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Experimental stark widths and shifts of V II spectral lines

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ABSTRACT

We have measured the Stark widths and shifts of V II spectral lines in the wavelength range 2000–4200 Å belonging to 75 multiplets. The spectra are emitted by laser-induced plasmas generated from fused glass discs prepared by borate fusion. The electron density and temperature are in the ranges (0.72–6.5) × 10^{17} cm^{-3} and (11 000–14 900) K, respectively. To avoid self-absorption, we have used seven samples with vanadium concentrations selected by the CSigma graph methodology. This has allowed to include strong and weak lines in the study, including resonance and forbidden lines. The experimental widths and shifts are compared with theoretical values available in the literature.

Key words: atomic data – line: profiles – plasmas.

1 INTRODUCTION

The knowledge of the broadening and shift produced by charged particles on the spectral lines is essential for spectroscopic diagnostic and atomic structure calculations. Therefore, accurate measurements of Stark broadening and shift parameters are required in laboratory plasma research, atomic structure calculations, and to analyse astrophysical data.

A large number of ionized vanadium spectral lines has been observed in stellar objects and solar plasma, being overabundant in some A-type stars (van't Veer-Mennert et al. 1985; Sadakane & Ueta 1989), where the Stark broadening is the main pressure broadening mechanism. In addition, some of these lines have been studied for the diagnostic of fusion plasmas because vanadium alloys are used in blanket applications.

Despite the interest of these applications, there are no accurate V II Stark width and shifts measurements and very few theoretical calculations. In 1983 and 1985, Lakićević made an estimation of the Stark broadening of some V II spectral lines only based on regularities and systematic trends. Until 2000 there have not been more sophisticated calculations, when Popović & Dimitrijević performed Stark broadening calculations for 14 transitions of VII. However, authors report that accuracy might be improved if reliable atomic data were available to perform the semiclassical method. Recently, Wood et al. (2014) have reported new experimental transition probability values for V II transitions, highlighting the need to measure simultaneously strong and very week lines in laboratory to explain unexpected trends in stellar abundance determinations. Also, Saloman & Kramida (2017) have critically evaluated atomic data for singly ionized vanadium.

The aim of this work is to provide for the first time experimental Stark width and shifts of V II lines. Our group has measured these parameters by using laser-induced breakdown spectroscopy (see for example Aguilera et al. 2014 and Manrique et al. 2019), developing a methodology that allows to perform simultaneously measurements for strong and weak spectral lines that are free of self-absorption through the selection of the suitable concentration of the studied element in the sample. In this work, we have studied Stark widths and shifts in the 2000–4200 Å wavelength range for strong and weak V II lines, including resonance and forbidden lines from 147 lines and 75 multiplets.

2 EXPERIMENT

A Q-switched Nd: YAG laser (wavelength 1064 nm, pulse width 4.5 ns, pulse energy 60 mJ) is used to generate the plasmas in air at atmospheric pressure. The focusing lens (focal length 126 mm) is placed at 122 mm from the sample surface. A pair of flat and concave (focal length 125 mm) mirrors is used to form a 1:1 image of the plasma on to the entrance of a Czerny–Turner spectrometer (focal length 0.75 m, grating of 3600 lines mm^{-1}, slit width 20 μm). Using an ICCD (1200 × 256 effective pixels), the system provides a spectral resolution of 15 pm at 3000 Å. The sample is rotated at 100 rpm during spectra measurements, which accumulate the emission from 100 laser shots. To reduce the statistical error, we repeated each measurement at five different sample locations, averaging the resulting line width values.

To avoid self-absorption, we select the atomic concentration of vanadium in the sample. By borate fusion of pure V_2O_5 in powder form, a set of seven samples were prepared as fused glass discs with

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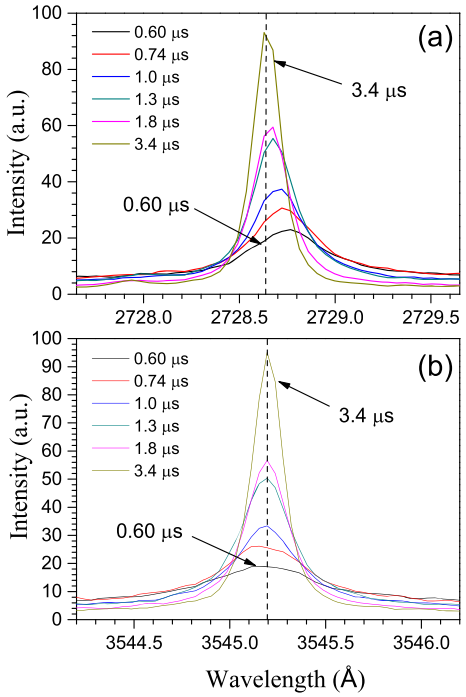

Figure 1. Spectra of the V II lines at 2728.637 Å (a) and 3545.196 Å (b), measured at different time windows centred at the indicated times.

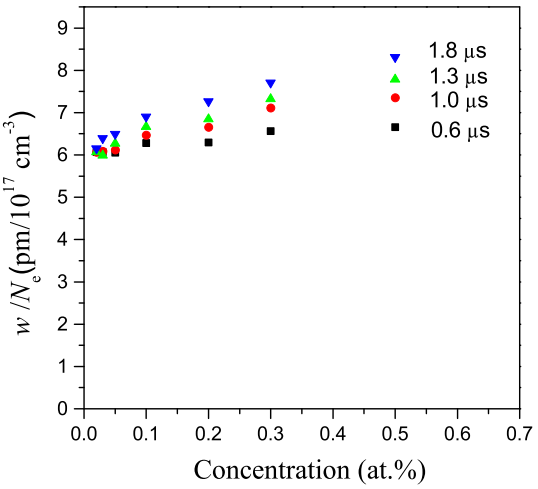

Figure 2. Ratio of the line width to the electron density against vanadium concentration for the V II line at 3102.300 Å measured at different time windows centred at the indicated times.

increasing vanadium concentration from 0.02 to 0.5 at. per cent. To determine the electron density from the Stark-broadened profile of a Ca II line, a trace amount of CaO was also added to each sample, leading to 0.008 at. per cent calcium concentration.

3 RESULTS AND DISCUSSION

For each line of interest, we measured spectra with time delays from the laser pulse ranging from 0.57 to 3.1 μs, whilst time widths varied correspondingly from 0.06 to 0.6 μs. Therefore, the resulting six time windows were centred at instants 0.60, 0.74, 1.0, 1.3, 1.8, and 3.4 μs. An additional spectrum with 7 μs delay is used as a reference to obtain the Stark shifts, as the electron density has decreased to a negligible value at this late time of the plasma evolution. Fig. 1 shows the temporal decrease of the widths for two spectral lines resulting from the drop of the plasma electron density. As can be seen, the line width is higher for the line at 3545.196 Å (Fig. 1b) whereas the line shift is high and positive for the line at 2728.637 Å (Fig. 1a) and smaller in absolute value and negative for the 3545.196 Å line (Fig. 1b).

In previous works (Manrique et al. 2019), the electron density was determined from the Stark broadening of the H α line, assuming that the measured value is valid for the whole set of samples. However, matrix effects could lead to a non-negligible variation of the electron density among the plasmas generated with the different samples. In this work, to be sure that the reference electron density and the profile of the line of interest are measured under the same conditions, we have determined the electron density in each measurement from the Stark width of the Ca II doublet at 3933.663 and 3968.469 Å. For the small calcium content in the samples (0.008 at. per cent), self-absorption of these lines is negligible. Moreover, for this low calcium concentration, the V II spectrum does not show overlapping with the Ca II spectrum. The electron density values deduced from the two Ca II lines are the same within the experimental error. For the Stark widths of these lines, we have used the experimental values reported by our group in a recent work (Aguilera et al. 2014) and checked again that the resulting electron densities are consistent with those determined using the H α line. The Lorentzian widths are obtained by fitting the experimental spectra to Voigt profiles with a fixed Gaussian component that includes the instrumental and Doppler profiles. The resulting electron density decreases from 6.5×10^{17} cm^{-3} at 0.6 μs to 0.72×10^{17} cm^{-3} at 3.4 μs. The plasma temperature at the different time windows is obtained from Boltzmann plots constructed using 22 V II lines with upper level energies in the range 4.29–6.88 eV and well-known transition probabilities (Kramida et al. 2019). The resulting temperature decreases from (14 900 $\pm$ 400) to (11 000 $\pm$ 100) K.

To reduce the systematic error due to self-absorption, we have used the method based on CSigma graphs developed in recent years by our group (Aragón & Aguilera 2014). Starting from a set of characteristic plasma parameters obtained in previous experiments performed in the same conditions, this method allows us to estimate the vanadium concentration required to limit self-absorption to a maximum value fixed at 10 per cent. In this way, the V II lines are classified in seven groups according to their intensity and, for each group, a sample with the required concentration is prepared and used for the Stark width measurements. Even though we have checked in many experiments that this procedure is reliable to control self-absorption, we have performed an additional test for the lines of higher intensity, consisting in the measurement of the line width for the whole set of samples prepared with different concentrations. Fig. 2 shows the ratio of the line width to the electron density as a function of the concentration in the sample for the most intense line among those investigated (3102.300 Å) measured at different time windows. As expected, the line width increases with concentration due to self-absorption. The effect is more pronounced at late-time windows at which the drop of the electron density leads to a decrease of the Lorentzian width, which results in stronger self-absorption. As can be seen, the data for different concentrations and time windows converge to the same w/N_e value at the 0.02 at. per cent concentration. This is an indication that, by using this small concentration, the effect

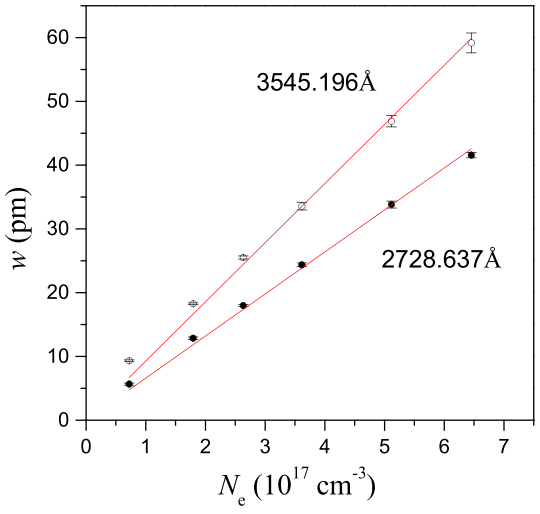

Figure 3. Line width as a function of the electron density for two V II lines. The error bars represent the standard deviation of the average for five repeated measurements. The final Stark widths are obtained as the slopes of the linear fittings with zero intercept.

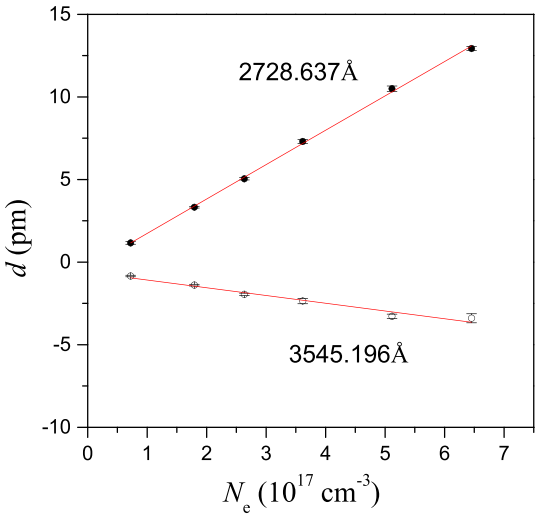

Figure 4. Line shift versus electron density for two V II lines. The error bars represent the standard deviation of the average for five repeated measurements. The Stark shifts are determined as the slopes of the linear plots with zero intercept.

of self-absorption has been reduced to a value lower than about 10 per cent for this intense line.

Fig. 3 shows the experimental results for the line width as a function of the electron density for the lines at 2728.637 and 3545.196 Å, obtained from the spectra of Fig. 1. The corresponding plots for the line shift are shown in Fig. 4. In both cases the data show a linear dependence with the electron density, and the final values of the Stark width and shift is obtained from the slopes of the linear fittings. Therefore, the weak dependence on temperature of the Stark width has not been observed for the temperature range in our plasmas taking into account our experimental errors.

Our results for the Stark widths and shifts of the V II lines are listed in Table 1, compared with previous results from the literature. To facilitate the location of the data for a given line we include Table 2, which collects the lines ordered by wavelength, showing the multiplet number for each line. Table S3, available online as Supporting Information, contains the list of measured Stark widths and shifts in ASCII code. The total experimental error is estimated as 15 per cent and 11 per cent for the Stark widths and shifts, respectively. It arises from the quadratic combination of the uncertainty due to the plasma inhomogeneity (3 per cent), the electron density measurement (11 per cent), and the self-absorption of the plasma (10 per cent only for width measurements). Nevertheless, for the shift measurements, the minimum shift detectable with our wavelength resolution leads to a minimum absolute error of 0.1 pm.

The comparison of our results with the data obtained by Popović & Dimitrijević (2000) shows good agreement for width measurements, which are 8 per cent higher for the resonance lines than ours and around 30 per cent higher for 4p–4s transitions. The disagreement increases to 60 per cent for 4d–4p transitions with lower values. Comparison with the same reference for shift measurements shows only good agreement, approximately 10 per cent, for 4d–4p transitions. Nevertheless, values for the resonance lines are lower by a factor 8 than ours whilst for the rest of the transitions are higher by a factor between 4 and 10, depending on the transition. It is worth noticing that shift calculations have lower accuracy than width values and that the authors of these calculations explained that, due to the lack of reliable atomic data, simpler methods for the calculations had to be applied. Width and shift Stark earlier estimations reported by Lakićević (1983 and 1985) using simple dependence on ionization potential, show good agreement (better than 30 per cent) to our width measurements whilst shift values have poor agreement.

CONCLUSIONS

We report measurements of Stark widths of V II spectral lines performed by laser-induced breakdown spectroscopy. Widths are provided for 147 lines, and shifts for 117 lines. Special care has been taken to avoid the systematic effect due to self-absorption, for which the CSigma graph methodology has been used to determine suitable vanadium concentrations in the samples. To our knowledge, these are the first experimental Stark width and shift data for VII, which are necessary for plasma diagnostics and for comparison with theoretical calculations.

ACKNOWLEDGEMENTS

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DATA AVAILABILITY

The data underlying this article are available in the article and in its online supplementary material.

Table 1. Stark widths and shifts (pm) of V II spectral lines at electron density 10^{17} cm^{-3}. Temperature ranged from 11 000 to 14 900 K. The estimated relative error is 15 and 11 per cent for width and shift, respectively, with a minimum absolute error for shifts of 0.1 pm.

No	Transition[a]	Multiplet[a]	λ(Å)[a]	*w*	*d*	*w*[b]	*d*[b]
1	$3d^4$–$3d^3(^4F)4p$	a^5D–z^3D^o	2702.177	6.8	1.80		
2		a^5D–z^5F^o	2700.927	6.4	2.33	7.00	0.257
			2706.156	6.2	2.11	7.00	0.257
			2715.655	6.5	2.04	7.00	0.257
			2728.637	6.6	2.09	7.00	0.257
3		a^5D-z^5D^o	2687.952	6.4	2.08	6.81	− 0.0792
			2679.316	6.5	1.80	6.81	− 0.0792
			2672.000	6.4	1.93	6.81	− 0.0792
			2677.796	6.4	1.90	6.81	− 0.0792
4	$3d^4$–$3d^3(^4P)4p$	a^5D–y^5D^o	2131.835	3.8			
5	$3d^3(^4F)4s$–$3d^3(^4F)4p$	a^5F–z^5G^o	3102.300	6.1	− 0.19	9.14	− 2.14
			3110.710	6.0	− 0.26	9.14	− 2.14
			3118.382	6.4	− 0.32	9.14	− 2.14
			3125.286	7.0	− 0.42	9.14	− 2.14
			3126.219	6.0	− 0.26	9.14	− 2.14
			3130.270	6.0	− 0.33	9.14	− 2.14
			3133.334	6.3	− 0.25	9.14	− 2.14
6		a^5F–z^3D^o	2906.458	6.9	− 0.44		
			2957.521	5.6	− 0.48		
			2896.206	6.1	− 0.54		
			2903.075	6.2	− 0.40		
7		a^5F–z^5F^o	2944.571	5.6	− 0.20	8.89	− 1.36
			2952.071	6.1	− 0.46	8.89	− 1.36
			2911.063	6.2		8.89	− 1.36
			2934.401	5.9	− 0.50	8.89	− 1.36
			2930.808	5.4	− 0.19	8.89	− 1.36
8		a^5F–z^5D^o	2880.028	5.6	− 0.46	8.63	− 1.72
			2882.499	6.0	− 0.35	8.63	− 1.72
			2889.619	6.6		8.63	− 1.72
9		a^3F–z^3D^o	3556.800	8.6	− 0.46	12.6	− 2.84
			3545.196	8.3	− 0.54	12.6	− 2.84
			3589.759	9.1	− 0.65	12.6	− 2.84
			3520.019	8.6	− 0.43	12.6	− 2.84
10		a^3F–z^5F^o	3593.333	8.1	− 0.31		
			3592.022	8.9	− 0.49		
			3530.772	8.5	− 0.40		
			3560.590	7.7			
			3566.176	8.5	− 0.60		
11		a^3F–z^5D^o	3517.300	8.7	− 0.43		
			3504.436	7.9	− 0.53		
			3493.163	8.3			
			3485.921	8.7	− 0.79		
12		a^3F–z^3G^o	3276.125	7.7	− 0.31	11.6	− 2.51
			3271.122	8.0	− 0.39	11.6	− 2.51
			3267.702	8.0	− 0.37	11.6	− 2.51
			3298.738	7.5	− 0.24	11.6	− 2.51
13		a^3F–z^3F^o	3190.682	8.6	− 0.23	11.1	− 2.32
			3188.513	8.8	− 0.58	11.1	− 2.32
			3187.712	8.5		11.1	− 2.32
			3214.746	9.3		11.1	− 2.32
			3208.346	8.4		11.1	− 2.32
			3164.839	7.8		11.1	− 2.32
14	$3d^3(^4F)4s$–$3d^3(^2G)4p$	a^3F–y^3G^o	2514.637	5.6			
15	$3d^3(^4F)4s$–$3d^3(^2P)4p$	a^3F–y^3D^o	2380.913	4.3			
16	$3d^4$–$3d^3(^4F)4p$	a^3P2–z^3D^o	3951.956	12.6	3.5		
17		a^3P2–z^5F^o	3973.629	12.8	3.4		
18		a^3P2–z^5D^o	3903.252	11.6	3.5		
19	$3d^4$–$3d^3(^2P)4p$	a^3P2–y^3P^o	2549.277	5.5	1.31		
20	$3d^4$–$3d^3(^4F)4p$	a^3H–z^3G^o	3715.464	11.2	3.2		
			3732.747	12.6	3.3		
			3745.799	11.9	3.1		
21	$3d^4$–$3d^3(^2G)4p$	a^3H–y^3G^o	2765.663	9.2	2.5		

Table 1 – *continued*

No	Transition[a]	Multiplet[a]	λ(Å)[a]	*w*	*d*	*w*[b]	*d*[b]
			2768.555	9.2	1.68		
22	$3d^4$–$3d^3(^2H)4p$	a^3H–x^3G^o	2352.178	4.7	1.73		
			2342.141	5.0	1.44		
23	$3d^4$–$3d^3(^4F)4p$	b^3F2–z^3F^o	3727.341	12.7	3.0		
			3770.966	12.5	3.3		
24	$3d^4$–$3d^3(^2G)4p$	b^3F2–y^3G^o	2836.516	7.2	1.99		
25		b^3F2–y^3F^o	2803.462	6.5	1.83		
			2802.792	7.9	1.81		
26	$3d^4$–$3d^3(^2D2)4p$	b^3F2–x^3F^o	2577.676	7.7	1.66		
27	$3d^3(^4P)4s$–$3d^3(^4P)4p$	a^5P–y^5D^o	2968.378	6.0			
			2983.562	6.2			
			3013.104	7.2	− 0.32		
			3022.580	6.8	− 0.52		
28	$3d^3(^4P)4s$–$3d^3(^4P)4p$	a^5P–z^5P^o	3001.204	5.8	− 0.26		
			3014.821	6.0	− 0.25		
			2988.025	6.2	− 0.32		
			2995.999	6.3	− 0.39		
29	$3d^4$–$3d^3(^4F)4p$	a^3G–z^3G^o	4005.702	12.1	1.96		
			4023.378	13.0	2.41		
			4035.620	13.6	2.42		
30		a^3G–z^3F^o	3878.707	13.0	2.30		
			3899.128	13.2	2.27		
			3914.321	13.5	2.07		
31	$3d^4$–$3d^3(^2G)4p$	a^3G–z^3H^o	3067.104	6.6	1.15		
32		a^3G–y^3F^o	2888.240	7.5	1.11		
33	$3d^4$–$3d^3(^2D2)4p$	a^3G–x^3F^o	2630.664	6.3	1.08		
			2642.206	6.7			
			2645.836	6.0	1.25		
34	$3d^4$–$3d^3(^2H)4p$	a^3G–x^3G^o	2465.270	5.2	1.04		
			2453.346	4.7	1.05		
35	$3d^3(^2G)4s$–$3d^3(^4F)4p$	b^3G–z^3F^o	4183.429	13.0	1.32		
36	$3d^3(^2G)4s$–$3d^3(^2G)4p$	b^3G–z^3H^o	3217.109	7.1	0.90		
			3237.870	8.6	1.01		
			3254.765	7.6	0.79		
37	$3d^3(^2G)4s$–$3d^3(^2G)4p$	b^3G–y^3G^o	3100.932	8.5	1.16		
38		b^3G–y^3F^o	3041.416	8.1	0.88		
			3042.261	7.8	0.95		
39		b^3G–z^1H^o	3023.884	7.3	0.85		
40		b^3G–z^1G^o	3012.016	7.8	0.93		
41	$3d^3(^2G)4s$–$3d^3(^2H)4p$	b^3G–y^3H^o	2797.012	7.1	0.73		
42	$3d^3(^2G)4s$–$3d^3(^2H)4p$	b^3G–x^3G^o	2584.951	5.7	0.89		
43	$3d^4$–$3d^3(^2G)4p$	a^1G2–z^1H^o	3155.404	9.0	1.73		
44		a^1G2–z^1G^o	3142.482	7.8	1.19		
45	$3d^4$–$3d^3(^2D2)4p$	a^3D–w^3D^o	2819.431	6.1	1.23		
			2817.495	6.4	1.25		
46	$3d^3(^2G)4s$–$3d^3(^2G)4p$	b^1G–z^1F^o	3282.529	7.9	0.34		
47		b^1G–z^1G^o	3265.886	9.4			
48	$3d^3(^2P)4s$–$3d^3(^4P)4p$	b^3P–z^3P^o	3621.208	11.5			
49	$3d^3(^2P)4s$–$3d^3(^2P)4p$	b^3P–z^3S^o	3024.983	6.6			
			3028.054	6.9			
50	$3d^3(^2P)4$ s–$3d^3(^2P)4p$	b^3P–x^3F^o	3005.813	6.9			
51		b^3P–x^3D^o	2972.263	7.2			
			2981.200	6.8			
52	$3d^4$–$3d^3(^2H)4p$	a^1I–x^3G^o	2775.760	7.1	2.6		
53	$3d^3(^4P)4s$–$3d^3(^4P)4p$	c^3P–z^3P^o	3787.239	10.5	0.35		
54	3d3(4P)4s–3d3(2P)4p	c3P–y3Do	3251.864	8.6	0.60		
			3257.886	8.4	0.20		
55	$3d^3(^4P)4s$–$3d^3(^4P)4p$	c^3P–x^3D^o	3083.209	8.6			
56	$3d^3(^2H)4s$–$3d^3(^2G)4p$	b^3H–z^3H^o	3669.423	9.3	− 0.27		
57	$3d^3(^2H)4s$–$3d^3(^2H)4p$	b^3H–y^3H^o	3136.517	7.1	− 0.35		
			3139.745	7.1	− 0.31		
58		b^3H–z^3I^o	3063.244	6.7			

Table 1 – *continued*

No	Transition[a]	Multiplet[a]	λ(Å)[a]	*w*	*d*	*w*[b]	*d*[b]
59		$b^3H–x^3G^o$	2869.132	6.3			
			2854.335	6.5			
			2847.572	6.4			
60		$b^3H–y^1H^o$	2845.244	8.4			
61	$3d^3(^2H)4s–3d^2(^3F)4s4p(^3P*)$	$b^3H–v^3G^o$	2023.565	3.1			
62	$3d^3(^2D2)4s–3d^3(^2D2)4p$	$b^3D–x^3F^o$	3151.316	7.7	1.33		
63	$3d^4–3d^3(^2G)4p$	$a^1D2–z^1F^o$	3497.027	10.1	1.58		
64	$3d^3(^2P)4s–3d^3(^2P)4p$	$a^1P–z^1S^o$	3847.334	10.1	− 1.40		
65		$a^1P–z^1D^o$	3618.927	10.2	− 1.17		
66		$a^1P–y^1P^o$	2949.176	6.3			
67	$3d^3(^2H)4s–3d^3(^2H)4p$	$a^1H–y^1G^o$	3250.773	7.9	− 0.65		
68		$a^1H–y^1H^o$	3113.564	7.9	− 0.35		
69	$3d^3(^2D2)4s–3d^3(^2D2)4p$	$b^1D–y^1F^o$	3337.823	9.7	1.98		
70	$3d^4–3d^3(^2H)4p$	$a^1F–y^1G^o$	3661.365	11.1	2.9		
71	$3d^3(^2F)4s–3d^3(^2F)4p$	$c^3F–w^3G^o$	2955.579	7.7	1.43		
72	$3d^4–3d^3(^2F)4p$	$d^3F1–w^3F^o$	3174.533	7.5			
73		$d^3F1–w^3G^o$	2973.971	6.9	1.29		
			2985.179	6.9	0.71		
74	$3d^3(^4F)4p–3d^3(^4F)4d$	$z^5G^o–e^5H$	2663.208	16.3	7.8	11.3	8.11
			2655.647	18.1	8.4	11.3	8.11
			2649.325	19.4	8.2	11.3	8.11
			2644.332	20.5	7.1	11.3	8.11
75		$z^3G^o–e^3H$	2781.399	22.5	11.6	14.7	9.87

[a]Data from Kramida et al. 2019.
[b]Popović & Dimitrijević (2000). Data reported for a temperature of 10 000 K.

Table 2. List of lines with reported Stark width or shift, ordered by wavelength and with multiplet number for faster location in Table 1.

λ (Å)	No
2023.565	61
2131.835	4
2342.141	22
2352.178	22
2380.913	15
2453.346	34
2465.270	34
2514.637	14
2549.277	19
2577.676	26
2584.951	42
2630.664	33
2642.206	33
2644.332	74
2645.836	33
2649.325	74
2655.647	74
2663.208	74
2672.000	3
2677.796	3
2679.316	3
2687.952	3
2700.927	2
2702.177	1
2706.156	2
2715.655	2
2728.637	2
2765.663	21
2768.555	21
2775.760	52
2781.399	75

Table 2 – *continued*

λ (Å)	No
2797.012	41
2802.792	25
2803.462	25
2817.495	45
2819.431	45
2836.516	24
2845.244	60
2847.572	59
2854.335	59
2869.132	59
2880.028	8
2882.499	8
2888.240	32
2889.619	8
2896.206	6
2903.075	6
2906.458	6
2911.063	7
2930.808	7
2934.401	7
2944.571	7
2949.176	66
2952.071	7
2955.579	71
2957.521	6
2968.378	27
2972.263	51
2973.971	73
2981.200	51
2983.562	27
2985.179	73
2988.025	28

Table 2 – *continued*

λ (Å)	No
2995.999	28
3001.204	28
3005.813	50
3012.016	40
3013.104	27
3014.821	28
3022.580	27
3023.884	39
3024.983	49
3028.054	49
3041.416	38
3042.261	38
3063.244	58
3067.104	31
3083.209	55
3100.932	37
3102.300	5
3110.710	5
3113.564	68
3118.382	5
3125.286	5
3126.219	5
3130.270	5
3133.334	5
3136.517	57
3139.745	57
3142.482	44
3151.316	62
3155.404	43
3164.839	13
3174.533	72
3187.712	13
3188.513	13
3190.682	13
3208.346	13
3214.746	13
3217.109	36
3237.870	36
3250.773	67
3251.864	54
3254.765	36
3257.886	54
3265.886	47
3267.702	12
3271.122	12
3276.125	12
3282.529	46
3298.738	12
3337.823	69
3485.921	11
3493.163	11
3497.027	63
3504.436	11
3517.300	11
3520.019	9
3530.772	10
3545.196	9
3556.800	9
3560.590	10
3566.176	10
3589.759	9
3592.022	10
3593.333	10
3618.927	65
3621.208	48

Table 2 – *continued*

λ (Å)	No
3661.365	70
3669.423	56
3715.464	20
3727.341	23
3732.747	20
3745.799	20
3770.966	23
3787.239	53
3847.334	64
3878.707	30
3899.128	30
3903.252	18
3914.321	30
3951.956	16
3973.629	17
4005.702	29
4023.378	29
4035.620	29
4183.429	35

REFERENCES

Aguilera J. A., Aragón C., Manrique J., 2014, MNRAS, 444, 1854

Aragón C., Aguilera J. A., 2014, JQSRT, 149, 90

Kramida A., , Ralchenko Yu., Reader J., and NIST ASD Team, 2019, NIST Atomic Spectra Database (ver. 5.7.1). National Institute of Standards and Technology, Gaithersburg, MD

Lakićević I. S., 1983, A&AS, 127, 37

Lakićević I. S., 1985, A&AS, 151, 457

Manrique J., Aragón C, Aguilera J. A., 2019, MNRAS, 482, 1931

Popović L. Č., Dimitrijević M. S., 2000, Phys. Scr., 61, 192

Sadakane K., Ueta M., 1989, Publ. Astron. Soc. Japan, 41, 279

Saloman E., Kramida A., 2017, ApJS, 231, 19

van't Veer-Mennert C., Coupry M. F., Bukhart C., 1985, A&AS, 146, 139

Wood M. P., Lawler J. E., Den Hartog E. A., Sneden C., Cowan J. J., 2014, ApJS, 214, 18

SUPPORTING INFORMATION

Supplementary data are available at *MNRAS* online.

Table S3. List of measured Stark widths and shifts, ordered by wavelength.

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