# Star formation rates in the L 1482 filament of the California molecular cloud 

Toshihiro Omodaka, ${ }^{1, *}$ Takumi Nagayama, ${ }^{2}$ Kazuhito Dobashi, ${ }^{3}$ James O. Chibueze, ${ }^{4,5}$ Akifumi Yamabi, ${ }^{3}$ Yoshito Shimajiri, ${ }^{6}$ Shinnosuke Inoue, ${ }^{1}$ Shota Hamada, ${ }^{1}$ Kazuyoshi Sunada, ${ }^{2}$ and Yuji Ueno ${ }^{2}$<br>${ }^{1}$ Department of Physics and Astronomy, Graduate School of Science and Engineering, Kagoshima University, 1-21-35 Korimoto, Kagoshima, Kagoshima 890-0065, Japan<br>${ }^{2}$ Mizusawa VLBI Observatory, National Astronomical Observatory of Japan, 2-21-1 Osawa, Mitaka, Tokyo 181-8588, Japan<br>${ }^{3}$ Department of Astronomy and Earth Sciences, Tokyo Gakugei University, Koganei, Tokyo 184-8501, Japan<br>${ }^{4}$ Centre for Space Research, Physics Department, North-West University, Potchefstroom 2520, South Africa<br>${ }^{5}$ Department of Physics and Astronomy, Faculty of Physical Sciences, University of Nigeria, Carver Building, 1 University Road, Nsukka, Nigeria<br>${ }^{6}$ National Astronomical Observatory of Japan, 2-21-1 Osawa, Mitaka, Tokyo 181-8588, Japan<br>*E-mail: omodaka.toshihiro@gmail.com

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#### Abstract

We measured the trigonometric parallax of the $\mathrm{H}_{2} \mathrm{O}$ maser source associated with the L1482 molecular filament hosting the most massive young star, LkH $\alpha$ 101, in the California molecular cloud. The measured parallax is $1.879 \pm 0.096 \mathrm{mas}$, corresponding to the distance of $532 \pm 28 \mathrm{pc}$. This parallax is consistent with that of the nearby star cluster $\operatorname{LkH} \alpha$ 101, which was recently measured with Gaia DR2. We found that the L 1482 molecular filament and the $\mathrm{LkH} \alpha 101$ cluster are located at the same distance within 3 $\pm 30 \mathrm{pc}$. We observed the southern parts of L 1482 molecular clouds including the $\mathrm{H}_{2} \mathrm{O}$ maser source, which is adjacent to $\mathrm{LkH} \alpha 101$, using the Nobeyama 45 m telescope in the $J=1-0$ transitions of both ${ }^{12} \mathrm{CO}$ and ${ }^{13} \mathrm{CO}$. The peak intensity of the ${ }^{12} \mathrm{CO}$ line revealed the high excitation temperature region ( $60-70 \mathrm{~K}$ ) due to heating by UV radiation from $\mathrm{LkH} \alpha 101$. We derived the column density of these molecular clouds assuming local thermodynamic equilibrium (LTE) from the ${ }^{13} \mathrm{CO}$ emission. Using Dendrogam, we searched for small-scale, dense structures (cores) and identified 337 cores in the ${ }^{13} \mathrm{CO}$ data. Gravitationally bound cores with a virial mass to LTE mass ratio $\leq 1.5$ and young stars are concentrated in the high excitation temperature region. The column density in the warm region is five to six times larger than that of the surrounding colder molecular region. This suggests that the warm region has been compressed by a high-pressure wave and successive radiation-driven star formation is in progress in this warm region. In the cold molecular cloud to the north of the warm region, the cores are likely gravitationally unbound, which may be the reason why star formation is not active there.


Key words: astrometry — ISM: clouds — masers — stars: formation — stars: individual (LkH $\alpha$ 101)

## 1 Introduction

The formation process of massive stars in giant molecular clouds is important but poorly understood phenomenon. As described by Lada, Lombardi, and Alves (2009), the California Molecular Cloud (CMC) rivals the well-known massive star-forming region, the Orion Molecular Cloud (OMC) as the most massive giant molecular cloud with similar mass, size, filamentary structure, and morphology. However, despite its large mass, the CMC is characterized by a star formation rate (SFR) that is an order of magnitude lower than that of the OMC. The most massive star that is forming in the CMC is an early B star, $\mathrm{LkH} \alpha 101$, embedded in a cluster of lower-mass young stars (Andrews \& Wolk 2008). The study of star formation in the CMC is very important due to its marked contrast to that of the OMC. The CMC is characterized by filamentary structures, with L 1482 being one of the prominent filaments in the region. $\mathrm{LkH} \alpha 101$ is a member of the embedded star cluster of young low-mass stars in the south-eastern part of L 1482 and is likely an early B star $\left(M \sim 15 M_{\odot}\right.$, Herbig et al. 2004). The most prominent feature is a dark lane of L1482 which runs across the southeastern corner of the $\mathrm{LkH} \alpha 101$ cluster (Andrews \& Wolk 2008). Since the $\mathrm{LkH} \alpha 101$ cluster is at the junction of two molecular filaments, L 1482 filament 1 and filament 2 on the plane of the sky, Li et al. (2014) suggested that the origin of the $\mathrm{LkH} \alpha 101$ cluster might be merging molecular filaments fed by converging inflows. However, Redman et al. (1986) argued that the dark lane is in the foreground and probably not associated with the $\mathrm{LkH} \alpha 101$ cluster. That argument is further supported by the interstellar $\mathrm{H}^{+}$chemistry constraints in the vicinity of an intense radiation source (Brittain et al. 2004). Accurate distances of both the $\mathrm{LkH} \alpha 101$ cluster and L 1482 filaments are required to understand star formation in the region.

Recently, the parallax of VLA J043001.15+337724.6, which is a member of the $\mathrm{LkH} \alpha 101$ cluster, was measured to be $1.87 \pm 0.10$ mas (corresponding to the distance of $532 \pm 28 \mathrm{pc}$ ) with Very Long Baseline Array (VLBA) radio continuum observations (Dzib et al. 2018). The astrometric result of Gaia DR2 (Gaia Collaboration 2018) also contains the measured parallax of the $\mathrm{LkH} \alpha 101$ cluster. $\mathrm{H}_{2} \mathrm{O}$ masers were detected in L 1482 filament 2 (Sunada et al. 2007), thus we made its parallax measurement with the VLBI astrometric observations to compare the distances of the $\mathrm{LkH} \alpha 101$ star cluster and the L 1482 molecular filament directly. If $\mathrm{LkH} \alpha 101$ is adjacent to L 1482 , $\mathrm{LkH} \alpha 101$ ( B star) and the intense radiation pressure from it will have a significant effect on the surrounding gas and form the $\mathrm{H}_{\text {II }}$ region. The cloud surface adjacent to the $\mathrm{H}_{\text {II }}$ region is photo-dissociated by the ultraviolet (UV) and farUV (FUV). Recently, Goicoechea et al. (2016) observed
the Orion Bright Bar photo-dissociated region (PDR) using the ALMA telescope; the warm cloud edge, exposed to energetic radiation from the Trapezium stars, revealed a fragmented ridge of high-density substructures and photoablative gas flows. They suggest that the cloud edge has been compressed by a high-pressure wave moving into the molecular cloud. The warm PDR adjacent to $\mathrm{LkH} \alpha 101$ is a similar physical region to the Orion Bar PDR. In this warm region, successive star formation may be triggered due to radiation-driven star formation.

In this paper, we study the star formation in the warm region facing to the $\mathrm{LkH} \alpha 101$ cluster. We observed molecular regions adjacent to $\mathrm{LkH} \alpha 101$ using the Nobeyama 45 m telescope in the $J=1-0$ transitions of ${ }^{12} \mathrm{CO}$ and ${ }^{13} \mathrm{CO}$. By using the ${ }^{13} \mathrm{CO}$ data, we derived the column density of the molecular clouds assuming local thermodynamic equilibrium (LTE). We executed Dendrogam analysis to identify gravitationally unstable dense cores. We explored the evolution of star formation in the California Molecular Cloud, We explored the evolution of star formation in the California Molecular Cloud, which Lada et al. (2009) referred to as a sleeping giant molecular cloud.

## 2 Observations and data reduction

### 2.1 VERA telescope

The VLBI observations of the $22.235080 \mathrm{GHz} \mathrm{H} \mathrm{H}_{2} \mathrm{O}$ maser in L1482 were carried out on 2010/333, 2011/016, 2011/113, 2011/228, 2011/297, and 2011/353 (year/day of the year) with VERA. The observing time was approximately eight hours at each epoch. VERA is a Japanese VLBI array, which consists of four 20 m diameter antennas in Mizusawa, Iriki, Ogasawara, and Ishigaki-jima, with 2300 km maximum and 1200 km minimum baseline lengths (Omodaka 2009) VERA has a dual-beam system, which allows simultaneously observing of adjacent sources (target maser and position reference sources within 2.2 separation) for the purpose of correcting for the tropospheric fluctuations (Kawaguchi et al. 2000). The target source, L1482, and the position reference source, $\mathrm{J} 0429+3319$, which is listed in the second realization of the International Celestial Reference Frame (ICRF2), were simultaneously observed using the dual-beam system of VERA. The separation angle between the target and the position reference source is 1.84 . The data correlation was carried out with the FX-type hardware correlator at the Mitaka campus of the National Astronomical Observatory of Japan (NAOJ) (Chikada et al. 1991). The phase tracking centers of L1482 and J0429+3319 are $(\alpha, \delta)_{\mathrm{J} 2000.0}=\left(04^{\mathrm{h}} 30^{\mathrm{m}} 27.4008,+35^{\circ} 09^{\prime} 17 . .649\right)$ and
$\left(04^{\mathrm{h}} 29^{\mathrm{m}} 52^{\mathrm{s}} \cdot 721121,+33^{\circ} 19^{\prime} 01^{\prime \prime} 85849\right)$, respectively. The accumulation time of the correlation process was set to 1 s . The frequency and velocity spacing of L1482 are 15.625 kHz and $0.21 \mathrm{~km} \mathrm{~s}^{-1}$, respectively. The total bandwidth and the frequency spacing of J0429+3319 are 240 MHz and 0.25 MHz , respectively. Data reduction was done using VERA Data Analyzer (VEDA), which is a software package developed by the Mizusawa VLBI Observatory of NAOJ. The fringe search of J0429+3319 was done with an integration time of 128 s and a time interval of 32 s . The fringe and the self-calibration solutions of J0429+3319 were applied accordingly to L1482. The peak intensity in the self-calibration map of $\mathrm{J} 0429+3319$ is approximately $0.07 \mathrm{Jy} \mathrm{beam}^{-1}$. The instrumental delay and phase between the dual-beam is calibrated by the horn-on-dish method using the artificial noise source and the phase correction detector (Honma et al. 2008b). The tropospheric delay is calibrated using the tropospheric zenith delay measured by GPS (Honma et al. 2008a) and the tropospheric mapping function (Niell 1996). The ionospheric delay is calibrated using the Global Ionosphere Map (GIM) produced by the University of Bern.

### 2.2 CO observations with the Nobeyama 45 m radio telescope

Using the Nobeyama 45 m radio telescope, the ${ }^{12} \mathrm{CO}(J=$ $1-0$ ) observations towards the $13^{\prime} \times 12^{\prime}$ region around L1482 was carried out between 2013 March and May, while the ${ }^{13} \mathrm{CO}(J=1-0)$ observations covering the $40^{\prime} \times 40^{\prime}$ region in the southern part of L 1482 were done between 2013 December and 2014 January. We used the new twobeam, two-polarization, side-band separating receiver, TZ, for the ${ }^{12} \mathrm{CO}(J=1-0)$ observations, and the 25 -element focal plane receiver, BEARS, installed in the 45 m telescope for the ${ }^{13} \mathrm{CO}(J=1-0)$ observations. At 110 GHz , the telescope has a beam size of $16^{\prime \prime}$ (half power beam width, HPBW). The beam separation of BEARS is 41.11 on the plane of the sky (Sunada et al. 2000). The typical noise temperature was $\sim 150 \mathrm{~K}$ for the TZ receiver and 300 K for BEARS. The on-the-fly (OTF) technique was adopted to map each observation box. The temperature scale was determined by the chopper-wheel method. The telescope pointing was checked every hour by observing SiO maser line from Orion KL. The parameters of the observations are summarized in table 1 .

## 3 Results

### 3.1 Annual parallax

A maser spot at $V_{\text {LSR }}=7.38 \mathrm{~km} \mathrm{~s}^{-1}$ was detected over all six VLBI observation epochs. The peak intensities

Table 1. Observation parameters with the Nobeyama 45 m telescope.

| Molecular line | ${ }^{12} \mathrm{CO}(J=1-0)$ | ${ }^{13} \mathrm{CO}(J=1-0)$ |
| :--- | :---: | :---: |
| Rest frequency $(\mathrm{GHz})$ | 115.271202 | 110.201353 |
| Beam size (") | 14 | 15 |
| Receiver | TZ | BEARS |
| Bandwidth (MHz) | 63 | 32 |
| Frequency resolution $(\mathrm{kHz})^{\text {Velocity resolution }\left(\mathrm{km} \mathrm{s}^{-1}\right)}$ | 15.26 | 35.75 |
| Map gird size | 0.05 | 0.1 |
| Noise level (K) | $7!^{\prime \prime} 5$ | $7!\prime 5$ |
| Map coverage | 2.53 | 1.17 |

from the first to the last epochs were 2.37, 1.34, $1.80,2.73,9.21$, and $7.24 \mathrm{Jy} \mathrm{beam}^{-1}$. Figure 1 (left-hand panel) shows the phase-referenced map of this spot. The detected position of this $\mathrm{H}_{2} \mathrm{O}$ maser spot, $(\alpha, \delta)_{\text {J2000.0 }}=$ $\left(04^{\mathrm{h}} 30^{\mathrm{m}} 27^{\mathrm{s}} 4008,+35^{\circ} 09^{\prime} 17^{\prime \prime} 649\right)$, is consistent within 0.2 with that of the class I young stellar object (YSO) J04302741+3509178 (Broekhoven-Fiene et al. 2015) in L 1482 filament 2 (Li et al. 2014), the bolometric luminosity of which is $9.26 L_{\odot}$ (Harvey et al. 2013). The $\mathrm{H}_{2} \mathrm{O}$ maser is most likely associated with this YSO. The positional variation of this maser spot during the six epochs is shown in figure 2 . Least-square fits were made to these positions, with the parallax, as well as the proper motion. The parallax and the proper motion were obtained to be $\pi=1.879 \pm 0.096 \mathrm{mas}(D=532 \pm 28 \mathrm{pc})$ and $\left(\mu_{\alpha} \cos \delta, \mu_{\delta}\right)$ $=(3.17 \pm 0.16,-8.55 \pm 0.16) \mathrm{mas} \mathrm{yr}^{-1}$, respectively. Systematic errors in (RA, Dec) of ( $0.127,0.138$ ) mas were added in quadrature to the formal errors calculated by $\sigma_{\text {formal }}=\theta_{\text {beam }} /(2 \times S N R)$ to achieve the reduced $\chi^{2}=1$.

### 3.2 CO excitation temperature

In order to demonstrate the observed ${ }^{12} \mathrm{CO}(J=1-0)$ spectra, we divided the $13^{\prime} \times 12^{\prime}$ observed region into 10 $\times 10$ grids, and in each grid we averaged the data. Figure 3 shows the averaged spectra. The peak intensity is 45 K and the average over the cloud is 23 K , and we show the ${ }^{12} \mathrm{CO}$ intensity distribution around $\mathrm{LkH} \alpha 101$ in figure 4 . We estimated the excitation temperature $T_{\text {ex }}$ on the assumption that the rotational levels of gas are in the LTE and that the ${ }^{12} \mathrm{CO}(J=1-0)$ emission line is optically thick. $T_{\text {ex }}$ is then given by
$T_{\text {ex }}=\frac{5.53}{\ln \left[1+5.53 /\left(T_{\text {peak }}+0.819\right)\right]} \mathrm{K}$,
where $T_{\text {peak }}$ is the peak intensity of the ${ }^{12} \mathrm{CO}$ line in Kelvin. The average $T_{\text {ex }}$ in this region is higher than 30 K (Kong et al. 2015) obtained over the vast majority of the CMC.


Fig. 1. Phase-referencing $\mathrm{H}_{2} \mathrm{O}$ maser map of L 1482 (left) and self-calibrated image of the position reference source (J0429+3319) (right) from the 2011 day-353 epoch.


Fig. 2. Parallax and proper motion of the $\mathrm{H}_{2} \mathrm{O}$ maser spot at $V_{\mathrm{LSR}}=7.38 \mathrm{~km} \mathrm{~s}^{-1}$ in L 1482 . (Left) Positions on the sky. (Middle) RA (black circles) and Dec (gray circles) offset versus time. (Right) Same as the middle panel, but the proper motion of the target has been subtracted, leaving only the effect of the parallax.

The highest excitation temperatures obtained in the CMC are nearly $\sim 70 \mathrm{~K}$.

Figure 5 shows the spatial distribution of the ${ }^{13} \mathrm{CO}$ $(J=1-0)$ spectra which we obtained by dividing the $40^{\prime} \times 40^{\prime}$ observed region into $10 \times 10$ grids and by averaging the data in each grid. It is noted that the presence of an arc-shaped region (southward) with high excitation temperatures centering around $\operatorname{LkH} \alpha 101$, and we show the ${ }^{13} \mathrm{CO}$ intensity distribution in figure 6 . A further interesting feature is the presence of a single selfabsorption over the region facing the $\mathrm{LkH} \alpha 101$ cluster. Figure 7 shows the comparison of the ${ }^{12} \mathrm{CO}$ and ${ }^{13} \mathrm{CO}$ spectra. Clear dips in the ${ }^{12} \mathrm{CO}$ spectra can be recognized
in the velocity range from $-1 \mathrm{~km} \mathrm{~s}^{-1}$ to $+1 \mathrm{~km} \mathrm{~s}^{-1}$. The velocities of these dips coincide with the peak velocities of the ${ }^{13} \mathrm{CO}$ spectra. This suggests that the ${ }^{12} \mathrm{CO}$ emission emitted from the warm region around $\mathrm{LkH} \alpha 101$ is absorbed by the colder gas in the foreground, because the line is optically thick.

As shown in the next subsection, we estimated the column densities of ${ }^{13} \mathrm{CO}, \mathrm{N}_{\mathrm{li}^{3} \mathrm{CO}}$, on the assumption of LTE. We used $T_{\text {ex }}$ derived from the ${ }^{12} \mathrm{CO}$ data. However, the area mapped in ${ }^{13} \mathrm{CO}$ is much larger than that mapped in ${ }^{12} \mathrm{CO}$. Li et al. (2014) carried out ${ }^{12} \mathrm{CO}(J=3-2)$ observations toward L1482, and covered an extent similar to our ${ }^{13} \mathrm{CO}$ map. They found that the strong ${ }^{12} \mathrm{CO}(J=3-2)$


Fig. 3. Distribution of the observed ${ }^{12} \mathrm{CO}$ emission. The region mapped in ${ }^{12} \mathrm{CO}$ is $13^{\prime} \times 12^{\prime}$ as shown in figure 4 . We divided the region into $10 \times 10$ grids and calculated the average spectra, in each grid.


Fig. 4. Integrated intensity map of the ${ }^{12} \mathrm{CO}(J=1-0)$ emission line towards the L 1482 molecular cloud. The velocity used for the integration is from $-5 \mathrm{~km} \mathrm{~s}^{-1}$ to $+5 \mathrm{~km} \mathrm{~s}^{-1}$, and the size of map is $13^{\prime} \times 12^{\prime}$. The contours start from $50 \mathrm{Kkm} \mathrm{s}^{-1}$ with an increment of $25 \mathrm{Kkm} \mathrm{s}^{-1}$.
emission is detected around the cloud (see, their figure 2) even in regions not covered by our ${ }^{12} \mathrm{CO} J=1-0$ observations. Since the energy necessary to excite the $J=3-2$ transition is $h v / k \simeq 32 \mathrm{~K}$ in temperature, the regions emitting the molecular line should have an excitation temperature higher than this value. For regions not covered in our ${ }^{12} \mathrm{CO}$ observations, we therefore assumed a uniform $T_{\mathrm{ex}}$ of 35 K .

### 3.3 Dense cores in L 1482 filament: core identification and mass derivation

In order to study the star formation activity in the L1482 filament, we identified cores using the fiducial core-finding algorithm Dendrogram, (Rosolowsky et al. 2008) implemented in the astrodendro python package. Dendrogram is a tree diagram that represents the hierarchy of structures


Fig. 5. Distribution of the observed ${ }^{13} \mathrm{CO}$ emission. The region mapped in ${ }^{13} \mathrm{CO}$ is $40^{\prime} \times 40^{\prime}$ as shown in figure 6 . We divided the region into $10 \times 10$ grids and calculated the average spectra, in each grid.


Fig. 6. Integrated intensity map of ${ }^{13} \mathrm{CO}(J=1-0)$ emission line towards the L 1482 molecular cloud. The velocity range used for the integration is $-2.9 \mathrm{~km} \mathrm{~s}^{-1}$ to $+2.9 \mathrm{~km} \mathrm{~s}^{-1}$, and the size of map is $40^{\prime} \times 40^{\prime}$. The contours start from $20 \mathrm{~K} \mathrm{~km} \mathrm{~s}^{-1}$ with an increment of $10 \mathrm{Kkm} \mathrm{s}^{-1}$.
within given data. It is composed of two types of structure, "branches" and "leaves." Branches are the structures that split into multiple sub-structures, and the leaves are the structures that have no sub-structure. In this study, we define the smallest structure "leaf" as a core, and focus on the detection of cores. We used the three-dimensional
${ }^{13} \mathrm{CO}$ data cube, which is an appropriate tracer to analyze internal structures through the entire molecular cloud, because the ${ }^{13} \mathrm{CO}$ emission line is moderately optically thin in the observed region, and is a good probe of the column density. To identify the cores, we applied Dendrogram to the ${ }^{13} \mathrm{CO}$ cube with parameters min_value of $5 \sigma$, min_delta


Fig. 7. Comparison of the ${ }^{12} \mathrm{CO}$ (dotted lines) and the ${ }^{13} \mathrm{CO}$ (solid lines) emission lines around $\mathrm{LkH} \alpha 101$. We divided the $13^{\prime} \times 12^{\prime}$ area observed in ${ }^{12} \mathrm{CO}$ into $5 \times 6$ grids, and calculated the ${ }^{12} \mathrm{CO}$ and ${ }^{13} \mathrm{CO}$ spectra in each grid.
of $3 \sigma$, and min_npix of 5.2 pixels $\left(=A_{\theta_{\text {beam }}} / A_{\text {pixel }}\right)$. We note that min_npix is the minimum number of pixels that the structures must contain. We used the signal-to-noise ratio map to avoid detection of fake sources due to the nonuniform noise distribution. The application of Dendrogram with these parameters to the ${ }^{13} \mathrm{CO}$ data cube resulted in the identification of 337 cores in a square region of $\sim 40^{\prime} \times 40^{\prime}$.

We derived the physical parameters of the cores assuming the LTE using the standard procedures (e.g., Wilson et al. 2009). Using $T_{\text {ex }}$ derived based on equation (1), we calculated the optical depth of the ${ }^{13} \mathrm{CO}$ line as a function of the coordinates and velocity as
$\tau_{13 \mathrm{CO}}(\alpha, \delta, V)=-\ln \left\{1-\frac{T_{13} \mathrm{CO}}{5.29\left[J\left(T_{\mathrm{ex}}\right)-0.164\right]}\right\}$,
where $T_{13} \mathrm{CO}$ is the brightness temperature of the ${ }^{13} \mathrm{CO}$ line and $J\left(T_{\text {ex }}\right)=5.29 /\left\{\exp \left(5.29 / T_{\text {ex }}\right)-1\right\}$. We calculated the ${ }^{13} \mathrm{CO}$ column density as
$N^{{ }^{13} \mathrm{CO}}(\alpha, \delta)=2.42 \times 10^{14}\left[\frac{T_{\mathrm{ex}} \int \tau_{13} \mathrm{CO} d V}{1-\exp \left(-5.29 / T_{\mathrm{ex}}\right)}\right] \mathrm{cm}^{-2}$.

The total molecular column density can be estimated as $N\left(\mathrm{H}_{2}\right)=5 \times 10^{5} \mathrm{~N}^{13} \mathrm{CO}$ (Dickman 1978).

We estimated the radius of each core assuming that the cores are spherical as

$$
\begin{equation*}
R_{\mathrm{core}}=\sqrt{S / \pi} \tag{4}
\end{equation*}
$$

where $S$ is the projected area of the core. The mass of each core was derived based on the following equation:
$M_{\mathrm{LTE}}=\mu m_{\mathrm{H}} \Sigma N\left(\mathrm{H}_{2}\right) \Delta x^{2}$,
where $\mu$ is the mean molecular weight $2.8, m_{\mathrm{H}}$ is the proton mass, and $\Delta x^{2}\left(=7 .!5 \times 7.15 \simeq 3.5 \times 10^{33} \mathrm{~cm}^{2}\right.$ at 532 pc$)$ is the pixel size. The summation was performed over each core within the boundary. Based on the mass and radius derived above, we estimated the mean molecular number density of the cores as
$n=\frac{3 M_{\mathrm{LTE}}}{4 \pi R_{\text {core }}^{3} \mu m_{\mathrm{H}}} \mathrm{H}_{2} \mathrm{~cm}^{-3}$.
As we discuss in the next subsection, virial masses are useful to investigate the gravitational stability of the cores. We estimated the virial mass of each core as
$M_{\text {vir }}=\frac{3 a^{-1} R_{\text {core }} \Delta V_{\text {comp }}^{2}}{8(\ln 2) G}$,
where $G$ is the gravitational constant and $a$ is the shape factor taken to be $3 / 5$ corresponding to a centrally condensed density distribution with $\rho \propto r^{-2}$ (Bertoldi $\&$ McKee 1992). $\Delta V_{\text {comp }}$ is the line width defined at full width at halfmaximum of the composite spectrum averaged over each core, which we measured by a Gaussian fit. We further estimated the virial ratio $\alpha_{\text {vir }}$ which is the ratio of the virial mass to the core mass, i.e.,
$\alpha_{\text {vir }}=\frac{3 a^{-1} R_{\text {core }} \Delta V_{\text {comp }}^{2}}{8(\ln 2) G M_{\text {LTE }}}$.
In table 2, we list the identified 337 cores in order of right ascension (J2000.0) and present their physical properties derived in the above. Figure 8 shows their locations. The radius of the identified cores ranges from 0.029 to 0.180 pc
Table 2. Properties of identified cores using dendrogram.

Table 2. (Continued)

| ID | RA | Dec | $\begin{gathered} R_{\text {core }} \\ {[\mathrm{pc}]} \end{gathered}$ | $\begin{aligned} & T_{\text {peak }} \\ & {[\mathrm{K}]} \end{aligned}$ | Background level [K] | $\begin{gathered} \Delta \mathrm{V} \\ {\left[\mathrm{~km} \mathrm{~s}^{-1}\right]} \end{gathered}$ | $\begin{gathered} M_{\mathrm{vir}} \\ {\left[M_{\odot}\right]} \end{gathered}$ | $\underset{\left[\times 10^{23} \mathrm{~cm}^{-2}\right]}{\Sigma N\left(\mathrm{H}_{2}\right)}$ | $\begin{aligned} & M_{\mathrm{LTE}} \\ & {\left[M_{\odot}\right]} \end{aligned}$ | $\stackrel{n}{\left[\times 10^{4} \mathrm{~cm}^{-3}\right]}$ | $\begin{aligned} & M_{\text {vir }} / M_{\text {LTE }} \\ & \quad\left(=\alpha_{\text {vir }}\right) \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 39 | $4^{\mathrm{h}} 29^{\mathrm{m}} 488^{\text {s }} 73$ | $35^{\circ} 26^{\prime} 50.13$ | 0.118 | 13.542 | 10.2 | 0.664 | 10.88 | 13.00 | 10.81 | 2.27 | 1.01 |
| 40 | $4^{\mathrm{h}} 29^{\mathrm{m}} 499.15$ | 35 ${ }^{\circ} 15^{\prime} 29.99$ | 0.045 | 26.501 | 8.3 | 1.135 | 12.12 | 11.07 | 9.20 | 34.80 | 1.32 |
| 41 | $4^{\mathrm{h}} 29^{\mathrm{m}} 4995$ | $35^{\circ} 11^{\prime} 38.13$ | 0.050 | 13.723 | 8.7 | 1.892 | 37.43 | 6.69 | 5.56 | 15.34 | 6.73 |
| 42 | $4^{\mathrm{h}} 29^{\mathrm{m}} 5055$ | $35^{\circ} 12^{\prime} 56.17$ | 0.035 | 23.538 | 7.3 | 1.613 | 18.77 | 5.76 | 4.79 | 38.48 | 3.92 |
| 43 | $4^{\mathrm{h}} 29^{\mathrm{m}} 51{ }^{\text {s }}$ 41 | $35^{\circ} 15^{\prime} 24.9$ | 0.039 | 30.578 | 13.6 | 1.223 | 12.30 | 7.61 | 6.32 | 36.72 | 1.95 |
| 44 |  | $35^{\circ} 28^{\prime} 20^{\prime \prime} 4$ | 0.044 | 10.939 | 9.2 | 0.722 | 4.75 | 1.59 | 1.32 | 5.36 | 3.59 |
| 45 | $4^{\mathrm{h}} 29^{\mathrm{m}} 53.47$ | $35^{\circ} 13^{\prime} 03.5$ | 0.049 | 25.132 | 20.1 | 1.504 | 23.08 | 11.57 | 9.61 | 28.16 | 2.40 |
| 46 | $4^{\mathrm{h}} 29^{\mathrm{m}} 54.33$ | $35^{\circ} 11^{\prime} 23.14$ | 0.056 | 15.220 | 27.6 | 1.390 | 22.48 | 7.54 | 6.27 | 12.30 | 3.59 |
| 47 | $4^{\mathrm{h}} 29^{\mathrm{m}} 54.435$ | 35 ${ }^{\circ} 17^{\prime} 23.19$ | 0.036 | 15.341 | 7.3 | 1.327 | 13.32 | 3.04 | 2.53 | 18.66 | 5.27 |
| 48 | $4^{\mathrm{h}} 29^{\mathrm{m}} 55927$ | $35^{\circ} 15^{\prime} 46.15$ | 0.035 | 26.291 | 6.8 | 1.730 | 21.60 | 7.62 | 6.34 | 50.93 | 3.41 |
| 49 | $4^{\mathrm{h}} 29^{\mathrm{m}} 55928$ | $35^{\circ} 14^{\prime} 05^{\prime \prime} 0$ | 0.042 | 29.807 | 7.5 | 1.444 | 18.42 | 12.30 | 10.22 | 47.54 | 1.80 |
| 50 |  | $35^{\circ} 14^{\prime} 08.17$ | 0.031 | 30.101 | 12.5 | 1.523 | 14.97 | 7.46 | 6.20 | 71.73 | 2.41 |
| 51 | $4^{\text {h }} 29^{\text {m }} 7{ }^{\text {5 }}$ ¢ 77 | $35^{\circ} 12^{\prime} 26.17$ | 0.039 | 20.297 | 11.5 | 1.249 | 12.84 | 4.92 | 4.09 | 23.75 | 3.14 |
| 52 | $4^{\text {h } 29}{ }^{\text {m }} 99574$ | $35^{\circ} 14^{\prime} 08.15$ | 0.035 | 23.588 | 12.1 | 2.320 | 38.83 | 9.69 | 8.05 | 64.71 | 4.82 |
| 53 | $4^{\mathrm{h}} 30^{\mathrm{m}} 00579$ | $35^{\circ} 09^{\prime} 10.5$ | 0.084 | 12.872 | 11.6 | 1.244 | 27.11 | 13.24 | 11.01 | 6.40 | 2.46 |
| 54 | $4^{\mathrm{h}} 30^{\mathrm{m}} 01924$ | $35^{\circ} 13^{\prime} 04.9$ | 0.038 | 21.199 | 13.2 | 1.506 | 17.93 | 6.08 | 5.05 | 31.73 | 3.55 |
| 55 | $4^{\mathrm{h}} 30^{\mathrm{m}} 02^{\text {s }}$. 98 | $35^{\circ} 11^{\prime} 52.16$ | 0.084 | 19.118 | 10.2 | 1.436 | 36.13 | 26.57 | 22.09 | 12.84 | 1.64 |
| 56 | $4^{\mathrm{h}} 30^{\mathrm{m}} 03.95$ | $35^{\circ} 13^{\prime} 37.10$ | 0.036 | 30.040 | 9.9 | 1.501 | 17.05 | 11.76 | 9.77 | 72.18 | 1.74 |
| 57 | $4^{\mathrm{h}} 30^{\mathrm{m}} 04.530$ | $35^{\circ} 07^{\prime 2} 25^{\prime \prime} 4$ | 0.031 | 11.178 | 8.6 | 1.853 | 22.16 | 2.13 | 1.77 | 20.49 | 12.51 |
| 58 | $4^{\mathrm{h}} 30^{\mathrm{m}} 04.56$ | $35^{\circ} 07^{\prime} 57^{\prime \prime} 4$ | 0.046 | 13.123 | 8.5 | 1.443 | 20.17 | 4.43 | 3.68 | 13.04 | 5.48 |
| 59 | $4^{\mathrm{h}} 30^{\mathrm{m}} 055^{5} 01$ | 35 ${ }^{\circ} 09^{\prime} 59.5$ | 0.050 | 14.954 | 9.8 | 1.304 | 17.79 | 6.46 | 5.37 | 14.80 | 3.31 |
| 60 | $4^{\mathrm{h}} 30^{\mathrm{m}} 05950$ | $35^{\circ} 38^{\prime} 10.9$ | 0.039 | 8.183 | 10.1 | 1.206 | 11.95 | 1.53 | 1.27 | 7.37 | 9.42 |
| 61 | $4^{\text {h }} 30^{\mathrm{m}} 055^{5} 78$ | $35^{\circ} 10^{\prime} 31.11$ | 0.046 | 17.389 | 8.9 | 1.285 | 15.99 | 6.34 | 5.27 | 18.65 | 3.03 |
| 62 | $4^{\text {h }} 30^{\text {m }} 066^{5} 29$ | $35^{\circ} 11^{\prime} 24.15$ | 0.031 | 21.277 | 7.9 | 1.490 | 14.32 | 4.57 | 3.80 | 43.90 | 3.77 |
| 63 | $4^{\mathrm{h}} 30^{\mathrm{m}} 06.9{ }^{\text {s }} 73$ | $35^{\circ} 13^{\prime} 30.17$ | 0.038 | 28.752 | 9.8 | 1.797 | 25.52 | 13.64 | 11.34 | 71.21 | 2.25 |
| 64 | $4^{\mathrm{h}} 30^{\mathrm{m}} 06.9 .95$ | $35^{\circ} 15^{\prime} 50.11$ | 0.079 | 8.542 | 13.5 | 1.549 | 39.45 | 7.67 | 6.38 | 4.46 | 6.19 |
| 65 | $4^{\mathrm{h}} 30^{\mathrm{m}} 07{ }^{\text {s }}$. 12 | $35^{\circ} 9^{\prime} 05^{\prime \prime} 3$ | 0.075 | 8.386 | 13.3 | 1.350 | 28.51 | 7.11 | 5.91 | 4.82 | 4.83 |
| 66 | $4^{\mathrm{h}} 30^{\mathrm{m}} 088^{\text {s }} 14$ | $35^{\circ} 07^{\prime} 51.7$ | 0.060 | 12.750 | 10.0 | 1.497 | 28.02 | 7.42 | 6.17 | 9.84 | 4.54 |
| 67 | $4^{\text {h }} 30^{\mathrm{m}} 08.8{ }^{\text {s }} 37$ | $35^{\circ} 13^{\prime} 04.0$ | 0.035 | 27.904 | 24.1 | 1.736 | 21.75 | 9.17 | 7.62 | 61.23 | 2.85 |
| 68 | $4^{\text {h }} 30^{\mathrm{m}} 08.598$ | $35^{\circ} 11^{\prime} 11{ }^{\prime \prime} 5$ | 0.035 | 22.539 | 14.4 | 1.225 | 10.82 | 4.63 | 3.85 | 30.91 | 2.81 |
| 69 | $4^{\text {h }} 30^{\text {m }} 099.35$ | $35^{\circ} 40^{\prime} 53.16$ | 0.082 | 8.071 | 8.4 | 1.217 | 25.27 | 6.26 | 5.21 | 3.25 | 4.85 |
| 70 | $4^{\text {h }} 30^{\text {m }} 10$ ¢ 16 | $35^{\circ} 11^{\prime 2} 28.2$ | 0.036 | 24.189 | 7.8 | 1.252 | 11.87 | 6.03 | 5.01 | 36.99 | 2.37 |
| 71 | $4^{\mathrm{h}} 30^{\mathrm{m}} 10.58$ | $35^{\circ} 13^{\prime} 36.13$ | 0.036 | 24.038 | 6.8 | 1.801 | 24.53 | 9.65 | 8.02 | 59.23 | 3.06 |
| 72 | $4^{\text {h }} 30^{\text {m }} 10595$ | $35^{\circ} 28^{\prime} 46.12$ | 0.044 | 6.260 | 8.3 | 1.907 | 33.18 | 2.64 | 2.20 | 8.89 | 15.10 |
| 73 | $4^{\mathrm{h}} 30^{\mathrm{m}} 10.595$ | $35^{\circ} 28^{\prime} 54.17$ | 0.049 | 6.364 | 8.4 | 1.871 | 35.73 | 2.94 | 2.44 | 7.16 | 14.62 |
| 74 | $4^{\mathrm{h}} 30^{\mathrm{m}} 11{ }^{\text {s }} 21$ | 35 ${ }^{\circ} 3^{\prime} 49.18$ | 0.079 | 8.293 | 9.0 | 1.346 | 30.07 | 5.64 | 4.69 | 3.28 | 6.41 |
| 75 | $4^{\mathrm{h}} 30^{\mathrm{m}} 12 .{ }^{\text {s }} 19$ | $35^{\circ} 10^{\prime} 03.16$ | 0.060 | 19.816 | 7.6 | 1.520 | 28.86 | 13.42 | 11.16 | 17.80 | 2.59 |
| 76 | $4^{\mathrm{h}} 30^{\mathrm{m}} 12{ }^{\text {s }}$. 46 | $35^{\circ} 14^{\prime} 59.18$ | 0.035 | 14.930 | 7.9 | 1.405 | 14.25 | 26.03 | 2.16 | 17.39 | 6.59 |

Table 2. (Continued)

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Table 2. (Continued)

| ID | RA | Dec | $\begin{aligned} & R_{\text {core }} \\ & {[\mathrm{pcc}} \end{aligned}$ | $\begin{gathered} T_{\text {peak }} \\ {[\mathrm{K}]} \end{gathered}$ | Background level [K] | $\begin{gathered} \Delta \mathrm{V} \\ {\left[\mathrm{~km} \mathrm{~s}^{-1}\right]} \end{gathered}$ | $\begin{gathered} M_{\mathrm{vir}} \\ {\left[M_{\odot}\right]} \end{gathered}$ | $\begin{gathered} \stackrel{\Sigma N\left(\mathrm{H}_{2}\right)}{\left[\times 10^{23} \mathrm{~cm}^{-2}\right]} . \end{gathered}$ | $\begin{aligned} & M_{\mathrm{LTE}} \\ & {\left[M_{\odot}\right]} \end{aligned}$ | $\begin{gathered} n \\ {\left[\times 10^{4} \mathrm{~cm}^{-3}\right]} \end{gathered}$ | $\begin{gathered} M_{\mathrm{vir}} / M_{\mathrm{LTE}} \\ \left(=\alpha_{\mathrm{vir}}\right) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 189 | $4^{\mathrm{h}} 30^{\text {m }} 39 ¢ 30$ | $35^{\circ} 14^{\prime} 14.18$ | 0.038 | 12.043 | 16.3 | 1.551 | 19.02 | 2.93 | 2.44 | 15.31 | 7.80 |
| 190 | $4^{\text {h }} 30^{\text {m }} 39953$ | $35^{\circ} 18^{\prime} 122^{\prime \prime} 4$ | 0.046 | 8.542 | 23.8 | 0.909 | 8.01 | 2.10 | 1.74 | 6.17 | 4.60 |
| 191 | $4^{\mathrm{h}} 30^{\mathrm{m}} 3995$ | $35^{\circ} 22^{\prime} 16^{\prime \prime} 4$ | 0.044 | 12.458 | 8.2 | 1.194 | 13.00 | 3.23 | 2.68 | 10.86 | 4.84 |
| 192 | $4^{\mathrm{h}} 30^{\text {m }} 399.74$ | $35^{\circ} 19^{\prime} 30{ }^{\prime \prime} 6$ | 0.051 | 8.917 | 26.2 | 1.248 | 16.66 | 3.07 | 2.56 | 6.64 | 6.52 |
| 193 | $4^{\mathrm{h}} 30^{\mathrm{m}} 400^{5} 07$ | $35^{\circ} 06^{\prime} 53.4$ | 0.046 | 9.052 | 28.3 | 1.432 | 19.86 | 2.60 | 2.16 | 7.64 | 9.20 |
| 194 | $4^{\mathrm{h}} 30^{\mathrm{m}} 40$ 5 09 | $35^{\circ} 33^{\prime} 444^{\prime \prime} 4$ | 0.042 | 13.840 | 22.9 | 1.614 | 23.03 | 3.85 | 3.20 | 14.87 | 7.20 |
| 195 | $4^{\mathrm{h}} 30^{\mathrm{m}} 40{ }^{\text {¢ }} 39$ | $35^{\circ} 05^{\prime} 59.9$ | 0.059 | 11.181 | 7.5 | 1.364 | 22.84 | 4.57 | 3.80 | 6.38 | 6.01 |
| 196 | $4^{\mathrm{h}} 30^{\mathrm{m}} 40{ }^{5} 52$ | $35^{\circ} 29^{\prime} 21.77$ | 0.036 | 18.353 | 12.4 | 1.486 | 16.71 | 4.45 | 3.70 | 27.32 | 4.52 |
| 197 | $4^{\mathrm{h}} 30^{\mathrm{m}} 40{ }^{\text {¢ }} 59$ | $35^{\circ} 27^{\prime} 56.10$ | 0.065 | 19.268 | 9.4 | 1.337 | 24.49 | 13.89 | 11.55 | 14.49 | 2.12 |
| 198 | $4^{\mathrm{h}} 30^{\mathrm{m}} 40{ }^{5} 66$ | $35^{\circ} 21^{\prime} 37^{\prime \prime} 1$ | 0.035 | 13.307 | 11.1 | 1.248 | 11.24 | 2.33 | 1.94 | 15.57 | 5.80 |
| 199 | $4^{\mathrm{h}} 30^{\mathrm{m}} 40{ }^{5} 91$ | $35^{\circ} 35^{\prime} 411^{\prime \prime} 0$ | 0.042 | 9.836 | 10.3 | 1.615 | 23.04 | 2.86 | 2.38 | 11.06 | 9.68 |
| 200 | $4^{\mathrm{h}} 30^{\mathrm{m}} 41^{\text {s }} 30$ | $35^{\circ} 40^{\prime} 51.77$ | 0.045 | 7.568 | 8.2 | 1.267 | 15.09 | 1.94 | 1.62 | 6.11 | 9.33 |
| 201 | $4^{\mathrm{h}} 30^{\mathrm{m}} 411^{\text {s }} 38$ | $35^{\circ} 28^{\prime} 39.14$ | 0.042 | 19.346 | 6.3 | 1.369 | 16.56 | 5.55 | 4.61 | 21.45 | 3.59 |
| 202 | $4^{\mathrm{h}} 30^{\mathrm{m}} 41^{1}$ ¢ 40 | $35^{\circ} 11^{\prime} 33.17$ | 0.085 | 6.693 | 6.9 | 2.224 | 87.42 | 10.86 | 9.03 | 5.06 | 9.69 |
| 203 | $4^{\mathrm{h}} 30^{\mathrm{m}} 41^{\text {s }}$ : 50 | $35^{\circ} 31^{\prime \prime} 01{ }^{\prime \prime} 9$ | 0.035 | 13.705 | 7.0 | 1.548 | 17.28 | 3.03 | 2.52 | 20.22 | 6.87 |
| 204 | $4^{\mathrm{h}} 30^{\mathrm{m}} 411^{\text {s }} 59$ | $35^{\circ} 43^{\prime} 09.9$ | 0.101 | 8.797 | 6.4 | 1.476 | 45.81 | 12.98 | 10.79 | 3.61 | 4.25 |
| 205 | $4^{\mathrm{h}} 30^{\mathrm{m}} 41{ }^{\text {s }}$ : 79 | $35^{\circ} 08^{\prime} 31.17$ | 0.036 | 8.145 | 8.1 | 1.967 | 29.27 | 2.21 | 1.84 | 13.58 | 15.91 |
| 206 | $4^{\mathrm{h}} 30^{\mathrm{m}} 41^{\text {s }}$ s 80 | $35^{\circ} 22^{\prime} 444^{\prime \prime} 7$ | 0.048 | 11.270 | 6.7 | 1.420 | 20.07 | 4.07 | 3.38 | 10.53 | 5.94 |
| 207 | $4^{\mathrm{h}} 30^{\mathrm{m}} 42.574$ | $35^{\circ} 36^{\prime} 15^{\prime \prime} 5$ | 0.057 | 10.861 | 6.8 | 1.377 | 22.46 | 4.95 | 4.12 | 7.66 | 5.46 |
| 208 | $4^{\mathrm{h}} 30^{\mathrm{m}} 43 \mathrm{~s}^{5} 21$ | $35^{\circ} 32^{\prime} 300^{\prime \prime}$ | 0.069 | 14.908 | 7.2 | 1.464 | 30.91 | 12.01 | 9.98 | 10.47 | 3.10 |
| 209 | $4^{\mathrm{h}} 30^{\mathrm{m}} 433^{5} 26$ | $35^{\circ} 21^{\prime} 32.17$ | 0.051 | 11.841 | 9.3 | 1.666 | 29.70 | 5.48 | 4.56 | 11.84 | 6.52 |
| 210 | $4^{\mathrm{h}} 30^{\mathrm{m}} 43 \mathrm{~s}^{5} 31$ | $35^{\circ} 44^{\prime} 16.15$ | 0.035 | 7.667 | 8.7 | 1.751 | 22.13 | 1.35 | 1.13 | 9.04 | 19.67 |
| 211 | $4^{\mathrm{h}} 30^{\mathrm{m}} 43546$ | $35^{\circ} 25^{\prime} 011^{\prime \prime} 0$ | 0.039 | 8.995 | 7.0 | 1.678 | 23.18 | 2.01 | 1.67 | 9.70 | 13.88 |
| 212 | $4^{\text {h }} 30^{\text {m }} 43$ S 59 | $35^{\circ} 33^{\prime} 55^{\prime \prime} 3$ | 0.038 | 12.503 | 6.9 | 1.513 | 18.09 | 2.27 | 1.89 | 11.87 | 9.57 |
| 213 | $4^{\text {h }} 30^{\text {m }} 433^{5} 66$ | $35^{\circ} 38^{\prime} 21.4$ | 0.051 | 9.039 | 7.4 | 1.511 | 24.42 | 2.78 | 2.31 | 5.99 | 10.59 |
| 214 | $4^{\mathrm{h}} 30^{\mathrm{m}} 43^{5} 74$ | $35^{\circ} 07^{\prime} 00^{\prime \prime}{ }^{\prime}$ | 0.041 | 8.773 | 14.2 | 1.427 | 17.39 | 2.44 | 2.03 | 10.14 | 8.58 |
| 215 | $4^{\text {h }} 30^{\text {m }} 433^{5} 79$ | $35^{\circ} 24^{\prime \prime} 17^{\prime \prime} 1$ | 0.062 | 8.801 | 13.5 | 1.532 | 30.29 | 4.76 | 3.95 | 5.72 | 7.66 |
| 216 | $4^{\mathrm{h}} 30^{\mathrm{m}} 43^{5} 99$ | $35^{\circ} 31^{\prime} 03.0$ | 0.033 | 15.312 | 7.7 | 1.447 | 14.33 | 2.19 | 1.82 | 17.43 | 7.88 |
| 217 | $4^{\mathrm{h}} 30^{\mathrm{m}} 44^{5} 502$ | $35^{\circ} 29^{\prime} 24.14$ | 0.070 | 18.105 | 10.4 | 1.474 | 31.74 | 13.79 | 11.46 | 11.51 | 2.77 |
| 218 | $4^{\mathrm{h}} 30^{\mathrm{m}} 44^{\text {a }} 10$ | $35^{\circ} 07^{\prime} 54.18$ | 0.041 | 8.785 | 9.8 | 2.038 | 35.45 | 2.48 | 2.06 | 10.32 | 17.17 |
| 219 | $4^{\mathrm{h}} 30^{\mathrm{m}} 44^{5}$ S 23 | $35^{\circ} 35^{\prime} 444^{\prime \prime} 0$ | 0.041 | 9.447 | 13.6 | 1.555 | 20.65 | 2.10 | 1.75 | 8.73 | 11.82 |
| 220 | $4^{\text {h }} 30^{\mathrm{m}} 44^{\text {s }}$. 67 | $35^{\circ} 19^{\prime} 32.19$ | 0.046 | 8.683 | 15.0 | 1.301 | 16.37 | 2.59 | 2.15 | 7.61 | 7.62 |
| 221 | $4^{\mathrm{h}} 30^{\mathrm{m}} 45^{5} 506$ | $35^{\circ} 14^{\prime} 499^{\prime \prime} 4$ | 0.049 | 5.808 | 17.7 | 1.955 | 39.00 | 2.84 | 2.36 | 6.92 | 16.51 |
| 222 | $4^{\mathrm{h}} 30^{\mathrm{m}} 455.20$ | $35^{\circ} 26^{\prime} 21.16$ | 0.029 | 18.356 | 19.8 | 1.060 | 6.78 | 1.04 | 0.86 | 12.17 | 7.87 |
| 223 | $4^{\text {h }} 30^{\text {m }} 45$ ! 36 | $35^{\circ} 34^{\prime} 31.15$ | 0.036 | 10.202 | 19.8 | 1.581 | 18.92 | 1.81 | 1.51 | 11.14 | 12.54 |
| 224 | $4^{\text {h }} 30^{\text {m }} 45540$ | $35^{\circ} 3^{\prime \prime} 033^{\prime \prime} 3$ | 0.067 | 9.925 | 11.3 | 1.523 | 32.63 | 7.42 | 6.17 | 7.07 | 5.29 |
| 225 | $4^{\text {h }} 30^{\text {m }} 45550$ | $35^{\circ} 31^{\prime} 23.11$ | 0.035 | 15.365 | 9.5 | 1.388 | 13.90 | 2.66 | 2.21 | 17.74 | 6.30 |
| 226 | $4^{\text {h }} 30^{\text {m }} 45$ ¢ 56 | $35^{\circ} 28^{\prime} 37{ }^{\prime \prime} 2$ | 0.039 | 13.084 | 13.0 | 1.598 | 21.00 | 3.24 | 2.69 | 15.63 | 7.80 |

Table 2. (Continued)

| ID | RA | Dec | $\begin{aligned} & R_{\text {core }} \\ & {[\mathrm{pc}]} \end{aligned}$ | $\begin{gathered} T_{\text {peak }} \\ {[\mathrm{K}]} \end{gathered}$ | Background level [K] | $\begin{gathered} \Delta \mathrm{V} \\ {\left[\mathrm{~km} \mathrm{~s}^{-1}\right]} \end{gathered}$ | $\begin{gathered} M_{\mathrm{vir}} \\ {\left[M_{\odot}\right]} \end{gathered}$ | $\begin{gathered} \Sigma N\left(\mathrm{H}_{2}\right) \\ {\left[\times 10^{23} \mathrm{~cm}^{-2}\right]} \end{gathered}$ | $\begin{aligned} & M_{\mathrm{LTE}} \\ & {\left[M_{\odot}\right]} \end{aligned}$ | $\begin{gathered} n \\ {\left[\times 10^{4} \mathrm{~cm}^{-3}\right]} \end{gathered}$ | $\begin{gathered} M_{\mathrm{vir}} / M_{\mathrm{LTE}} \\ \left(=\alpha_{\mathrm{vir}}\right) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 227 | $4^{\text {h }} 30^{\text {m }} 45.60$ | $35^{\circ} 22^{\prime} 48 . \prime 3$ | 0.053 | 10.510 | 11.3 | 1.535 | 26.32 | 4.44 | 3.69 | 8.54 | 7.13 |
| 228 | $4^{\text {h }} 30^{\text {m }} 45^{\text {s }}$. 76 | $35^{\circ} 33^{\prime} 14 . \prime 2$ | 0.046 | 14.315 | 8.6 | 1.437 | 20.00 | 4.31 | 3.58 | 12.67 | 5.59 |
| 229 | $4^{\text {h }} 30^{\text {m }} 46.5 .11$ | $35^{\circ} 13^{\prime} 42.15$ | 0.044 | 8.768 | 9.2 | 1.368 | 17.07 | 3.19 | 2.65 | 10.73 | 6.43 |
| 230 | $4^{\text {h }} 30^{\text {m }} 46.5$ | $35^{\circ} 26^{\prime} 59 . \prime 3$ | 0.073 | 18.542 | 15.4 | 1.124 | 19.34 | 12.80 | 10.64 | 9.43 | 1.82 |
| 231 | $4^{\text {h }} 30^{\text {m }} 46.53$ | $35^{\circ} 20^{\prime} 36.16$ | 0.041 | 7.521 | 8.8 | 2.149 | 39.42 | 2.53 | 2.11 | 10.53 | 18.72 |
| 232 | $4^{\text {h }} 30^{\text {m }} 46.51$ | $35^{\circ} 36^{\prime} 20.10$ | 0.042 | 9.091 | 9.2 | 1.570 | 21.79 | 2.75 | 2.29 | 10.64 | 9.52 |
| 233 | $4^{\text {h }} 30^{\text {m }} 46.583$ | $35^{\circ} 06^{\prime} 08^{\prime \prime} 7$ | 0.060 | 9.789 | 14.5 | 1.272 | 20.21 | 4.41 | 3.67 | 5.85 | 5.51 |
| 234 | $4^{\mathrm{h}} 30^{\mathrm{m}} 46^{\text {s }} 98$ | $35^{\circ} 43^{\prime} 53 . \prime 3$ | 0.100 | 7.600 | 11.4 | 1.641 | 56.31 | 12.23 | 10.17 | 3.50 | 5.54 |
| 235 | $4^{\mathrm{h}} 30^{\mathrm{m}} 47^{\text {s }} .14$ | $35^{\circ} 20^{\prime} 31.12$ | 0.065 | 7.028 | 10.3 | 2.184 | 64.36 | 5.90 | 4.91 | 6.16 | 13.11 |
| 236 | $4^{\text {h }} 30^{\text {m }} 47^{\text {s }}$. 43 | $35^{\circ} 34^{\prime} 54 . \prime 4$ | 0.071 | 8.921 | 9.8 | 1.761 | 45.84 | 7.20 | 5.99 | 5.76 | 7.66 |
| 237 | $4^{\text {h }} 30^{\text {m }} 47{ }^{\text {s }}$. 45 | $35^{\circ} 10^{\prime} 43 . \prime 7$ | 0.083 | 8.332 | 6.8 | 1.712 | 50.95 | 12.44 | 10.34 | 6.23 | 4.93 |
| 238 | $4^{\text {h }} 30^{\text {m }} 47.57$ | $35^{\circ} 08^{\prime} 36.18$ | 0.029 | 9.236 | 8.5 | 2.100 | 26.63 | 1.46 | 1.21 | 17.11 | 21.99 |
| 239 | $4^{\text {h }} 30^{\text {m }} 48 .{ }^{\text {s }} 12$ | $35^{\circ} 32^{\prime} 39^{\prime \prime} 8$ | 0.036 | 14.603 | 10.0 | 1.594 | 19.23 | 3.09 | 2.56 | 18.94 | 7.50 |
| 240 | $4^{\text {h }} 30^{\text {m }} 48 .{ }^{\text {s }} 15$ | $35^{\circ} 08^{\prime} 04.10$ | 0.051 | 9.326 | 13.2 | 1.844 | 36.38 | 4.99 | 4.15 | 10.78 | 8.77 |
| 241 | $4^{\mathrm{h}} 30^{\mathrm{m}} 48^{\text {s }} 58$ | $35^{\circ} 24^{\prime} 16^{\prime \prime} 8$ | 0.052 | 7.488 | 9.1 | 1.755 | 33.70 | 3.02 | 2.51 | 6.15 | 13.43 |
| 242 | $4^{\text {h }} 30^{\text {m }} 48 .{ }^{\text {s }} 63$ | $35^{\circ} 38^{\prime} 52.12$ | 0.065 | 9.583 | 12.6 | 1.263 | 21.83 | 5.21 | 4.33 | 5.44 | 5.04 |
| 243 | $4^{\mathrm{h}} 30^{\mathrm{m}} 48^{\text {s }}$. 84 | $35^{\circ} 37^{\prime} 11.13$ | 0.044 | 11.175 | 8.4 | 1.455 | 19.33 | 3.13 | 2.60 | 10.53 | 7.43 |
| 244 | $4^{\mathrm{h}} 30^{\mathrm{m}} 49.36$ | $35^{\circ} 22^{\prime} 28.12$ | 0.058 | 10.779 | 8.6 | 1.650 | 32.89 | 7.00 | 5.82 | 10.28 | 5.65 |
| 245 | $4^{\mathrm{h}} 30^{\mathrm{m}} 49.62$ | $35^{\circ} 41^{\prime} 56{ }^{\prime \prime} 8$ | 0.090 | 7.073 | 8.7 | 1.485 | 41.52 | 6.86 | 5.70 | 2.69 | 7.29 |
| 246 | $4^{\text {h }} 30^{\text {m }} 50.33$ | $35^{\circ} 37^{\prime} 47.10$ | 0.042 | 11.625 | 8.9 | 1.414 | 17.67 | 2.92 | 2.43 | 11.29 | 7.28 |
| 247 | $4^{\text {h }} 30^{\text {m }} 50.65$ | $35^{\circ} 07^{\prime} 14.12$ | 0.060 | 8.358 | 20.0 | 1.494 | 27.91 | 5.51 | 4.58 | 7.30 | 6.10 |
| 248 | $4^{\text {h }} 30^{\text {m }} 50.83$ | $35^{\circ} 20^{\prime} 10 . \prime 2$ | 0.065 | 6.586 | 12.7 | 1.774 | 42.50 | 4.34 | 3.61 | 4.53 | 11.77 |
| 249 | $4^{\text {h }} 30^{\text {m }} 50.54$ | $35^{\circ} 08^{\prime} 29.15$ | 0.081 | 8.321 | 6.9 | 1.697 | 48.76 | 12.45 | 10.35 | 6.71 | 4.71 |
| 250 | $4^{\text {h }} 30^{\text {m }} 50.54$ | $35^{\circ} 20^{\prime} 32.18$ | 0.051 | 7.282 | 19.8 | 1.869 | 37.37 | 2.90 | 2.41 | 6.26 | 15.51 |
| 251 | $4^{\mathrm{h}} 30^{\mathrm{m}} 51.53$ | $35^{\circ} 24^{\prime} 16^{\prime \prime} .9$ | 0.039 | 7.124 | 15.6 | 1.820 | 27.24 | 1.64 | 1.36 | 7.90 | 20.04 |
| 252 | $4^{\text {h }} 30^{\text {m }} 51 . \mathrm{s} 67$ | $35^{\circ} 34^{\prime} 23.10$ | 0.039 | 10.308 | 10.7 | 1.434 | 16.93 | 1.96 | 1.63 | 9.47 | 10.39 |
| 253 | $4^{\mathrm{h}} 30^{\mathrm{m}} 51.58$ | $35^{\circ} 10^{\prime} 48.19$ | 0.033 | 9.772 | 8.4 | 1.727 | 20.41 | 2.35 | 1.95 | 18.71 | 10.46 |
| 254 | $4^{\mathrm{h}} 30^{\mathrm{m}} 51.93$ | $35^{\circ} 26^{\prime} 43.17$ | 0.035 | 14.081 | 11.6 | 1.272 | 11.67 | 2.58 | 2.15 | 17.25 | 5.44 |
| 255 | $4^{\text {h }} 30^{\text {m }} 52.521$ | $35^{\circ} 39^{\prime} 45.16$ | 0.095 | 10.201 | 10.0 | 1.183 | 27.66 | 11.12 | 9.24 | 3.72 | 2.99 |
| 256 | $4^{\text {h }} 30^{\mathrm{m}} 52 . \mathrm{s} 21$ | $35^{\circ} 26^{\prime} 09^{\prime \prime} 0$ | 0.036 | 12.004 | 10.9 | 1.345 | 13.70 | 1.55 | 1.29 | 9.53 | 10.62 |
| 257 | $4^{\text {h }} 30^{\text {m }} 52.536$ | $35^{\circ} 38^{\prime} 39.15$ | 0.036 | 10.691 | 11.1 | 1.310 | 12.99 | 1.78 | 1.48 | 10.91 | 8.79 |
| 258 | $4^{\text {h }} 30^{\text {m }} 52.539$ | $35^{\circ} 38^{\prime} 04{ }^{\prime \prime} 3$ | 0.048 | 11.447 | 8.3 | 1.365 | 18.54 | 3.54 | 2.94 | 9.17 | 6.30 |
| 259 | $4^{\text {h }} 30^{\text {m }} 52 .{ }^{\text {s }} 60$ | $35^{\circ} 32^{\prime} 53 . \prime 2$ | 0.042 | 12.147 | 8.0 | 1.375 | 16.70 | 2.60 | 2.16 | 10.06 | 7.72 |
| 260 | $4^{\mathrm{h}} 30^{\mathrm{m}} 53.5$ | $35^{\circ} 29^{\prime} 22.15$ | 0.052 | 14.017 | 16.6 | 1.100 | 13.24 | 5.03 | 4.18 | 10.24 | 3.17 |
| 261 | $4^{\text {h }} 30^{\text {m }} 53.37$ | $35^{\circ} 27^{\prime} 57.19$ | 0.033 | 14.905 | 14.9 | 1.108 | 8.41 | 1.67 | 1.39 | 13.33 | 6.05 |
| 262 | $4^{\mathrm{h}} 30^{\mathrm{m}} 53.58$ | $35^{\circ} 22^{\prime} 56.10$ | 0.036 | 9.497 | 6.9 | 1.876 | 26.63 | 2.14 | 1.78 | 13.13 | 14.98 |
| 263 | $4^{\mathrm{h}} 30^{\mathrm{m}} 53 .{ }^{\mathrm{s}} 88$ | $35^{\circ} 21^{\prime} 08 . \prime 6$ | 0.042 | 8.843 | 10.5 | 1.921 | 32.60 | 2.68 | 2.23 | 10.37 | 14.62 |
| 264 | $4^{\text {h }} 30^{\text {m }} 54.5$ | $35^{\circ} 28^{\prime} 33 . \prime 9$ | 0.044 | 13.970 | 16.1 | 1.137 | 11.80 | 3.47 | 2.89 | 11.68 | 4.09 |

Table 2．（Continued）

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Table 2．（Continued）

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Fig. 8. Integrated intensity map of the ${ }^{13} \mathrm{CO}(J=1-0)$ emission line towards the L 1482 molecular cloud. The circles indicate the locations and diameters of the identified cores.

Table 3. Summary of the physical properties of the core.

|  | Minimum | Mean | Maximum |
| :--- | :---: | :---: | :---: |
| $R_{\text {core }}[\mathrm{pc}]$ | 0.029 | 0.054 | 0.180 |
| $T_{\text {peak }}[\mathrm{K}]$ | 5.808 | 12.230 | 30.578 |
| Background level $[\mathrm{K}]$ | 6.2 | 11.3 | 28.3 |
| $\Delta \mathrm{~V}\left[\mathrm{~km} \mathrm{~s}^{-1}\right]$ | 0.606 | 1.426 | 2.674 |
| $M_{\text {vir }}\left[M_{\odot}\right]$ | 3.19 | 23.55 | 91.65 |
| $\Sigma N\left(\mathrm{H}_{2}\right)\left[\times 10^{23} \mathrm{~cm}^{-2}\right]$ | 1.11 | 6.05 | 42.69 |
| $M_{\mathrm{LTE}}\left[M_{\odot}\right]$ | 0.76 | 4.58 | 35.49 |
| $n\left[\times 10^{4} \mathrm{~cm}^{3}\right]$ | 0.35 | 12.96 | 72.18 |
| $M_{\text {vir }} / M_{\mathrm{LTE}}\left(=\alpha_{\text {vir }}\right)$ | 1.03 | 6.94 | 33.23 |

with a mean value of $\sim 0.054 \mathrm{pc}$. The core mass ranges from 0.76 to $35.49 M_{\odot}$. We show the summary of the physical parameters of the core in table 3.

Finally, we also derived the total molecular mass of the entire region mapped in ${ }^{13} \mathrm{CO}$ (figure 6) to be $\sim 7600 M_{\odot}$ using the derived $T_{\text {ex }}$ and equations (2) and (3).

## 4 Discussion

### 4.1 The relation between $L 1482$ filament and LkH $\alpha 101$

The measured parallax of the $\mathrm{H}_{2} \mathrm{O}$ maser in L 1482 filament 2 is consistent with that of the $\mathrm{LkH} \alpha 101$ cluster measured by VLBA and Gaia DR2. The difference is only
$0.01 \pm 0.10$ mas. The parallax and the proper motion of VLA J043001.15+337724.6 in the $\mathrm{LkH} \alpha 101$ cluster is measured to be $\pi=1.87 \pm 0.10$ mas and $\left(\mu_{\alpha} \cos \delta, \mu_{\delta}\right)=$ $(1.86 \pm 0.04,-5.70 \pm 0.05) \mathrm{mas}_{\mathrm{yr}}{ }^{-1}$ with VLBA radio continuum observations (Dzib et al. 2018). The parallax and the proper motion of $\mathrm{LkH} \alpha 101$ cluster using Gaia DR2 data (Gaia Collaboration 2018) are $\pi=1.876 \pm 0.051$ mas and $\left(\mu_{\alpha} \cos \delta, \mu_{\delta}\right)=(2.44 \pm 0.55,-5.27 \pm 0.54) \mathrm{mas} \mathrm{yr}^{-1}$, respectively. These are the averages of the parallaxes and the proper motions of 35 sources which are selected from the following three steps:
(i) We selected 178 sources located within $5^{\prime}$ from the center of the $\mathrm{LkH} \alpha 101$ cluster $(\alpha, \delta)_{\mathrm{J} 2000.0}=$ $\left(04^{\mathrm{h}} 30^{\mathrm{m}} 11^{\mathrm{s}} 04,+35^{\circ} 16^{\prime} 44^{\prime \prime} 0\right)$.
(ii) We selected 38 sources whose parallax errors are less than $10 \%$.
(iii) We excluded three foreground sources with parallaxes $3.8,5.0$, and 13.5 mas.

Finally, 35 sources were selected. Figures 9 and 10 show distributions of the parallax and proper motion of the 35 sources. There is a difference of approximately 3 mas $\mathrm{yr}^{-1}$ $\left(8 \mathrm{~km} \mathrm{~s}^{-1}\right.$ at 532 pc$)$ in the proper motions between the $\mathrm{H}_{2} \mathrm{O}$ maser spot and the $\mathrm{LkH} \alpha 101$ cluster. While the $\mathrm{H}_{2} \mathrm{O}$ maser spot is detected at $V_{\text {LSR }}=7.38 \mathrm{~km} \mathrm{~s}^{-1}$, the ambient molecular gas in L 1482 filament 2 is detected at $V_{\text {LSR }} \simeq-1 \mathrm{~km} \mathrm{~s}^{-1}$


Fig. 9. Distributions of the parallax and proper motion of 35 sources measured by Gaia DR2 (Gaia Collaboration 2018).


Fig. 10. (Left) Red, blue, and white circles mark the sources observed with VERA (this work), VLBA (Dzib et al. 2018), and Gaia DR2 (Gaia Collaboration 2018), respectively. The background is the Digitized Sky Survey 2 (DSS2) optical image. The cyan contours show the ${ }^{13} \mathrm{CO}(\mathrm{J}=1-0)$ integrated intensity (Li et al. 2014). The origin ( 0,0 ) corresponds to $(\alpha, \delta)_{\mathrm{J} 2000.0}=04^{\mathrm{h}} 30^{\mathrm{m}} 11^{\mathrm{s}} .04,+35^{\circ} 16^{\prime} 44 .{ }^{\prime \prime} 0$. (Right) Proper motions of the objects represented
 at the assumed distance 532 pc .
in the ${ }^{13} \mathrm{CO}(J=1-0)$ line (Li et al. 2014). The difference in velocity is approximately $8 \mathrm{~km} \mathrm{~s}^{-1}$, which may originate from the internal motion of the $\mathrm{H}_{2} \mathrm{O}$ maser spot with respect to the associated YSO and/or the peculiar motion of the associated YSO itself. The standard deviations of proper motions of the 35 sources in the $\mathrm{LkH} \alpha 101$ cluster are calculated to be $3.27 \mathrm{mas} \mathrm{yr}^{-1}$ and 3.20 mas $\mathrm{yr}^{-1}$ in RA and Dec, respectively. This indicates a velocity dispersion of approximately $8.2 \mathrm{~km} \mathrm{~s}^{-1}$. However, two outlier sources which have large peculiar motions of $36 \mathrm{~km} \mathrm{~s}^{-1}$ and $50 \mathrm{~km} \mathrm{~s}^{-1}$ are included in the 35 sources, and these would lead to the overestimation of the velocity dispersion. We excluded two
outlier sources and calculated the standard deviations to
 tively, indicating the velocity dispersion of approximately $3.3 \mathrm{~km} \mathrm{~s}^{-1}$. This is close to the average velocity dispersion of $4.5 \mathrm{~km} \mathrm{~s}^{-1}$ calculated for 28 OB-associations (Melnik \& Dambis 2020).

### 4.2 The arc-shaped warm region

It is noted that the warm arc-shaped region surrounds $\mathrm{LkH} \alpha 101$ in the center. The L 1482 molecular filament and the $\mathrm{LkH} \alpha 101$ star cluster are located at the same distance


Fig. 11. Integrated intensity map of ${ }^{13} \mathrm{CO}(J=1-0)$ towards the L 1482 molecular cloud overlaid with the VLA 1.4-GHz continuum emission (Condon et al. 1998). The contours start from $0.33 \mathrm{mJy} \mathrm{beam}^{-1}$ with an increment of $0.4 \mathrm{mJy} \mathrm{beam}^{-1}$.
within $3 \pm 30 \mathrm{pc}$. This suggests that the warm arc-shaped region is located at the periphery of the $\mathrm{H}_{\text {II }}$ region ionized by $\mathrm{LkH} \alpha 101$. In order to study the relationship between the warm molecular region and the $\mathrm{H}_{\text {II }}$ region, we superposed the VLA 1.4 GHz map (taken from Condon et al. 1998) on the ${ }^{13} \mathrm{CO}$ integrated intensity map (see figure 11), which clearly shows that the arc-shaped region is heated by the FUV $(6 \mathrm{eV}<\mathrm{h} v<13.6 \mathrm{eV})$ radiation from $\mathrm{LkH} \alpha 101$. The arc-shaped region appears similar to the bright bar in Orion, and can be an ideal region to study photondominated region and star formation induced by the interaction with the $\mathrm{H}_{\text {II }}$ regions created by massive stars.

The thickness of the warm region is about $3^{\prime}$, corresponding to $\sim 0.46 \mathrm{pc}$ at the distance of 532 pc . Early theoretical model of PDRs (Tielens \& Hollenbach 1985a, 1985b; Köster et al. 1994) estimated that the thickness of the region heated by the FUV corresponds to $A_{\mathrm{V}}=10 \mathrm{mag}$ in visual extinction, being only 0.03 pc assuming hydrogen gas densities $n_{\mathrm{H}_{2}}$ of $10^{5} \mathrm{~cm}^{-3}$. As the observed ${ }^{12} \mathrm{CO}$ emission line in PDR regions requires very high hydrogen gas densities, $n_{\mathrm{H}_{2}}=10^{6}-10^{7} \mathrm{~cm}^{-3}$ (Burton et al. 1990). The emission from the PDR regions facing $\mathrm{LkH} \alpha 101$ may be produced in dense cores embedded in lower-density gas. In the warm region, we often observe a clear dip in the ${ }^{12} \mathrm{CO}$ spectra. As shown in figure 7 , the radial velocity of the dip nicely matches with the peak velocity of the optically thinner ${ }^{13} \mathrm{CO}$ emission, indicating that the dip is due to the absorption by a large amount of colder gas


Fig. 12. $M_{\text {vir }}$ vs. $M_{\text {LTE }}$ relation of the 337 cores found in this study. Blue and red dots indicate cores with $\alpha_{\text {vir }} \leq 1.5$ and $>1.5$, respectively.
lying in the foreground of the warm region. Because the dip is widely seen around the arc-shaped region, the hightemperature gas should be distributed over a considerable extent.

### 4.3 Star formation activities in the L 1482 filament

As noted in subsection 3.3, we identified 337 cores using Dendrogram analysis. Figure 12 shows the relation between $M_{\text {vir }}$ and $M_{\text {LTE }}$ of the 337 cores. As seen in the figure, the virial ratio $\alpha_{\text {vir }}\left(=M_{\text {vir }} / M_{\mathrm{LTE}}\right)$ is greater than 1 . In general,


Fig. 13. $M_{\text {LTE }}-\alpha$ relation of the cores in logarithmic scale. The solid line shows the slope of -0.57 from our data and the dashed line shows the slope of -0.67 of Bertoldi \& McKee (1992).
in case there is no external pressure ( $P_{\text {ext }}=0$ ), cores with $\alpha_{\text {vir }}<1$ are gravitationally unstable and are to collapse, while cores with $\alpha_{\text {vir }}>1$ are gravitationally unbound and are to disperse. Shimoikura et al. (2019) who studied massive cores in M17 suggested that the cores with $\alpha_{\text {vir }}>1$ can also be gravitationally bound, if they are surrounded by high external pressure.

Since the region observed in this study is adjacent to $\mathrm{LkH} \alpha 101$ and a part of the region should be located in the high external pressure of the $\mathrm{H}_{\text {II }}$ region, some of the cores we detected may be gravitationally bound. Virial ratios of spherical cores with uniform density can reach $\alpha_{\text {vir }}=$ 2 at most, if they are in the gravitational equilibrium in high external pressure. However, the cores we detected are not perfectly spherical or uniform, and thus we tentatively assume in this paper that cores with $\alpha_{\text {vir }} \leq 1.5$ are gravitationally bound and that cores with $\alpha_{\text {vir }}>1.5$ are gravitationally unbound. As seen in figure 8 , there are seven cores with $\alpha_{\text {vir }} \leq 1.5$ which are indicated by the blue circles in the figure, and the other cores with $\alpha_{\text {vir }}>1.5$ are indicated by the red circles. Three of the seven cores are located in the photo-dissociated region facing $\mathrm{LkH} \alpha$ 101, and the other four cores are found in the northwest of $\mathrm{LkH} \alpha 101$. As is clear from figure 8 , the ${ }^{13} \mathrm{CO}$ emission is relatively weak in the northwest region between $\mathrm{LkH} \alpha 101$ and the four cores in the northwest, and the molecular density and therefore the external pressure is likely to be low there. It is noted that the size of the other three cores with $\alpha_{\text {vir }} \leq 1.5$ adjacent to $\mathrm{LkH} \alpha 101$ is very small compared with the four cores in the northwest, suggesting that the cores close to $\mathrm{LkH} \alpha 101$ are compressed by the high pressure in the $\mathrm{H}_{\text {II }}$ region. We imagine that these cores may eventually form stars when they become unstable due to increase of the external pressure, e.g., the impact of
ionization front of the $\mathrm{H}_{\text {II }}$ region or strong stellar wind from $\mathrm{LkH} \alpha 101$.

Filaments are noted as the earliest stages of star formation. Theories and models for instabilities and fragmentation of the filament have been extensively studied (e.g., Ostriker 1964; Inutsuka \& Miyama 1992). In contrast to the results of those studies, our results show that star formation is unlikely to occur in the massive filamentary structure L 1482 except for the southeast part facing $\mathrm{LkH} \alpha 101$. In the southeast part of the filament adjacent to $\mathrm{LkH} \alpha 101$ and the associated young cluster, the photo-dissociated warm region is formed, and the three cores with $\alpha_{\text {vir }} \leq 1.5$ are located therein. In this region, there are also protostars associated with $\mathrm{H}_{2} \mathrm{O}$ masers. These results suggest that radiation-driven star formation triggered by the expansion of the $\mathrm{H}_{\text {II }}$ region is proceeding in the southeastern part of the CMC.

Bertoldi and McKee (1992) proposed an $\alpha_{\text {vir }} \sim M^{-0.67}$ relation. To test the validity of the exponent of the virial parameter, we made the plot of the core masses against their corresponding $\alpha_{\text {vir }}$ (see, figure 13). Our regression fit produced an exponent of -0.57 . This is much closer to the -0.54 obtained for the Cepheus cloud (Bertoldi \& McKee 1992). This is an indication that L 1482 could be a pressureconfined clump with the possibility of becoming a site of active star formation in the future.

Recently, an ALMA view of the Orion bright bar directly exposed to the Trapezium stars reveals a fragmented ridge of high-density substructures, photo-ablative gas flows, and instabilities at the molecular cloud surface (Goicoechea et al. 2016). Our results show the existence of the hightemperature gas in the warm region where clumpy structures are detected. The possible excitation scenarios of the warm gas are as follows;



Fig. 14. $350 \mu \mathrm{~m}$ dust continuum emission of the OMC (left) and CMC (right).
(i) the ejection of energetic electrons from dust grains by FUV photons,
(ii) turbulent heating, where the energy stored in local velocity dispersion is thermalized in small-scale shocks (Guesten \& Fiebig 1988), and
(iii) photo-ablative high-temperature gas flows as observed in the Orion bar.

At the moment, it remains uncertain which mechanism is the most effective. Also, Goicoechea et al. (2016) suggested that the cloud edge has been compressed by a high-pressure wave that currently moves into the cloud. Our results are consistent with this interpretation. The reasons are as follows;
(1) the warm regions are formed adjacent to $B$ star,
(2) the column density in the warm region is 5-6 times higher than that of surrounding colder regions, and
(3) protostellar cores, an $\mathrm{H}_{2} \mathrm{O}$ maser source, and young stars are concentrated in this region.

This suggests that radiation-driven star formation is proceeding in this region.

Finally, we discuss the differences in star formation rates in the CMC and OMC. As noted by Lada et al. (2009), the CMC is very similar in morphology to the Orion A molecular cloud. However, the CMC shows an SFR that is an order of magnitude lower than that of the OMC. Figure 14 shows the $350 \mu \mathrm{~m}$ maps of the CMC (right-hand panel) and OMC (left-hand panel). The map of the CMC is enlarged to the scale of the distance of the OMC. As can be seen in figure 14 , there are a number of dense and prominent filaments in the OMC which are known sites of massive star formation, whereas filaments in the CMC are apparently much less dense. According to Krumholz and McKee (2008), a high column density greater than $\sim 1 \mathrm{~g} \mathrm{~cm}^{-2}$ [or
$\left.N\left(\mathrm{H}_{2}\right) \simeq 10^{23} \mathrm{~cm}^{-2}\right]$ is needed for the formation of massive stars. Formation of massive stars should be closely related to the high SFR through the compression of the natal clouds by the powerful stellar wind and the creation of the PDR regions. The critical density $\left(\sim 1 \mathrm{~g} \mathrm{~cm}^{-2}\right)$ corresponds to $N\left({ }^{13} \mathrm{CO}\right) \simeq 2 \times 10^{17} \mathrm{~cm}^{-2}$. Nakamura et al. (2019) recently analyzed the ${ }^{13} \mathrm{CO}$ data observed toward the OMC and found that roughly $\sim 10 \%$ of the total mass is confined in dense regions with $N\left({ }^{13} \mathrm{CO}\right)>2 \times 10^{17} \mathrm{~cm}^{-2}$ (see, their figure 27). The fraction of the high-density regions is much larger in the northern part of the OMC than in the diffuser, and more extended southern part which forms much less massive stars. A similar trend is also found in clouds associated with the M17 H II region and in rather quiescent infrared dark clouds in the vicinity known as M17-SWex (Nakamura et al. 2019; Nguyen-Luong et al. 2020). In the case of the CMC, we analyzed the frequency distribution of $N\left({ }^{13} \mathrm{CO}\right)$ using our data, and found that there are a very limited number of pixels (only 39 pixels) exceeding the critical value $N\left({ }^{13} \mathrm{CO}\right)=2 \times 10^{17} \mathrm{~cm}^{-2}$, and also found that such pixels are located only around the $\mathrm{LkH} \alpha 101$ cluster. This indicates that the mean density of the CMC is not high enough to produce massive stars, which can be one of the major reasons for the low SFR of the CMC compared with the OMC.

## 5 Conclusions

We have measured the trigonometric parallax of the $\mathrm{H}_{2} \mathrm{O}$ maser source associated with the L 1482 molecular filament to be $1.879 \pm 0.096$ mas, corresponding to the distance of $532 \pm 28 \mathrm{pc}$. This parallax is consistent with that of the nearby star cluster, $\mathrm{LkH} \alpha 101$. Our ${ }^{12} \mathrm{CO}$ observations revealed the clumpy arc-shaped, warm structure illuminated
by the $\mathrm{LkH} \alpha 101$ cluster. Based on the ${ }^{13} \mathrm{CO}$ observations, we performed core identification, and found 337 cores. Most of the cores are likely to be gravitationally bound and thus they are unlikely to form stars except for some cores located along the arc-shaped structure. The column density in the warm arc-shaped region is 5-6 times higher than that of surrounding colder molecular region, which supports the idea that the radiation-driven star formation triggered by the expansion of the $\mathrm{H}_{\text {II }}$ region is occurring around the $\mathrm{LkH} \alpha 101$ cluster.

We have also compared the CMC and the OMC, and found the presence of more filamentary structures in the OMC. The OMC has a significantly higher star formation rate compared to the CMC, and the difference may be caused by the fact that nearly all the cores in the L 1482 filament of the CMC are gravitationally stable or unbound and thus they are unable to form stars. We also discussed that the mean density of the CMC is not high enough to produce massive stars, which can be one of the main reasons for the low SFR of the CMC compared with the OMC.

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