi(y_i) &= q^{\rho_1} t^{-\rho_1} \underline{x}_{i+1} \quad (i = 1, \dots, n), \\ \pi(x_{n+1}) &= q^{-2\rho_1} t^{\rho_1} \underline{y}_{n+1}, & \pi(y_{n+1}) &= q^{\rho_1+1} t^{-\rho_1-1} \underline{x}_1, & \pi(t) &= \frac{q^2}{t},\end{aligned}\tag{2.5}$$

where

$$q^{\rho_1} = (q^n a_1 b_1 \dots a_{n+1} b_{n+1})^{1/(n+1)}.$$

Then the system  $q$ - $P_{(n+1, n+1)}$  is invariant under actions of the transformations  $r_0, \dots, r_{2n+1}$  and  $\pi$ . Furthermore, the group of symmetries  $\langle r_0, \dots, r_{2n+1}, \pi \rangle$  is isomorphic to the extended affine Weyl group of type  $A_{2n+1}^{(1)}$ . Namely those transformations satisfy the fundamental relations

$$\begin{aligned}r_i^2 &= 1, & (r_i r_j)^{2-a_{ij}} &= 1 \quad (i, j = 0, \dots, 2n+1; i \neq j), \\ \pi^{2n+2} &= 1, & \pi r_i &= r_{i+1} \pi, & \pi r_{2n+1} &= r_0 \pi \quad (i = 0, \dots, 2n),\end{aligned}$$

where

$$\begin{aligned}a_{i,i} &= 2 & (i = 0, \dots, 2n+1), \\ a_{i,i+1} &= a_{2n+1,0} = a_{i+1,i} = a_{0,2n+1} = -1 & (i = 0, \dots, 2n), \\ a_{i,j} &= 0 & (\text{otherwise}).\end{aligned}$$

The system  $q$ - $P_{(2,2)}$  can be reduced to  $q$ - $P_{VI}$ .

**THEOREM 2.4 ([16])** If, in the system  $q$ - $P_{(2,2)}$ , we set

$$f = \frac{t(x_2 - x_1)\xi_1}{\xi_2}, \quad g = \frac{x_2(qt + x_1 y_2)\psi_1}{(1 + x_2 y_2)\psi_2},\tag{2.6}$$

where

$$\begin{aligned}\xi_1 &= (x_1 - x_2)(y_0 - y_1) - (a_1 - b_1), \\ \xi_2 &= (tx_2 - x_1)(x_1 - x_2)(y_0 - y_1) + (a_1 - b_1)x_1 + \{(b_1 - a_2)t - (a_1 - a_2)\}x_2, \\ \psi_1 &= q^{1/2}(q^{1/2} - a_1 b_1 t)\underline{x}_2 y_2 + (1 - q^{1/2} a_1 b_1)t, \\ \psi_2 &= q^{1/2} a_2 (q^{1/2} - a_1 b_1 t)\underline{x}_1 \underline{x}_2 y_2 + a_1 (1 - q^{1/2} b_1 a_2)t\underline{x}_1 - (a_1 - a_2)t\underline{x}_2,\end{aligned}$$

then they satisfy the  $q$ -Painlevé VI equation

$$\frac{f\bar{f}}{\alpha_3 \alpha_4} = \frac{(\bar{g} - t\beta_1)(\bar{g} - t\beta_2)}{(\bar{g} - \beta_3)(\bar{g} - \beta_4)}, \quad \frac{g\bar{g}}{\beta_3 \beta_4} = \frac{(f - t\alpha_1)(f - t\alpha_2)}{(f - \alpha_3)(f - \alpha_4)},$$

with parameters

$$\alpha_1 = 1, \quad \alpha_2 = q^{1/2}a_1b_1, \quad \alpha_3 = 1, \quad \alpha_4 = \frac{1}{q^{1/2}a_2b_2},$$

$$\beta_1 = q^{1/2}b_1, \quad \beta_2 = q^{1/2}a_1, \quad \beta_3 = \frac{1}{qa_2}, \quad \beta_4 = \frac{1}{b_2}.$$

REMARK 2.5 In [16] we defined the transformation  $\pi$  by

$$\pi(a_i) = b_i, \quad \pi(b_i) = a_{i+1} \quad (i = 1, \dots, n), \quad \pi(a_{n+1}) = b_{n+1}, \quad \pi(b_{n+1}) = \frac{a_1}{q},$$

$$\pi(x_i) = \underline{y}_i, \quad \pi(y_i) = \underline{x}_{i+1} \quad (i = 1, \dots, n), \quad \pi(x_{n+1}) = \underline{y}_{n+1}, \quad \pi(y_{n+1}) = \frac{q}{t}x_1, \quad \pi(t) = \frac{q^2}{t}.$$

As a matter of fact, unless we replace constraint (2.1) with

$$\prod_{i=1}^{n+1} \frac{a_i^{1/2}}{b_i^{1/2}} \frac{1 + \underline{x}_i y_i}{1 + \underline{x}_i y_{i-1}} = q^{1/4},$$

the system  $q$ - $P_{(n+1,n+1)}$  is not invariant under an action of  $\pi$ . If we do so, then the system  $q$ - $P_{(2,2)}$  seems to reduce not to  $q$ - $P_{V1}$  but to a  $q$ -analogue of the Painlevé V equation. This  $q$ -difference equation was derived in [7] from a binational representation of the extended affine Weyl group of type  $A_1^{(1)} \times A_3^{(1)}$  given in [6]. Afterward that equation was found to be a subsystem of  $q$ - $P_{V1}$  in [17].

The system  $q$ - $P_{(n+1,n+1)}$  admits a particular solution in terms of the  $q$ -hypergeometric function  ${}_{n+1}\phi_n$  defined by the formal power series

$${}_{n+1}\phi_n \left[ \begin{matrix} \alpha_1, \dots, \alpha_n, \alpha_{n+1} \\ \beta_1, \dots, \beta_n \end{matrix} ; q, t \right] = \sum_{k=0}^{\infty} \frac{(\alpha_1; q)_k \dots (\alpha_n; q)_k (\alpha_{n+1}; q)_k}{(\beta_1; q)_k \dots (\beta_n; q)_k (q; q)_k} t^k,$$

where  $(\alpha; q)_k$  stands for the  $q$ -shifted factorial

$$(\alpha; q)_0 = 1, \quad (\alpha; q)_k = (1 - \alpha)(1 - q\alpha) \dots (1 - q^{k-1}\alpha) \quad (k \geq 1).$$

THEOREM 2.6 ([16]) If, in the system  $q$ - $P_{(n+1,n+1)}$ , we assume that

$$y_i = 0 \quad (i = 1, \dots, n + 1), \quad \prod_{i=1}^{n+1} a_i = q^{-n/2},$$

then a vector of the variables  $\mathbf{x} = {}^t[x_1, \dots, x_{n+1}]$  satisfies a system of linear  $q$ -difference equations

$$\bar{\mathbf{x}} = \left( A_0 + \frac{A_1}{1 - qt} \right) \mathbf{x}, \tag{2.7}$$

with  $(n + 1) \times (n + 1)$  matrices

$$A_0 = \begin{bmatrix} b_1 & b_2 - a_2 & b_3 - a_3 & \dots & b_n - a_n & b_{n+1} - a_{n+1} \\ & b_2 & b_3 - a_3 & \dots & b_n - a_n & b_{n+1} - a_{n+1} \\ & & & \ddots & \vdots & \vdots \\ & & & & b_n - a_n & b_{n+1} - a_{n+1} \\ & O & & & b_n & b_{n+1} - a_{n+1} \\ & & & & & b_{n+1} \end{bmatrix},$$

$$A_1 = \begin{bmatrix} 1 \\ 1 \\ \vdots \\ 1 \\ 1 \\ 1 \end{bmatrix} [a_1 - b_1 \quad a_2 - b_2 \quad a_3 - b_3 \quad \dots \quad a_n - b_n \quad a_{n+1} - b_{n+1}].$$

Furthermore system (2.7) admits a solution

$$\mathbf{x} = t^{-\log_q a_1} \begin{bmatrix} \phi_1 \\ \vdots \\ \phi_{n+1} \end{bmatrix},$$

$$\phi_j = \prod_{i=1}^{j-1} \frac{b_i - a_1}{a_{i+1} - a_1} {}_{n+1}\phi_n \left[ \begin{matrix} q^{\frac{a_1}{b_1}}, \dots, q^{\frac{a_1}{b_{j-1}}}, \frac{a_1}{b_j}, \dots, \frac{a_1}{b_n}, \frac{a_1}{b_{n+1}} \\ q^{\frac{a_1}{a_2}}, \dots, q^{\frac{a_1}{a_j}}, \frac{a_1}{a_{j+1}}, \dots, \frac{a_1}{a_n} \end{matrix} ; q, q^{(n+2)/2} b_1 \dots b_{n+1} t \right].$$

Therefore, we want to regard the system  $q\text{-}P_{(n+1, n+1)}$  as a generalization of  $q\text{-}P_{\text{VI}}$  from a viewpoint of a particular solution in terms of the  $q$ -hypergeometric function. However, as is seen in Section 1, we have two problems. In [16], we derived  $q\text{-}P_{\text{VI}}$  by reducing linear system (2.2) to the one with  $2 \times 2$  matrices given in [4] from  $q\text{-}P_{(2,2)}$ . We could not use a similar method in a general case, although we reduced linear system (2.2) to the one with  $(n + 1) \times (n + 1)$  matrices in [2]. Definition of dependent variables (2.6) is complicated and hence unsuitable for a generalization. In this article, we choose a suitable set of  $2n$  dependent variables  $(f_i, g_i)$  and reduce the system  $q\text{-}P_{(n+1, n+1)}$  to a one of  $2n$ -th order. In the obtained system, the forward  $q$ -shifts  $(\bar{f}_i, \bar{g}_i)$  is expressed as functions in  $(f_i, g_i)$ .

### 3. Main result

The key to solving is the affine Weyl group symmetry of  $q\text{-}P_{(n+1, n+1)}$ . We can simplify definition of dependent variables (2.6) as

$$r_1 r_2(f) = -t \frac{x_1 - x_2}{tx_2 - x_1}, \quad r_1 r_2(g) = \frac{a_2 t}{q^{1/2}} \frac{x_2(1 + \underline{x_1}y_1)}{x_1(1 + \underline{x_2}y_1)}.$$

This fact suggests a choice of dependent variables of a  $2n$ -th order system.

**THEOREM 3.1** If, in the system  $q$ - $P_{(n+1,n+1)}$ , we set

$$f_i = t \frac{x_i - x_{i+1}}{tx_{n+1} - x_1}, \quad g_i = a_{i+1} \frac{x_{i+1}(1 + x_i y_i)}{x_i(1 + x_{i+1} y_i)} \quad (i = 1, \dots, n), \tag{3.1}$$

then they satisfy a system of  $q$ -difference equations

$$f_i \bar{f}_i = qt \frac{F_i F_{i+1} \bar{g}_0 (b_i - \bar{g}_i) (\bar{g}_i - a_{i+1})}{F_{n+1} F_1 \bar{g}_i (b_0 - \bar{g}_0) (\bar{g}_0 - a_1)} \quad (i = 1, \dots, n), \tag{3.2}$$

$$g_i \bar{g}_i = \frac{F_{i+1} G_i}{F_i G_{i+1}} \quad (i = 1, \dots, n), \tag{3.3}$$

where

$$b_0 = qb_{n+1}, \quad g_0 = \frac{1}{q^{(n-2)/2} t} \prod_{i=1}^n \frac{1}{g_i} = \frac{qa_1 x_1 (1 + x_{n+1} y_{n+1})}{t x_{n+1} (1 + x_1 y_0)} = a_1 \frac{x_1 (1 + x_0 y_0)}{x_0 (1 + x_1 y_0)},$$

and

$$\begin{aligned} F_i &= \sum_{j=1}^{i-1} f_j + t \sum_{j=i}^n f_j + t, \\ G_i &= \sum_{j=i}^n \prod_{k=i}^{j-1} b_k a_{k+1} \frac{\prod_{l=j+1}^n g_l}{\prod_{l=1}^{j-1} g_l} f_j + q^{n/2} t \prod_{k=i}^n b_k a_{k+1} \\ &\quad + q^n t \sum_{j=1}^{i-1} \frac{b_{n+1} a_1 \prod_{k=1}^n b_k a_{k+1}}{\prod_{k=j}^{i-1} b_k a_{k+1}} \frac{\prod_{l=j+1}^n g_l}{\prod_{l=1}^{j-1} g_l} f_j \quad (i = 1, \dots, n+1). \end{aligned}$$

Furthermore, if we set  $h = tx_{n+1} - x_1$ , then it satisfies a linear  $q$ -difference equation

$$\bar{h} = -\frac{F_{n+1} F_1 (b_0 - \bar{g}_0) (\bar{g}_0 - a_1)}{t(t-1)^2 \bar{g}_0} h. \tag{3.4}$$

System (3.2), (3.3) is equivalent to  $q$ - $P_{VI}$  in the case of  $n = 1$ . In this section, we prove this theorem. We also discuss the relationship between the dependent variables  $f_i, g_i, h$  ( $i = 1, \dots, n$ ) and the ones  $x_j, y_j$  ( $j = 1, \dots, n+1$ ) in more detail at the end of this section.

### 3.1 Proof of equation (3.2)

Definition of dependent variables (3.1) implies

$$g_i - a_{i+1} = -a_{i+1} \frac{x_i - x_{i+1}}{x_i(1 + x_{i+1} y_i)} \quad (i = 0, 1, \dots, n). \tag{3.5}$$

The first equation of  $q$ - $P_{(n+1,n+1)}$  can be rewritten to

$$b_i - g_i = \frac{1 + \underline{x}_i y_i}{\underline{x}_i} (x_i - x_{i+1}) \quad (i = 0, 1, \dots, n). \quad (3.6)$$

Combining them, we obtain

$$\begin{aligned} \frac{(b_0 - g_0)(g_0 - a_1)}{g_0} &= -\frac{(t\underline{x}_{n+1} - \underline{x}_1)(t\underline{x}_{n+1} - x_1)}{t\underline{x}_{n+1}\underline{x}_1}, \\ \frac{(b_i - g_i)(g_i - a_{i+1})}{g_i} &= -\frac{(x_i - \underline{x}_{i+1})(x_i - x_{i+1})}{\underline{x}_i \underline{x}_{i+1}} \quad (i = 1, \dots, n). \end{aligned} \quad (3.7)$$

Hence we can derive equation (3.2) by using

$$F_i = \frac{t(t-1)x_i}{t\underline{x}_{n+1} - x_1} \quad (i = 1, \dots, n+1).$$

### 3.2 Proof of equations (3.3) and (3.4)

Equation (3.5) can be rewritten to

$$\begin{aligned} y_0 &= -\frac{x_0 g_0 - a_1 x_1}{x_0 x_1 (g_0 - a_1)} = -\frac{t\underline{x}_{n+1} g_0 - q a_1 x_1}{t\underline{x}_{n+1} x_1 (g_0 - a_1)} = \frac{q}{t} y_{n+1}, \\ y_i &= -\frac{\underline{x}_i g_i - a_{i+1} x_{i+1}}{\underline{x}_i x_{i+1} (g_i - a_{i+1})} \quad (i = 1, \dots, n). \end{aligned}$$

Substituting it to the second equation of  $q$ - $P_{(n+1,n+1)}$ , we obtain

$$\begin{aligned} \underline{y}_{i-1} - \underline{y}_i &= \frac{\underline{x}_{i-1} g_{i-1}}{\underline{x}_i (\underline{x}_{i-1} - \underline{x}_i)} - \frac{a_i}{\underline{x}_{i-1} - \underline{x}_i} + \frac{b_i a_{i+1} \underline{x}_{i+1}}{\underline{x}_i (\underline{x}_i - \underline{x}_{i+1}) g_i} - \frac{b_i}{\underline{x}_i - \underline{x}_{i+1}} \quad (i = 1, \dots, n), \\ \underline{y}_n - \underline{y}_{n+1} &= \frac{x_n g_n}{x_{n+1} (\underline{x}_n - \underline{x}_{n+1})} - \frac{a_{n+1}}{\underline{x}_n - \underline{x}_{n+1}} + \frac{q b_{n+1} a_1 x_1}{x_{n+1} (t\underline{x}_{n+1} - q x_1) g_0} - \frac{b_{n+1} t}{t\underline{x}_{n+1} - q x_1}. \end{aligned}$$

It implies

$$\begin{aligned} &b_i(x_{i-1} - x_i) + a_i(x_i - x_{i+1}) + (x_{i-1} - x_i)(x_i - x_{i+1})(y_{i-1} - y_i) \\ &= \frac{x_{i-1}}{x_i} (x_i - x_{i+1}) \overline{g_{i-1}} + b_i a_{i+1} \frac{x_{i+1}}{x_i} (x_{i-1} - x_i) \frac{1}{g_i} \quad (i = 1, \dots, n), \\ &b_{n+1} t(x_n - x_{n+1}) + a_{n+1}(t\underline{x}_{n+1} - x_1) + (x_n - x_{n+1})(t\underline{x}_{n+1} - x_1)(y_n - y_{n+1}) \\ &= \frac{x_n}{x_{n+1}} (t\underline{x}_{n+1} - x_1) \overline{g_n} + b_{n+1} a_1 \frac{x_1}{x_{n+1}} (x_n - x_{n+1}) \frac{1}{g_0}. \end{aligned} \quad (3.8)$$



On the other hand, equation (3.6) can be rewritten to

$$y_0 = -\frac{g_0 - b_0}{x_0 - x_1} - \frac{1}{x_0} = -\frac{g_0 - qb_{n+1}}{tx_{n+1} - x_1} - \frac{q}{tx_{n+1}} = \frac{q}{t}y_{n+1},$$

$$y_i = -\frac{g_i - b_i}{x_i - x_{i+1}} - \frac{1}{x_i} \quad (i = 1, \dots, n).$$

Combining it with equation (3.7), we obtain

$$y_{i-1} - y_i = \frac{g_i - b_i}{x_i - x_{i+1}} + \frac{1}{x_i} - \frac{g_{i-1} - b_{i-1}}{x_{i-1} - x_i} - \frac{1}{x_{i-1}}$$

$$= \frac{g_i}{x_i - x_{i+1}} - \frac{b_i}{x_i - x_{i+1}} + \frac{b_{i-1}a_i}{(x_{i-1} - x_i)g_{i-1}} - \frac{a_i}{x_{i-1} - x_i} \quad (i = 1, \dots, n),$$

$$y_n - y_{n+1} = \frac{tg_0 - qb_{n+1}t}{q(tx_{n+1} - x_n)} + \frac{1}{x_{n+1}} - \frac{g_n - b_n}{x_n - x_{n+1}} - \frac{1}{x_n}$$

$$= \frac{tg_0}{q(tx_{n+1} - x_1)} - \frac{b_{n+1}t}{tx_{n+1} - x_1} + \frac{b_n a_{n+1}}{(x_n - x_{n+1})g_n} - \frac{a_{n+1}}{x_n - x_{n+1}}.$$

We can rewrite it to

$$b_i(x_{i-1} - x_i) + a_i(x_i - x_{i+1}) + (x_{i-1} - x_i)(x_i - x_{i+1})(y_{i-1} - y_i)$$

$$= (x_{i-1} - x_i)g_i + b_{i-1}a_i(x_i - x_{i+1})\frac{1}{g_{i-1}} \quad (i = 1, \dots, n),$$

$$b_{n+1}t(x_n - x_{n+1}) + a_{n+1}(tx_{n+1} - x_1) + (x_n - x_{n+1})(tx_{n+1} - x_1)(y_n - y_{n+1})$$

$$= \frac{t}{q}(x_n - x_{n+1})g_0 + b_n a_{n+1}(tx_{n+1} - x_1)\frac{1}{g_n}. \tag{3.9}$$

Combining equations (3.8) and (3.9), we obtain

$$f_1 \frac{F_{n+1}}{F_1} \bar{g}_0 + tb_1 a_2 \frac{F_2}{F_1} \frac{1}{\bar{g}_1} = tg_1 + qb_{n+1} a_1 f_1 \frac{1}{g_0},$$

$$f_i \frac{F_{i-1}}{F_i} \bar{g}_{i-1} + b_i a_{i+1} f_{i-1} \frac{F_{i+1}}{F_i} \frac{1}{\bar{g}_i} = f_{i-1} g_i + b_{i-1} a_i f_i \frac{1}{g_{i-1}} \quad (i = 2, \dots, n),$$

$$t \frac{F_n}{F_{n+1}} \bar{g}_n + b_{n+1} a_1 f_n \frac{F_1}{F_{n+1}} \frac{1}{\bar{g}_0} = \frac{t}{q} f_n g_0 + b_n a_{n+1} t \frac{1}{g_n}. \tag{3.10}$$

Since equation (3.10) is equivalent to (3.3) in the case of  $n = 1$ , we consider the case of  $n \geq 2$ . Then equation (3.10) is transformed to

$$\begin{aligned} G_{n+1} - q^{n-1}t \left( t \frac{g_0 g_1}{f_1} + q b_{n+1} a_1 \right) G_1 + q^{n-1} b_1 a_2 t^2 \frac{g_0 g_1}{f_1} G_2 &= 0, \\ G_{i-1} - \left( \frac{f_{i-1} g_{i-1} g_i}{f_i} + b_{i-1} a_i \right) G_i + b_i a_{i+1} \frac{f_{i-1} g_{i-1} g_i}{f_i} G_{i+1} &= 0 \quad (i = 2, \dots, n), \\ G_n - \left( \frac{1}{q} f_n g_n g_0 + b_n a_{n+1} \right) G_{n+1} + q^{n-1} b_{n+1} a_1 t f_n g_n g_0 G_1 &= 0, \end{aligned} \quad (3.11)$$

via a transformation

$$\begin{aligned} \bar{g}_i &= \frac{F_{i+1} G_i}{g_i F_i G_{i+1}} \quad (i = 1, \dots, n), \\ \bar{g}_0 &= \frac{1}{q^{n/2} t} \prod_{i=1}^n \frac{1}{\bar{g}_i} = \frac{1}{q^{n-1} t^2} \frac{F_1 G_{n+1}}{g_0 F_{n+1} G_1}. \end{aligned}$$

Furthermore equation (3.11) is reduced to a system of linear equations

$$\begin{bmatrix} 1 & -\alpha_1 & \beta_1 & & & & & & & & \\ & 1 & -\alpha_2 & \beta_2 & & & & & & & \\ & & 1 & -\alpha_3 & \beta_3 & & & & & & \\ & & & & & \ddots & & & & & \\ & & & & & & & & & & \\ & & & & & & & & & & \\ & & & & & & & & & & \\ & & & & & & -\alpha_{n-2} & \beta_{n-2} & & & \\ & & & & & & 1 & -\alpha_{n-1} & & & \\ & & & & & & & 1 & & & \\ \beta_n & & & & & & & & & & \end{bmatrix} \begin{bmatrix} G_1 \\ G_2 \\ G_3 \\ \vdots \\ G_{n-2} \\ G_{n-1} \\ G_n \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ \vdots \\ 0 \\ -\beta_{n-1} G_{n+1} \\ \alpha_n G_{n+1} \end{bmatrix}, \quad (3.12)$$

where

$$\begin{aligned} \alpha_i &= \frac{f_i g_i g_{i+1}}{f_{i+1}} + b_i a_{i+1}, & \beta_i &= b_{i+1} a_{i+2} \frac{f_i g_i g_{i+1}}{f_{i+1}} \quad (i = 1, \dots, n-1), \\ \alpha_n &= \frac{1}{q} f_n g_n g_0 + b_n a_{n+1}, & \beta_n &= q^{n-1} b_{n+1} a_1 t f_n g_n g_0. \end{aligned}$$

In fact equation (3.11) can be rewritten to

$$\begin{aligned} G_1 - b_1 a_2 G_2 &= \frac{q f_1}{t g_0 g_1} \left( \frac{1}{q^n t} G_{n+1} - b_{n+1} a_1 G_1 \right), \\ G_i - b_i a_{i+1} G_{i+1} &= \frac{f_i}{f_{i-1} g_{i-1} g_i} (G_{i-1} - b_{i-1} a_i G_i) \quad (i = 2, \dots, n), \\ \frac{1}{q^n t} G_{n+1} - b_{n+1} a_1 G_1 &= \frac{1}{q^{n-1} t f_n g_n g_0} (G_n - b_n a_{n+1} G_{n+1}), \end{aligned} \quad (3.13)$$



via a cofactor expansion. The determinant of the tridiagonal matrix  $\Delta_i$  satisfies a recurrence relation

$$\Delta_0 = 1, \quad \Delta_1 = -\alpha_1, \quad \Delta_i = -\alpha_i \Delta_{i-1} - \beta_{i-1} \Delta_{i-2} \quad (i = 2, 3, \dots).$$

By solving this relation, we obtain

$$\Delta_i = (-1)^i \sum_{j=1}^{i+1} \prod_{k=1}^{j-1} b_k a_{k+1} \prod_{l=j}^i g_l g_{l+1} \frac{f_j}{f_{i+1}} \quad (i = 2, 3, \dots).$$

It implies

$$\begin{aligned} (-1)^{n-1} \beta_n \Delta_{n-1} &= q^{n-1} b_{n+1} a_1 t f_n g_n g_0 \sum_{j=1}^n \prod_{k=1}^{j-1} b_k a_{k+1} \prod_{l=j}^{n-1} g_l g_{l+1} \frac{f_j}{f_n} \\ &= q^{n/2} \sum_{j=1}^n b_{n+1} a_1 \prod_{k=1}^{j-1} b_k a_{k+1} \frac{\prod_{l=j+1}^n g_l}{\prod_{l=1}^{j-1} g_l} f_j. \end{aligned}$$

□

Thanks to Lemma 3.2, we can find that system (3.12) admits only one solution for  $G_i$  ( $i = 1, \dots, n$ ). Hence we only have to verify that

$$G_i = \sum_{j=i}^n \prod_{k=i}^{j-1} b_k a_{k+1} \frac{\prod_{l=j+1}^n g_l}{\prod_{l=1}^{j-1} g_l} f_j + q^{n/2} t \prod_{k=i}^n b_k a_{k+1} + q^n t \sum_{j=1}^{i-1} \frac{b_{n+1} a_1 \prod_{k=1}^n b_k a_{k+1}}{\prod_{k=j}^{i-1} b_k a_{k+1}} \frac{\prod_{l=j+1}^n g_l}{\prod_{l=1}^{j-1} g_l} f_j,$$

satisfy system (3.12). It is shown since

$$\begin{aligned} G_i - b_i a_{i+1} G_{i+1} &= \left( 1 - q^n t b_{n+1} a_1 \prod_{k=1}^n b_k a_{k+1} \right) \frac{\prod_{l=i+1}^n g_l}{\prod_{l=1}^{i-1} g_l} f_i \quad (i = 1, \dots, n), \\ \frac{1}{q^n t} G_{n+1} - b_{n+1} a_1 G_1 &= \frac{1}{q^{n/2}} \left( 1 - q^n t b_{n+1} a_1 \prod_{k=1}^n b_k a_{k+1} \right), \end{aligned}$$

satisfy equation (3.13). From the above, we have derived equation (3.3).

In the last, we prove equation (3.4). The first equation of (3.7) is rewritten to

$$\frac{1}{tx_{n+1} - x_1} = -\frac{(b_0 - \bar{g}_0)(\bar{g}_0 - a_1)}{\bar{g}_0} \frac{tx_{n+1}x_1}{tx_{n+1} - x_1},$$

from which we obtain equation (3.4) immediately.

### 3.3 Relationship between two types of dependent variables

In equation (3.1), we define the dependent variables  $(f_i, g_i, h)$  as rational functions in  $(x_j, y_j)$  with constraint (2.1). Conversely, the dependent variables  $(x_j, y_j)$  are given as rational functions in  $(f_i, g_i, h)$ . We obtain

$$\frac{1}{x_1} - \frac{1}{tx_{n+1}} = -\frac{(b_0 - g_0)(g_0 - a_1)}{g_0h},$$

from equation (3.4) and

$$\frac{1}{x_{i+1}} - \frac{1}{x_i} = -\frac{t(b_i - g_i)(g_i - a_{i+1})}{f_i g_i h} \quad (i = 1, \dots, n),$$

from equation (3.7). They imply

$$\begin{aligned} \frac{1}{x_i} &= \frac{t}{q-t} \frac{(b_0 - g_0)(g_0 - a_1)}{g_0h} \\ &+ \frac{t}{q-t} \sum_{j=1}^{i-1} \frac{t(b_j - g_j)(g_j - a_{j+1})}{f_j g_j h} + \frac{q}{q-t} \sum_{j=i}^n \frac{t(b_j - g_j)(g_j - a_{j+1})}{f_j g_j h} \quad (i = 1, \dots, n+1). \end{aligned} \quad (3.14)$$

We also obtain

$$y_i = -\frac{g_i/x_{i+1} - a_{i+1}/x_i}{g_i - a_{i+1}} \quad (i = 1, \dots, n),$$

from equation (3.5). The rest variable  $y_{n+1}$  is given by constraint (2.1).

**REMARK 3.3** Equation (3.14) allows us to express the backward  $q$ -shifts  $f_1, \dots, f_n$  as rational functions in  $(f_i, g_i)$ . Also we can express the backward  $q$ -shifts  $g_1, \dots, g_n$  as rational functions in  $(f_i, g_i)$  by solving equation (3.10) for  $g_1, \dots, g_n$ . We don't give its detail here.

We next state a relationship between the variables  $(f_i, g_i, h)$  and the ones  $(x_j, y_j)$ . Equation (3.9) is rewritten to

$$\begin{aligned} \frac{q}{t}y_{n+1} - y_1 &= t(g_1 - b_1)\frac{1}{f_1h} + a_1 \left( qb_{n+1}\frac{1}{g_0} - 1 \right) \frac{1}{h} \\ y_{i-1} - y_i &= t(g_i - b_i)\frac{1}{f_ih} + a_it \left( b_{i-1}\frac{1}{g_{i-1}} - 1 \right) \frac{1}{f_{i-1}h} \quad (i = 2, \dots, n), \\ y_n - y_{n+1} &= \frac{t}{q}(g_0 - qb_{n+1})\frac{1}{h} + a_{n+1}t \left( b_n\frac{1}{g_n} - 1 \right) \frac{1}{f_nh}. \end{aligned} \quad (3.15)$$

Recall that

$$tx_{n+1} - x_1 = h, \quad x_i - x_{i+1} = \frac{f_i h}{t} \quad (i = 1, \dots, n).$$

Hence we can give the variables  $(x_j, y_j)$  as rational functions in  $(f_i, g_i, h)$ . Conversely, since equation (3.15) admits only one solution for  $g_1, \dots, g_n$ , the variables  $(f_i, g_i, h)$  are given as rational functions in  $(x_j, y_j)$ . Then one constraint between the variables  $(x_j, y_j)$  is obtained together. We don't give their explicit formulas here. Those facts allow us to express the backward  $q$ -shifts  $\underline{x}_j, \underline{y}_j$  as functions in  $(x_j, y_j)$ .

#### 4. Affine Weyl group symmetry

As is seen in Section 2, the system  $q$ - $P_{(n+1, n+1)}$  is invariant under the action of the group of symmetries  $(r_0, \dots, r_{2n+1}, \pi) \simeq \widetilde{W}(A_{2n+1}^{(1)})$ . This action can be restricted to systems (3.2) and (3.3).

**THEOREM 4.1** The birational transformations  $r_0, \dots, r_{2n+1}, \pi$  defined by (2.3), (2.4) and (2.5) act on the dependent variables  $f_i, g_i$  ( $i = 1, \dots, n$ ) as follows.

$$r_{2j-2}(f_i) = f_i, \quad r_{2j-2}(g_i) = g_i, \quad (4.1)$$

for  $j = 1, \dots, n+1$ ,

$$\begin{aligned} r_1(f_1) &= f_1 \frac{R_1^{a,a,a}}{R_1^{b,a,b}}, & r_1(g_1) &= g_1 \frac{R_1^{b,a,a}}{R_1^{b,b,b}}, & r_1(f_i) &= f_i \frac{R_1^{b,a,a}}{R_1^{b,a,b}}, & r_1(g_i) &= g_i \quad (i \neq 1), \\ r_{2j-1}(f_{j-1}) &= f_{j-1} \frac{R_j^{b,a,b}}{R_j^{b,a,a}}, & r_{2j-1}(g_{j-1}) &= g_{j-1} \frac{R_j^{b,b,b}}{R_j^{b,a,a}}, & r_{2j-1}(f_j) &= f_j \frac{R_j^{a,a,a}}{R_j^{b,a,a}}, & r_{2j-1}(g_j) &= g_j \frac{R_j^{b,a,a}}{R_j^{b,b,b}}, \\ r_{2j-1}(f_i) &= f_i, & r_{2j-1}(g_i) &= g_i \quad (i \neq j-1, j), \\ r_{2n+1}(f_n) &= f_n \frac{R_{n+1}^{b,a,b}}{R_{n+1}^{a,a,a}}, & r_{2n+1}(g_n) &= g_n \frac{R_{n+1}^{b,b,b}}{R_{n+1}^{b,a,a}}, & r_{2n+1}(f_i) &= f_i \frac{R_{n+1}^{b,a,a}}{R_{n+1}^{a,a,a}}, & r_{2n+1}(g_i) &= g_i \quad (i \neq n) \end{aligned} \quad (4.2)$$

for  $j = 2, \dots, n$  and

$$\begin{aligned} \pi(f_i) &= \frac{q^2 (g_i R_i^* - b_{i+1} R_{i+1}^*)(b_{i+1} - g_{i+1})(R_{i+1}^* + 1 - \frac{t}{q}) f_1}{t (g_0 R_0^* - b_1 R_1^*)(b_1 - g_1)(R_1^* + 1 - \frac{t}{q}) f_{i+1}}, & \pi(g_i) &= \frac{a_{i+1} b_{i+1} R_{i+1}^*}{q^{\rho_1} g_i R_i^*} \quad (i \neq n), \\ \pi(f_n) &= q \frac{(g_n R_n^* - b_{n+1} R_{n+1}^*)(b_0 - g_0)(R_0^* + 1 - \frac{t}{q}) f_1}{(g_0 R_0^* - b_1 R_1^*)(b_1 - g_1)(R_1^* + 1 - \frac{t}{q}) f_0}, & \pi(g_n) &= \frac{a_{n+1} b_{n+1} R_0^*}{q^{\rho_1} g_n R_n^*}, \end{aligned} \quad (4.3)$$

where

$$\begin{aligned} R_j^{\alpha, \beta, \gamma} &= (g_j - \alpha_j) \frac{1}{f_j} + \left( \beta_j b_{j-1} \frac{1}{g_{j-1}} - \gamma_j \right) \frac{1}{f_{j-1}} \quad (j \neq n+1), \\ R_{n+1}^{\alpha, \beta, \gamma} &= \frac{1}{q} (g_0 - q \alpha_{n+1}) + \left( \beta_{n+1} b_n \frac{1}{g_n} - \gamma_{n+1} \right) \frac{1}{f_n}, \end{aligned}$$

and

$$\begin{aligned}
 R_i^* &= -\frac{f_i}{b_i - g_i} \left( \frac{t}{q} \sum_{j=1}^i R_j^{b,a,a} + \sum_{j=i+1}^{n+1} R_j^{b,a,a} \right) \\
 &= \frac{f_i}{b_i - g_i} \left( \sum_{j=0}^{i-1} \frac{\frac{t}{q} (1 - \frac{a_{j+1}}{g_j})(b_j - g_j)}{f_j} + \frac{(\frac{t}{q} - \frac{a_{i+1}}{g_i})(b_i - g_i)}{f_i} + \sum_{j=i+1}^n \frac{(1 - \frac{a_{j+1}}{g_j})(b_j - g_j)}{f_j} \right).
 \end{aligned}$$

Here we set  $f_0 = t$ .

Recall that

$$b_0 = qb_{n+1}, \quad g_0 = \frac{1}{q^{(n-2)/2} t g_1 \dots g_n}, \quad q^{\rho_1} = (q^n a_1 b_1 \dots a_{n+1} b_{n+1})^{1/(n+1)}.$$

In this section, we prove this theorem.

#### 4.1 Proof of action (4.1)

The action of  $r_{2j-2}$  on the dependent variables with the exception of  $y_{j-1}$  is trivial for any  $j = 1, \dots, n + 1$ . Hence we only have to verify that  $r_{2j-2}(g_{j-1}) = g_{j-1}$  is satisfied. And it is shown by using the first equation of  $q$ - $P_{(n+1,n+1)}$  as follows

$$\begin{aligned}
 r_{2j-2}(g_{j-1}) &= b_{j-1} \frac{\underline{x}_j}{\underline{x}_{j-1}} \frac{1 + \underline{x}_{j-1} y_{j-1} - \frac{b_{j-1} - a_j}{\underline{x}_{j-1} - \underline{x}_j} \underline{x}_{j-1}}{1 + \underline{x}_j y_{j-1} - \frac{b_{j-1} - a_j}{\underline{x}_{j-1} - \underline{x}_j} \underline{x}_j} \\
 &= b_{j-1} \frac{\underline{x}_j}{\underline{x}_{j-1}} \frac{a_j (1 + \underline{x}_{j-1} y_{j-1})}{b_{j-1} (1 + \underline{x}_j y_{j-1})} \\
 &= g_{j-1}.
 \end{aligned}$$

#### 4.2 Proof of action (4.2)

The action of  $r_{2j-1}$  on the dependent variables with the exception of  $x_j$  is trivial for any  $j = 1, \dots, n + 1$ . Hence we have to investigate the following actions:

$$\begin{aligned}
 &r_1(f_1), \quad r_1(g_1), \quad r_1(f_i) \quad (i \neq 1), \\
 &r_{2j-1}(f_{j-1}), \quad r_{2j-1}(g_{j-1}), \quad r_{2j-1}(f_j), \quad r_{2j-1}(g_j), \\
 &r_{2n+1}(f_n), \quad r_{2n+1}(g_n), \quad r_{2n+1}(f_i) \quad (i \neq n).
 \end{aligned}$$

We first consider the action  $r_{2j-1}(f_j)$ . It is described in terms of the variables  $x_i, y_i$  as

$$r_{2j-1}(f_j) = t \frac{x_j - \frac{a_j - b_j}{y_{j-1} - y_j} - x_{j+1}}{t x_{n+1} - x_1} = t \frac{(x_j - x_{j+1})(y_{j-1} - y_j) - (a_j - b_j)}{(t x_{n+1} - x_1)(y_{j-1} - y_j)}. \tag{4.4}$$

Substituting equation (3.15) to (4.4), we obtain

$$r_{2j-1}(f_j) = f_j \frac{(g_j - a_j) \frac{1}{f_j} + a_j (b_{j-1} \frac{1}{g_{j-1}} - 1) \frac{1}{f_{j-1}}}{(g_j - b_j) \frac{1}{f_j} + a_j (b_{j-1} \frac{1}{g_{j-1}} - 1) \frac{1}{f_{j-1}}} \quad (j = 1, \dots, n+1).$$

The other actions on the variables  $f_1, \dots, f_n$  can be shown in a similar way.

We next consider the action  $r_{2j-1}(g_j)$ . By using the system  $q$ - $\mathcal{P}_{(n+1, n+1)}$  twice, we can describe it in terms of the variables  $x_i, y_i$  as

$$\begin{aligned} r_{2j-1}(g_j) &= a_{j+1} \frac{x_{j+1} \left( 1 + \underline{x_j y_j} - \frac{a_j - b_j}{y_{j-1} - y_j} y_j \right)}{\left( \underline{x_j} - \frac{a_j - b_j}{y_{j-1} - y_j} \right) (1 + \underline{x_{j+1} y_j})} \\ &= g_j \frac{a_j \underline{x_j} (y_{j-1} - y_j)}{b_j (1 + \underline{x_j y_{j-1}}) - a_j (1 + \underline{x_j y_j})} \\ &= g_j \frac{(x_{j-1} - x_j)(y_{j-1} - y_j)}{(x_{j-1} - x_j)(y_{j-1} - y_j) - (a_j - b_j)(b_{j-1} \frac{1}{g_{j-1}} - 1)}. \end{aligned} \quad (4.5)$$

Then, substituting equation (3.15) to (4.5), we obtain

$$r_{2j-1}(g_j) = g_j \frac{(g_j - b_j) \frac{1}{f_j} + a_j (b_{j-1} \frac{1}{g_{j-1}} - 1) \frac{1}{f_{j-1}}}{(g_j - b_j) \frac{1}{f_j} + b_j (b_{j-1} \frac{1}{g_{j-1}} - 1) \frac{1}{f_{j-1}}} \quad (j = 1, \dots, n+1).$$

The other actions on the variables  $g_1, \dots, g_n$  can be shown in a similar way.

### 4.3 Proof of action (4.3)

The action  $\pi(f_i)$  is rewritten by the second equation of  $q$ - $\mathcal{P}_{(n+1, n+1)}$  and equation (3.6) to

$$\begin{aligned} \pi(f_i) &= \frac{q^2 \underline{y_i} - \underline{y_{i+1}}}{t \underline{y_0} - \underline{y_1}} = \frac{q^2 \underline{x_1}}{t \underline{x_{i+1}}} \frac{g_i f_i y_i - \frac{b_{i+1} f_{i+1}}{b_{i+1} - g_{i+1}} y_{i+1}}{\frac{g_0 t}{b_0 - g_0} y_0 - \frac{b_1 f_1}{b_1 - g_1} y_1} \quad (i = 1, \dots, n-1), \\ \pi(f_n) &= \frac{q^2 \underline{y_n} - \underline{y_{n+1}}}{t \underline{y_0} - \underline{y_1}} = q \frac{\underline{x_1} \frac{g_n f_n}{b_n - g_n} y_n - \frac{b_{n+1} t}{b_0 - g_0} y_0}{\underline{x_0} \frac{g_0 t}{b_0 - g_0} y_0 - \frac{b_1 f_1}{b_1 - g_1} y_1}. \end{aligned} \quad (4.6)$$

On the other hand, we obtain

$$\frac{\left(\frac{q}{i} - 1\right)(x_0 - x_1) \frac{1}{x_i}}{q} = \sum_{j=0}^{i-1} \frac{\frac{1}{q} \left(1 - \frac{a_{j+1}}{g_j}\right) (b_j - g_j)}{f_j} + \sum_{j=i}^n \frac{\left(1 - \frac{a_{j+1}}{g_j}\right) (b_j - g_j)}{f_j} \quad (i = 0, \dots, n), \quad (4.7)$$



from equation (3.7) and

$$\begin{aligned} \frac{(\frac{t}{q} - 1)(x_0 - x_1)}{t} y_i &= \sum_{j=0}^{i-1} \frac{\frac{t}{q} (1 - \frac{a_{j+1}}{g_j})(b_j - g_j)}{f_j} \\ &+ \frac{(\frac{t}{q} - \frac{a_{i+1}}{g_i})(b_i - g_i)}{f_i} + \sum_{j=i+1}^n \frac{(1 - \frac{a_{j+1}}{g_j})(b_j - g_j)}{f_j} \quad (i = 0, \dots, n), \end{aligned} \quad (4.8)$$

from equation (3.15). Substituting equations (4.7) and (4.8) to (4.6), we can show the action  $\pi(f_i)$ .

The action  $\pi(g_i)$  is rewritten by equation (3.6) to

$$\begin{aligned} \pi(g_i) &= \frac{b_{i+1}}{q^{\rho_1}} \frac{y_{i+1}(1 + x_i y_i)}{y_i(1 + x_{i+1} y_{i+1})} = \frac{a_{i+1}}{q^{\rho_1}} \frac{b_{i+1} f_{i+1} y_{i+1}}{b_{i+1} - g_{i+1} y_i} \quad (i = 1, \dots, n - 1), \\ \pi(g_n) &= \frac{b_{n+1}}{q^{\rho_1}} \frac{y_{n+1}(1 + x_n y_n)}{y_n(1 + x_{n+1} y_{n+1})} = \frac{a_{n+1}}{q^{\rho_1}} \frac{b_{n+1} t y_0}{b_n - g_n y_n}. \end{aligned} \quad (4.9)$$

Substituting equation (4.8) to (4.9), we can show the action  $\pi(g_i)$ .

### Acknowledgement

The author would like to express his gratitude to Professors Masatoshi Noumi and Yasuhiko Yamada for helpful comments and advice.

### Funding

This work was supported by JSPS KAKENHI Grant Number 15K04911.

### REFERENCES

1. FUJI, K. & SUZUKI, T. (2010) Drinfeld-Sokolov hierarchies of type  $A$  and fourth order Painlevé systems. *Funkcial. Ekvac.*, **53**, 143–167.
2. SUZUKI, T. & FUJI, K. (2012) Higher order Painlevé systems of type  $A$ , Drinfeld-Sokolov hierarchies and Fuchsian systems. *RIMS Kokyuroku Bessatsu*, **B30**, 181–208.
3. GARNIER, R. (1912) Sur des équations différentielles du troisième ordre dont l'intégrale est uniforme et sur une classe d'équations nouvelles d'ordre supérieur dont l'intégrale générale a ses points critiques fixés. *Ann. Sci. École Norm. Sup.*, **29**, 1–126.
4. JIMBO, M. & SAKAI, H. (1996) A  $q$ -analogue of the sixth Painlevé equation. *Lett. Math. Phys.*, **38**, 145–154.
5. KAWAKAMI, H. (2015) Matrix Painlevé systems, *J. Math. Phys.*, **56**, 033503.
6. KAJIWARA, K., NOUMI, M. & YAMADA, Y. (2002) Discrete dynamical systems with  $W(A_{m-1}^{(1)} \times A_{n-1}^{(1)})$  symmetry. *Lett. Math. Phys.*, **60**, 211–219.
7. MASUDA, T. (2003) On the rational solutions of  $q$ -Painlevé V equation. *Nagoya Math. J.*, **169**, 119–143.
8. MASUDA, T. (2015) A  $q$ -analogue of the higher order Painlevé type equations with the affine Weyl group symmetry of type  $D$ . *Funkcial. Ekvac.*, **58**, 405–430.
9. NAGAO, H. & YAMADA, Y. Study of  $q$ -Garnier system by Padé method. arXiv:1601.01099.
10. SAKAI, H. (2005) A  $q$ -analogue of the Garnier system. *Funkcial. Ekvac.*, **48**, 237–297.

11. SAKAI, H. Isomonodromic Deformation and 4-Dimensional Painlevé Type Equations. *UTMS*, **2010-17**, (Univ. of Tokyo, 2010) 1–21.
12. SASANO, Y. (2006) Higher order Painlevé equations of type  $D_l^{(1)}$ . *RIMS Koukyuroku*, **1473**, 143–163.
13. SUZUKI, T. (2010) A particular solution of a Painlevé system in terms of the hypergeometric function  ${}_{n+1}F_n$ . *SIGMA*, **6**, 078.
14. SUZUKI, T. (2013) A class of higher order Painlevé systems arising from integrable hierarchies of type A. *AMS Contemp. Math.*, **593**, 125–141.
15. SUZUKI, T. (2014) Six-dimensional Painlevé systems and their particular solutions in terms of rigid systems. *J. Math. Phys.*, **55**, 102902.
16. SUZUKI, T. (2015) A  $q$ -analogue of the Drinfeld-Sokolov hierarchy of type A and  $q$ -Painlevé system. *AMS Contemp. Math.*, **651**, 25–38.
17. TAKENAWA, T. (2003) Weyl group symmetry of type  $D_5^{(1)}$  in the  $q$ -Painlevé V equation. *Funkcial. Ekvac.*, **46**, 173–186.
18. TSUDA, T. (2010) On an integrable system of  $q$ -difference equations satisfied by the universal characters: its Lax formalism and an application to  $q$ -Painlevé equations. *Comm. Math. Phys.*, **293**, 347–359.
19. TSUDA, T. (2014) UC hierarchy and monodromy preserving deformation. *J. Reine Angew. Math.*, **690**, 1–34.
20. TSUDA, T. (2012) Hypergeometric solution of a certain polynomial Hamiltonian system of isomonodromy type. *Q. J. Math.*, **63**, 489–505.