

The nature of 500 micron risers I: SMA observations

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ABSTRACT

We present SMA observations at resolutions from 0.35 to 3 arcsec of a sample of 34 candidate high redshift dusty star forming galaxies (DSFGs). These sources were selected from the HerMES *Herschel* survey catalogues to have SEDs rising from 250 to 350 to 500 μm , a population termed 500-risers. We detect counterparts to 24 of these sources, with four having two counterparts. We conclude that the remaining ten sources that lack detected counterparts are likely to have three or more associated sources which blend together to produce the observed *Herschel* source. We examine the role of lensing, which is predicted to dominate the brightest ($F_{500} > 60$ mJy) half of our sample. We find that while lensing plays a role, at least 35 per cent of the bright sources are likely to be multiple sources rather than the result of lensing. At fainter fluxes we find a blending rate comparable to, or greater than, the predicted 40 per cent. We determine far-IR luminosities and star formation rates for the non-multiple sources in our sample and conclude that, in the absence of strong lensing, our 500-risers are very luminous systems with $L_{\text{FIR}} > 10^{13} L_{\odot}$ and star formation rates $> 1000 M_{\odot} \text{yr}^{-1}$.

Key words: galaxies: high-redshift – galaxies: starburst – infrared: galaxies – submillimetre: galaxies.

1 INTRODUCTION

The role and nature of dusty star forming galaxies (DSFGs) at high redshifts ($z > 4$) is currently unclear. Observations with *Herschel* have identified an unexpectedly large population of candidate high z DSFGs through colour selection in the three bands of its SPIRE instrument (Dowell et al. 2014; Asboth et al. 2016; Ivison et al. 2016). This selection is based on finding a rising spectral energy distribution (SED) from 250 to 500 μm , leading to these sources being termed ‘500-risers’. Since the far-IR SED of a typical galaxy peaks at around 100 μm in the rest frame, sources whose SED is still rising at 500 μm are likely to lie at $z > 4$. These 500-risers are found with densities on the sky of $3.3 \pm 0.8 \text{ deg}^{-2}$ down to 500 μm fluxes of 30 mJy (Dowell et al. 2014). Such sources, if at $z > 4$, would have luminosities $> 10^{13} L_{\odot}$ and star formation rates $> 1000 M_{\odot} \text{yr}^{-1}$. They would thus be extreme objects and as such

could represent the high luminosity end of a larger population of DSFGs at high redshift (see e.g. Greenslade et al. 2019). However, spectroscopic redshifts have been obtained for only a tiny fraction of this population (Capak et al. 2008; Coppin et al. 2009; Daddi et al. 2009; Riechers et al. 2010; Capak et al. 2011; Cox et al. 2011; Combes et al. 2012; Walter et al. 2012; Riechers et al. 2013; Dowell et al. 2014; Yun et al. 2015; Asboth et al. 2016; Ivison et al. 2016; Oteo et al. 2016; Fudamoto et al. 2017; Oteo et al. 2017a; Riechers et al. 2017; Strandet et al. 2017; Marrone et al. 2018; Zavala et al. 2018) and so its properties are only weakly constrained, and there is disagreement on whether this population contributes significantly to the cosmic SFR-density (SFRD) above $z > 4$ (Rowan-Robinson et al. 2016; Novak et al. 2017; Liu et al. 2018; Yan et al. 2019) or not (Koprowski et al. 2017; Michałowski et al. 2017). It has even been suggested that current data are unable to tell the difference between models where high z DSFGs dominate the SFRD or not (Casey et al. 2018).

The nature of the individual 500-risers that make up this population is also in dispute. Simulations by Béthermin et al. (2017) suggest that most 500 μm risers are in fact blends, where the

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brightest galaxy contributes on average 60 percent of the total *Herschel* 500 μm flux density, though this is also a function of 500 μm flux, and above $S_{500} = 60 \text{ mJy}$ they find almost all sources should be strongly lensed single galaxies (see fig. 6 of Béthermin et al. 2017). Gravitationally lensed sources will appear as pairs of sources separated by a few arcseconds, or as arc-like structures on similar scales (see Bussmann et al. (2013) for examples of lensed DSFGs observed by the SMA). Donevski et al. (2018) further suggest that confusion and instrumental noise also boost the number of 500-risers seen by *Herschel*. Such multiple sources could be chance alignments along the line of sight or might be physically associated systems within the same large-scale structure. They might also be the result of galaxy interactions and mergers that could also trigger starburst activity leading to their high far-IR luminosity. Observations with single dish submm telescopes e.g. Duivenvoorden et al. (2018) can go some way to addressing this since they can achieve somewhat higher angular resolutions than the 36 arcsec of the *Herschel* 500 μm beam. However, the idea that 500-risers are blends of multiple sources can best be tested with submm interferometry observations which can achieve angular resolutions of a few arcseconds or better, depending on configuration. Numerous interferometric studies of *Herschel* detected sources have been performed (Hodge et al. 2013; Karim et al. 2013; Swinbank et al. 2014; Simpson et al. 2015; Oteo et al. (2017b); Hill et al. 2018), but most of these have targeted sources selected at 850 μm and not as 500-risers.

The ideal test of the nature of the 500-riser population is to conduct submm interferometric observations with resolutions of a few arcseconds or better. This will be able to determine the multiplicity fraction and, at higher resolutions, examine the role of gravitational lensing. Oteo et al. (2017b) performed ALMA observations with a resolution of 0.12 arcsec at 850 μm on 500 μm risers selected from (Ivison et al. 2016), reaching average sensitivities of $0.1 \text{ mJy beam}^{-1}$, and with a primary beam full width at half-maximum (FWHM) of 17 arcsec. They find that 27 of their sample of 44 single dish sources (61 per cent) are single sources, whereas the remaining 16 sources (39 per cent) split into multiple galaxies. Additionally, they find several sources with multiplicities >3 , and one source where a single *Herschel* source resolves into five galaxies in the ALMA maps. They also find that 18 of their targets are gravitationally lensed to some extent. These results are interpreted as reducing the tension between models and the 500-riser counts from *Herschel*, with flux boosting thanks to confusion and multiplicity accounting for the discrepancy. However, this is not a complete sample, so statistical conclusions are somewhat uncertain, and the small primary beam means that some of the 500 μm flux measured by *Herschel* may be accounted for by further companion sources that lie outside the small ALMA beam.

It is thus clear that high resolution observations of 500 μm risers can provide valuable constraints on models, particularly the multiplicity of such sources and the relative contribution of multiple sources to the single dish observed flux density. There are strong predictions regarding the multiplicity of 500-risers, in particular that sources with $S_{500} > 60 \text{ mJy}$, should almost entirely be lensed single galaxies, something not borne out by Oteo et al. (2017b). This paper describes Sub-Millimeter Array (SMA; Ho, Moran & Lo 2004) observations of 34 500-risers to determine the multiplicity of these sources. A future paper will also use this data to obtain cross identifications at other wavelengths (Clements et al. in preparation). Objects are selected using a slightly less stringent method to those of Oteo et al. (2017b) (see below). Similar to Oteo et al. (2017b), our resulting set of observations does not represent a complete

sample, but we expect it to be reasonably representative. Results for a number of our sources have been presented elsewhere (Riechers et al. 2013; Dowell et al. 2014), but we here present the first uniform reduction of this data to examine the statistical properties of this population at resolutions of $\sim 2 \text{ arcsec}$.

The rest of this paper is structured as follows. In the next section we describe the sample of 500-risers observed. After that we describe the observations and data reduction, and present the results obtained. The subsequent section discusses the source properties derived from these observations and our other results. Finally, we draw our conclusions. Throughout this paper, we assume a concordance Λ -CDM cosmology, with $H_0 = 67.74 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_\Lambda = 0.69$, and $\Omega_m = 0.31$.

2 SAMPLE SELECTION

Our 500-riser targets are all selected from the HerMES survey (Oliver et al. 2012) and include sources chosen from both preliminary and finally published maps and catalogues. All our sources were selected as candidate high redshift DSFGs based on their *Herschel* fluxes and colours, but their selection is heterogenous in the sense that there is no single consistent flux or colour criterion used. This means that while it is a representative sample of sources, rigorous statistical conclusions may be difficult to draw since it is not a complete sample. There are 34 sources in total, with 500 μm flux densities ranging from 27 to 160 mJy, and an average of $67 \pm 29 \text{ mJy}$. For consistency, we extract the SPIRE fluxes of our sources from the latest versions of the HerMES maps¹ and catalogues. The methods used to detect sources and extract the final fluxes for these objects are described in Wang et al. (2013). The *Herschel* fluxes can be seen in Table 1, which also includes the number of SMA detections (see Section 3). Further details of our sources, as well as details of the SMA observations, can be found in Table 2. All sources were detected by *Herschel* at $>4\sigma$ with the majority at much higher levels of significance. Colour selection was on the basis of flux rising to 500 μm , i.e. $F_{250} < F_{350} < F_{500}$, the defining characteristic of a 500-riser, based on the data available at the time of selection and observation. However, as can be seen, refinements to the data reduction in later versions of the maps and catalogues mean that a few sources, such as Lock2, HeLMS35, and HFLS5, can no longer strictly be classed as 500 μm risers. However, we include them since their initial selection was as a 500 μm riser, and investigations by Donevski et al. (2018) and Béthermin et al. (2017) have shown that there can be a significant contamination of selected 500 μm risers by flux boosted sources with $S_{500} \lesssim S_{350}$.

3 OBSERVATIONS AND DATA ANALYSIS

3.1 Observations and data reduction

Our targets were observed in a number of programmes at the SMA over the period from 2010 May 19 to 2015 May 7. The observations were made at a variety of different tunings but generally at frequencies close to either 345 or 265 GHz. The most common configuration used was Compact but observations of specific sources in SubCompact, Extended and, in two cases, Very Extended were also used. A mixture of full tracks, partial tracks, and track sharing was used for these observations, resulting in different sampling of the uv plane. A summary of the observations made,

¹Specifically data release 4

Table 1. The *Herschel* fluxes of our targets, together with the number of detections in each of the maps, and the threshold used for detections on the dirty map, D_{mapLim} . SPIRE flux densities, in mJy, are given by extractions on the SPIRE maps at the positions of the sources. For sources where we have observations at multiple frequencies (HFLS3, Lock2, and Lock5), we split our maps into different frequencies and detect independently in both of them. HFLS3, Lock2, and Lock5, are observed at both 270 and 345 GHz, labelled as $_{270}$ and $_{345}$, respectively.

Map	Detections	D_{mapLim} (mJy)	F250 (mJy)	F350 (mJy)	F500 (mJy)	E250 (mJy)	E350 (mJy)	E500 (mJy)
Bootes13	2	4.8	54.1	80.6	81.8	7.8	8.1	8.3
Bootes15	0	7.0	51.2	68.3	66.1	6.0	7.0	8.0
Bootes24	0	9.0	28.7	49.7	55.8	6.0	7.0	8.0
Bootes27	0	3.9	37.8	0.0	52.2	6.0	7.0	8.0
Bootes33	1	4.1	50.0	60.9	57.0	8.1	8.7	8.0
Cosmos9	1	2.8	20.6	32.9	47.6	6.1	5.9	6.9
HFLS1	1	10.4	51.6	82.4	82.9	9.8	10.3	9.7
HFLS3_270	1	7.1	15.2	33.4	52.8	6.7	8.0	9.3
HFLS3_345	1	11.3	15.2	33.4	52.8	6.7	8.0	9.3
HFLS5	1	6.2	24.8	53.0	45.8	8.4	9.2	9.1
HeLMS11	1	7.7	83.9	89.6	110.4	15.6	14.9	17.5
HeLMS12	1	17.7	76.9	122.5	138.7	16.7	16.6	17.3
HeLMS28	1	16.8	91.2	137.8	160.5	14.9	14.2	16.9
HeLMS32	1	8.1	36.0	61.5	72.4	16.0	14.8	16.9
HeLMS34	1	13.1	66.7	90.7	106.7	14.7	14.4	16.0
HeLMS35	2	11.1	54.2	91.0	68.6	15.9	15.5	17.9
HeLMS45	0	10.8	32.0	59.0	82.0	14.0	15.0	16.0
HeLMS493	0	7.9	35.0	56.0	68.0	14.0	15.0	16.0
HeLMS62	1	10.6	36.0	30.6	41.0	27.4	17.0	17.1
HeLMS783	1	6.2	42.0	46.6	61.0	15.0	15.0	17.0
LSW28	1	7.0	41.0	65.1	66.9	11.8	11.9	13.2
LSW52	0	4.3	16.3	33.0	40.2	6.0	7.0	8.0
LSW73	2	3.9	19.3	27.4	35.7	8.4	8.6	7.7
Lock2_270	2	5.1	41.4	67.0	65.1	7.6	7.9	8.1
Lock2_345	1	11.6	41.4	67.0	65.1	7.7	8.0	8.1
Lock5_270	1	4.6	17.5	38.1	44.1	8.3	9.1	8.9
Lock5_345	1	11.0	17.5	38.1	44.1	8.3	9.1	8.9
Lock_102	1	15.4	64.0	105.5	122.7	6.3	6.7	7.1
XMM-26	1	7.3	45.7	67.8	81.6	11.9	9.6	8.8
XMM-65	0	6.3	0.0	24.6	42.3	6.0	7.0	8.0
XMM-81	0	6.8	0.0	9.4	28.5	6.0	7.0	8.0
XMM-M2	1	9.4	44.5	54.1	67.6	8.7	9.9	9.4
XMM-M5	0	9.1	26.6	44.0	46.1	6.0	7.0	8.0
XMM-M7	1	9.8	39.7	60.3	57.6	7.4	8.2	9.3
XMM-R2	0	9.2	23.0	32.8	40.3	6.0	7.0	8.0
XMM-R3	0	9.0	22.3	34.9	37.6	8.0	8.7	9.7
XMM-R7	0	8.1	14.3	24.9	37.3	9.8	9.3	9.0

the number of antennae used, and the median opacity at 225 GHz during them is given in Table 2.

Each of the maps was reduced using the IDL based SMA data reduction package MIR,² uniformly in the following way: The data were manually inspected, and any place where individual integrations, spectral bands, baselines, or antennae showed significant flux density offsets or phase changes were flagged and removed from the data. Bandpass calibration used the nearest accessible bright flat spectrum QSO selected from the following sources: 3c279, 3c454.3, 3c84, and BL-Lac. Flux calibration sources were well-characterized Solar System sources selected on the basis of availability from Titan, Uranus, Neptune, Callisto, and Ganymede. These are standard band-pass and flux calibrators used extensively for SMA observations. Gain calibration was done periodically throughout an observation by observing one to two nearby reasonably bright QSOs with known flux densities at 270 and/or 345 GHz. These sources were chosen

from the list of gain calibrators maintained by the SMA for this purpose. All calibration observations used the same tuning as the science observations. The total calibration process results in flux uncertainties of ~ 10 per cent. After calibration, the continuum was generated by averaging all the spectral channels which were not flagged in the inspection step, and the maps were then exported to the data analysis package MIRIAD (Sault, Teuben & Wright 2011) for further reduction.

We used the MIRIAD command INVERT to transform the visibilities into (dirty) maps, choosing a natural weighting (i.e. constant weight to all visibilities) since this maximizes the signal to noise of any sources present. We then CLEANed our maps, using 250 CLEAN iterations. The noise level varies between maps, but on the CLEANed maps was typically around 1.5 ± 0.5 mJy, with minimum and maximum values of 0.5 and 2.9 mJy, respectively. We extract sources and fluxes from the dirty maps rather than the CLEANed maps (see Section 4), and for the dirty maps our noise values had a mean of 2.2 ± 0.9 mJy, with minimum and maximum values of 0.8 and 4.9 mJy, respectively. The beamshape also varied between

²<https://github.com/qi-molecules/sma-mir>

Table 2. The positions, dates, frequencies, array configurations, and project codes for the observations of HerMES selected 500 μm risers, ordered roughly by the date of observation. A list of the SPIRE fluxes at the positions of these sources is given in Table 1. * Also described in Wardlow et al. (2013) with the source name HLock09.

Map	RA 2000	DEC 2000	Date	Freq (GHz)	Config	Project Code	# of Ants	Median Opacity
HFLS1			05/19/10	270	Compact	2010A-S035	7	0.05
			03/14/11	265	SubCompact	2010B-S042	7	0.11
			03/23/11	265	SubCompact	2010B-S042	6	0.11
			08/02/11	265	Extended	2011A-S053	8	0.09
HFLS3			05/12/10	270	Compact	2010A-S035	7	0.08
			03/14/11	265	SubCompact	2010B-S042	7	0.11
			03/23/11	265	SubCompact	2010B-S042	6	0.11
			08/02/11	265	Extended	2011A-S053	8	0.09
			08/03/11	345	Extended	2011A-S053	8	0.09
			07/06/11	336	SubCompact	2011A-S053	7	0.08
Lock2			09/06/11	336	Very Extended	2011A-S053	8	0.06
			05/12/10	270	Compact	2010A-S035	7	0.08
			03/14/11	265	SubCompact	2010B-S042	7	0.11
HFLS5			03/23/11	265	SubCompact	2010B-S042	6	0.11
			05/19/10	270	Compact	2010A-S035	7	0.05
			03/14/11	265	SubCompact	2010B-S042	7	0.11
Lock5			03/23/11	265	SubCompact	2010B-S042	6	0.11
			08/02/11	265	Extended	2011A-S053	8	0.09
			05/19/10	270	Compact	2010A-S035	7	0.05
			03/14/11	265	SubCompact	2010B-S042	7	0.11
XMM_M2			03/23/11	265	SubCompact	2010B-S042	6	0.11
			02/14/12	345	Extended	2011B-S020	8	0.07
			12/08/11	345	Compact	2011B-S038	8	0.07
			12/08/11	345	Compact	2011B-S038	8	0.07
			12/08/11	345	Compact	2011B-S038	8	0.07
			12/09/11	345	Compact	2011B-S038	8	0.08
			12/09/11	345	Compact	2011B-S038	8	0.08
			12/09/11	345	Compact	2011B-S038	8	0.08
			11/27/12	265	Compact	2012B-S018	7	0.12
			11/27/12	265	Compact	2012B-S018	7	0.12
			11/27/12	265	Compact	2012B-S018	7	0.12
			11/27/12	265	Compact	2012B-S018	7	0.12
			11/27/12	265	Compact	2012B-S018	7	0.12
			11/27/12	265	Compact	2012B-S018	7	0.12
			11/27/12	265	Compact	2012B-S018	7	0.12
			11/27/12	265	Compact	2012B-S018	7	0.12
LSW_28			11/27/12	265	Compact	2012B-S018	7	0.06
			11/27/12	265	Compact	2012B-S018	7	0.06
LSW_73			11/27/12	265	Compact	2012B-S018	7	0.06
			11/27/12	265	Compact	2012B-S018	7	0.06
LWS_52			11/27/12	265	Compact	2012B-S018	7	0.06
			11/27/12	265	Compact	2012B-S018	7	0.06
Bootes_27			03/19/13	265	Compact	2012B-S018	8	0.06
			03/19/13	265	Compact	2012B-S018	8	0.06
Bootes_33			03/19/13	265	Compact	2012B-S018	8	0.06
			03/19/13	265	Compact	2012B-S018	8	0.06
Bootes_13			03/19/13	265	Compact	2012B-S018	8	0.06
			03/02/13	302	Compact	2012B-S078	8	0.04
Lock_102*			01/19/13	302	Extended	2012B-S078	7	0.03
			02/03/13	302	Very Extended	2012B-S078	8	0.07
			07/02/13	265	Compact	2013A-S005	6	0.13
HeLMS_28	00:44:09.94	01:18:28.40	07/02/13	265	Compact	2013A-S005	6	0.13
HeLMS_12	00:52:58.52	06:13:18.80	07/02/13	265	Compact	2013A-S005	6	0.13
HeLMS_34	00:22:21.16	-01:55:21.70	07/02/13	265	Compact	2013A-S005	6	0.13
HeLMS_493	23:54:11.79	-08:29:12.00	07/15/13	265	Compact	2013A-S005	6	0.10
HeLMS_45	00:03:04.39	02:40:49.80	07/15/13	265	Compact	2013A-S005	6	0.10
HeLMS_62	00:55:17.18	02:04:01.20	07/15/13	265	Compact	2013A-S005	6	0.10
HeLMS_32	23:24:35.09	-05:24:51.80	07/09/13	265	Compact	2013A-S005	6	0.09
HeLMS_783	00:41:29.86	-02:47:48.10	07/09/13	265	Compact	2013A-S005	6	0.09
HeLMS_11	00:29:36.26	02:07:09.80	07/09/13	265	Compact	2013A-S005	6	0.09
Bootes_15	14:40:09.66	34:37:55.70	06/30/14	346	Compact	2014A-S092	8	0.06
Bootes_24	14:36:21.30	33:02:29.00	06/30/14	346	Compact	2014A-S092	8	0.06
Cosmos_9	10:02:49.11	02:32:57.40	11/28/14	231	Compact	2014A-S092	7	0.08
HeLMS_35	00:21:15.55	01:32:58.60	05/07/15	345	Compact	2014A-S092	7	0.06

maps, and by assuming a Gaussian shape for the central part of the beam, had an average FWHM of 2.5 ± 0.8 arcsec. At $z = 4$, the expected redshift of our 500 μm risers, this corresponds to 17.5 kpc. The DSFGs observed by Oteo et al. (2017b) all have sizes less than 0.6 arcsec so it is unlikely that our sources will be resolved. Multiple

sources or signs of lensing within an ~ 2.5 arcsec radius would also not necessarily be resolved. Experiments were performed using different weighting schemes and cleaning parameters, using both the data and simulations (see Section 3.2). The modelling discussed in Section 3.2 found that the process described above generally per-

formed best at recovering the flux density of injected fake sources. The noise levels and beamsizes in our maps are in reasonable agreement with what Hill et al. (2018) found in their recent SMA reduction of a number of SCUBA-2 selected sources. The resulting dirty maps, beamshapes, and CLEANed maps for all of our 500 μm risers can be found in Appendix A. The CLEANed maps are not used for the scientific analysis but are included for presentational purposes and to show how the complex beam shapes in some of the dirty maps can mimic the presence of multiple sources.

3.2 Source extraction

We extract sources from the dirty maps on the basis of SNR with the noise calculated from the whole dirty map, without correction for the primary beam response. We also examined using noise calculated within 1–2 beam FWHM radii of the source instead of the global value.

To determine the reliability of source extraction, we perform the following steps. First, for each of our maps, we create 1000 fake maps filled with Gaussian noise. These are then convolved with the beam for that map, and this fake dirty map is scaled until the noise properties match those of the original dirty map (i.e. same mean and standard deviation). This provides 1000 realizations of each map but without sources. The same extraction method as on the real data is then applied to these empty maps.

The best extraction method will minimize the number of fake sources detected, whilst simultaneously minimizing the SNR threshold used for detection, so as to maximize the depth of our catalogue. We found that using the global SNR measured on the dirty map performed best, with only 0.09 per cent fake sources detected at $>5\sigma$, compared to the local measurement on the dirty map, which had 2 per cent fake sources detected at $>5\sigma$. We also experimented with source extraction on the CLEANed maps but found this to have even higher rates of fake sources.

The average number of fake sources detected as a function of SNR varies significantly between maps. This is likely because the synthesized beams can be very different. We thus select an SNR cutoff for each map such that there is an average of 0.1 fake sources per map. This SNR cutoff varies from map to map, ranging from 3.54 to 3.86σ , but is on average 3.75σ , where σ is the global noise on the dirty map. Over all our 34 sources, we would thus expect 3.4 ± 1.8 spurious, low SNR sources.

We test the completeness of our extraction method by injecting fake sources into the real maps and then attempting to recover them. To do this, we copy each of our observed maps 1000 times, and inject a point source convolved with the beam with a flux distributed uniformly between 0 and 30 mJy. We require the injected sources to be at least 5 arcsec away from any detected source in the map. These maps then have sources extracted and we determine what fraction of injected sources are recovered within one synthesized beamwidth of the injected position. We find significant variations between the maps, with maps that have more complex beams (e.g. HeLMS28) having worse source recovery than maps with better UV coverage (e.g. Bootes13). We find on average that we recover 65.3 per cent of sources. Of the remaining 34.7 per cent, most are faint, with only 6.6 per cent having flux densities >10 mJy. These are usually in maps where a bright source is already present, raising the global noise.

We also use our injected sources to examine two alternative methods for extracting flux densities: using Gaussian fitting, or fitting the beam to the dirty map at the position of the peak pixel, in addition to measuring the peak flux from the map. We find that

extracting the peak flux from the map minimized the difference between injected and extracted flux densities, compared to both Gaussian and beam fitting. We note that we only injected point sources, so the alternate methods may perform better for resolved sources. We find good agreement between injected and recovered fluxes for flux densities >10 – 15 mJy. Below this, the injected flux is within a few standard deviations of the noise of most of our maps, and while the recovered flux is generally within 2σ of the injected value, the (recovery – injected) to injected flux density ratio diverges.

Where we have additional extended or very extended observations (EXT or VEXT) of a given source, we find that maps combining these with compact or subcompact (COM or SUB-COM) observations produced significantly lower recovered fluxes compared to compact (or subcompact) observations on their own. The reason for this remains unclear, but phase decoherence or partially resolving the source in extended configurations are likely explanations. We thus extract flux densities only using compact or subcompact observations. In Fig. A1, these are indicated with the addition of ‘_COM_’ to the map name.

4 RESULTS

4.1 Detected sources

We apply the source detection and extraction methods described above to our maps of 500-risers, finding 26 sources associated with our 34 *Herschel* selected 500-risers. Three of these sources are detected at two different frequencies in the SMA maps. The number of SMA sources associated with each 500-riser is given in Table 1 while the individual SMA sources are detailed in Table 3.

Of our 34 500-risers, 12 (35 per cent) have no detections with the SMA, 18 (53 per cent) have a single detection, and 4 (12 per cent) have two detections, indicating that at these flux density limits and resolutions, most 500 μm risers appear to be single sources. The sample size is fairly small, but there is no evidence at these flux densities to suggest any correlation between the number of detections and the detection threshold used in the map, with the two, one, and no detection subsets having mean detection thresholds of 6.2 ± 3.3 , 9.4 ± 4.0 , and 7.6 ± 2.0 mJy, respectively. There are also similar numbers of detections for the 345 and 270 GHz maps, though there are only 13 maps at 345 GHz and 24 at 270 GHz. The flux densities of detected sources range between 3.62 to 35.7 mJy at 270 GHz, with a mean of 12.3 mJy, and between 13.0 and 72.1 mJy at 345 GHz, with a mean of 22.96 mJy.

Of the 17 SPIRE sources with $S_{500} > 60$ mJy, which Béthermin et al. (2017) predict to be almost entirely single sources, we find 12 that appear to be single sources, 3 have multiple detections, and 2 have zero detections. One of the multiple sources, Lock2, has one counterpart at 345 GHz but two at 270 GHz. Curiously, the source detected at both 345 and 270 GHz is the fainter of the two sources at 270 GHz, with only 5.68 ± 0.94 mJy at 270 GHz compared to the other source which has 9.07 ± 0.94 mJy but no counterpart at 345 GHz. As we discuss later, the maps where we detect no sources are likely be made up of multiple faint sources below the detection limit. This means that of our 17 SPIRE sources with $S_{500} > 60$ mJy, 5 (30 per cent) are multiple galaxies. Of the 18 sources in Oteo et al. (2017b) with $S_{500} > 60$ mJy, they find 6 (33 per cent) that are resolved into multiple galaxies, in good agreement with our SMA result. Such a high fraction would appear to disagree with Béthermin et al. (2017), who suggest such sources are almost entirely lensed single sources.

Table 3. Properties of our 29 detected sources. The redshift, if known, is given in the final column. In three cases sources are detected at two different frequencies. * SNR extracted from the dirty map. Redshift sources are as follows: HFLS1, 3, 5, and Lock_102 from Riechers et al. (2013); Lock5 from Dowell et al. (2014); HeLMS34 from Asboth et al. (2016), HeLMS12 and 28 from Riechers et al. (in preparation), and Duivenvoorden et al. (2018). * Lock_102_ID.1 is our most extended source so this flux is derived not from the point source extraction applied to all of the other sources but instead extracts flux from a 4 arcsec diameter aperture on the COM configuration map. The flux extracted from the EXT map in a similar way is 68.8 mJy, which is within $\sim 1\sigma$, but we use the COM flux to minimize the possibility of flux being resolved out. The point source extracted flux for this source is 55.0 mJy.

Name	RA 2000	DEC 2000	Frequency (GHz)	Flux (mJy)	Flux_Err (mJy)	SNR*	z
Bootes13ID.1	218.93331	34.79416	265	10.4	1.25	8.32	–
Bootes13ID.2	218.93281	34.79787	265	9.35	1.25	7.48	–
Bootes33ID.1	217.88329	34.35276	265	6.23	1.08	5.77	–
Cosmos9ID.1	150.70492	2.54869	231	3.62	0.66	5.48	–
HFLS1_ID.1	257.07377	58.47777	265	15.8	2.87	5.50	4.29
HFLS3_270_ID.1	256.69914	58.77330	270	18.4	1.36	13.5	6.34
	256.69926	58.77330	336	24.6	1.79	13.7	6.34
HFLS5_ID.1	260.20754	59.77442	270	13.3	1.32	10.1	4.44
HeLMS11ID.1	7.40011	2.12036	265	10.0	2.07	4.83	–
HeLMS12ID.1	13.24405	6.22189	265	38.6	3.80	10.1	4.37
HeLMS28ID.1	11.04294	1.30615	265	34.4	3.75	9.17	4.17
HeLMS32ID.1	351.14402	– 5.41560	265	12.6	1.90	6.63	–
HeLMS34ID.1	5.58661	– 1.92247	265	26.0	3.25	8.00	5.16
HeLMS35ID.1	5.31452	1.54934	345	19.1	2.61	7.32	–
HeLMS35ID.3	5.31506	1.55261	345	13.0	2.61	4.98	–
HeLMS62ID.2	13.82280	2.06797	265	11.9	2.91	4.09	–
HeLMS783ID.1	10.37588	– 2.79840	265	7.86	1.72	4.57	–
LSW28ID.1	165.36617	56.32630	265	19.2	1.71	11.2	–
LSW73ID.1	164.08463	56.90126	265	4.19	1.00	4.19	–
LSW73ID.2	164.07629	56.89420	265	4.38	1.00	4.38	–
Lock2_270ID.1	163.29714	56.70437	270	9.07	0.94	9.65	–
Lock2_270ID.2	163.29509	56.70212	270	5.68	0.94	6.04	–
	163.29503	56.70220	345	13.1	3.0	4.36	–
Lock5_270ID.1	162.88248	56.60535	270	8.0	0.80	10.0	3.36
	162.88288	56.60542	346	14.2	2.92	4.86	3.36
Lock_102_ID.1	160.21099	56.11496	302	72.1*	2.99	24.1	5.29
XMM-26ID.1	36.43590	– 2.98359	265	7.62	1.92	3.96	–
XMM-M2ID.1	36.31398	– 2.78553	345	14.4	2.51	5.73	–
XMM-M7ID.1	36.68632	– 3.41962	345	13.3	2.55	5.21	–

We can only say with certainty that the maps where we detect two sources are definitely multiples; in maps where we detect a single source, there may be a secondary fainter source, and in maps where we detect no sources, we can at best put upper limits to any detection. For maps with one detection, the ratio between the brightest (detected) source and any secondary source can be a maximum of $\frac{S_{\text{Lim}}}{S_{\text{Detected}}}$, where S_{Lim} is the upper limit for detection, and S_{Detected} is the flux density of the detected source. We find that in this case, any secondary source would have a flux density *at most* between 98 per cent and 28 per cent that of the detected source, with an average of (63 ± 19) per cent. These upper limits are reasonably consistent with Donevski et al. (2018), who find that the brightest source that makes up a Herschel detected 500 μm riser contributes around 60 per cent of the 500 μm flux density. However, at the current sensitivities our estimates remain upper limits, though we note that in the four maps where we have two detections the brighter source contributes on average 60 per cent to the total flux density from both sources. For now, the simplest solution to the maps where we detect only one source is that this sole source is responsible for the SPIRE flux, and we move ahead under this assumption. The same cannot be said for those maps where we detect no sources.

4.2 Maps without detections

In twelve maps, we detect no sources at all. This is surprising: all but one of these maps have Herschel 500 μm flux densities >40 mJy, while Bootes15, HeLMS493, and HeLMS45 have 500 μm flux densities >60 mJy. The detection threshold in these dirty maps is $\sim 7\text{--}10$ mJy, making it difficult to explain these non-detections as single sources that are below our detection threshold. Examining Fig. A1, there is some evidence of a CLEANed source in HeLMS493, with a CLEANed signal-to-noise ratio of 5.61 and a flux density of 9.85 ± 1.76 mJy, but this source does not pass our threshold in the dirty map, as it has a signal-to-noise ratio of only 2.1 and a flux density of 4.65 ± 2.18 mJy. Numerous (>2) DSFGs which appear as single sources in the Herschel maps are known to be multiple at higher resolutions (Oteo et al. 2017a; Oteo et al. 2018). To determine how many sources would reproduce the typical SPIRE flux in the maps where we do not have any SMA detections, we take the ALESS average SED (da Cunha et al. 2015) at redshifts between 4 and 6, and normalizable its flux density at 500 μm to 60 mJy. If a single source were responsible for the SPIRE flux, we find that it should have 345 and 270 GHz flux densities of $\sim 24\text{--}38$ and $14\text{--}25$ mJy, respectively, and that we should have detected

Table 4. The derived properties of our sources from fitting single modified blackbodies, *uncorrected for lensing* μ . Lock2 was fitted only using the source detected at both 270 and 345 GHz.

Name	$\log_{10}(L_{\text{FIR}}/L_{\odot})$	$\log_{10}(M_{\text{Dust}}/M_{\odot})$	β	SFR ($M_{\odot}\text{yr}^{-1}$)
Bootes33ID_1	13.4 ^{+0.26} _{-0.86}	8.75 ^{+0.76} _{-0.52}	2.29 ^{+0.43} _{-0.47}	2800 ⁺¹⁴⁰⁰ ₋₁₉₀₀
Cosmos9ID_1	13.3 ^{+0.26} _{-0.85}	8.64 ^{+0.98} _{-0.55}	2.55 ^{+0.31} _{-0.45}	2200 ⁺¹²⁰⁰ ₋₁₄₀₀
HFLS1	13.5 ^{+0.07} _{-0.13}	9.28 ^{+0.24} _{-0.13}	1.80 ^{+0.56} _{-0.46}	3500 ⁺⁷⁰⁰ ₋₆₀₀
HFLS3	13.5 ^{+0.2} _{-0.21}	9.05 ^{+0.28} _{-0.13}	1.88 ^{+0.69} _{-0.60}	3500 ⁺⁴⁰⁰ ₋₄₀₀
HFLS5	13.3 ^{+0.16} _{-0.14}	9.14 ^{+0.29} _{-0.20}	1.56 ^{+0.67} _{-0.39}	2200 ⁺⁷⁰⁰ ₋₅₀₀
HeLMS11ID_1	13.5 ^{+0.33} _{-0.76}	8.96 ^{+0.80} _{-0.52}	2.32 ^{+0.45} _{-0.56}	3500 ⁺²²⁰⁰ ₋₂₄₀₀
HeLMS12ID_1	13.7 ^{+0.12} _{-0.13}	9.62 ^{+0.23} _{-0.14}	1.49 ^{+0.68} _{-0.35}	5500 ⁺¹⁴⁰⁰ ₋₁₀₀₀
HeLMS28ID_1	13.7 ^{+0.09} _{-0.09}	9.64 ^{+0.18} _{-0.17}	1.66 ^{+0.60} _{-0.40}	5500 ⁺¹²⁰⁰ ₋₉₀₀
HeLMS32ID_1	13.5 ^{+0.38} _{-0.82}	8.93 ^{+0.82} _{-0.47}	2.06 ^{+0.59} _{-0.63}	3500 ⁺³⁰⁰⁰ ₋₂₃₀₀
HeLMS34ID_1	13.7 ^{+0.17} _{-0.07}	9.26 ^{+0.17} _{-0.14}	1.51 ^{+0.66} _{-0.36}	5500 ⁺¹⁸⁰⁰ ₋₁₁₀₀
HeLMS62ID_2	13.6 ^{+0.65} _{-0.93}	9.04 ^{+0.87} _{-0.56}	1.86 ^{+0.72} _{-0.61}	4400 ⁺⁷⁴⁰⁰ ₋₂₂₀₀
HeLMS783ID_1	13.3 ^{+0.41} _{-0.88}	8.79 ^{+0.75} _{-0.48}	2.02 ^{+0.65} _{-0.65}	2200 ⁺²⁶⁰⁰ ₋₁₅₀₀
LSW28ID_1	13.6 ^{+0.37} _{-1.2}	9.20 ^{+0.70} _{-0.44}	1.50 ^{+0.64} _{-0.36}	4400 ⁺⁵⁰⁰⁰ ₋₃₁₀₀
Lock2	13.4 ^{+0.18} _{-0.84}	8.70 ^{+1.07} _{-0.56}	2.57 ^{+0.3} _{-0.48}	2800 ⁺⁸⁰⁰ ₋₂₇₀₀
Lock5	13.8 ^{+0.15} _{-0.15}	9.39 ^{+0.39} _{-0.36}	2.06 ^{+0.64} _{-0.68}	6900 ⁺¹¹⁰⁰ ₋₁₈₀₀
Lock_102	13.8 ^{+0.05} _{-0.02}	10.21 ^{+0.73} _{-0.42}	2.01 ^{+0.69} _{-0.76}	6900 ⁺⁹⁰⁰ ₋₄₀₀
XMM-26ID_2	13.6 ^{+0.21} _{-0.74}	8.82 ^{+0.99} _{-0.55}	2.65 ^{+0.25} _{-0.40}	3900 ⁺²²⁰⁰ ₋₃₀₀₀
XMM-M2ID_1	13.4 ^{+0.22} _{-0.84}	8.75 ^{+0.81} _{-0.49}	2.51 ^{+0.35} _{-0.56}	4400 ⁺¹²⁰⁰ ₋₁₆₀₀
XMM-M7ID_1	13.4 ^{+0.21} _{-0.70}	8.66 ^{+0.90} _{-0.48}	2.41 ^{+0.40} _{-0.58}	4400 ⁺¹⁰⁰⁰ ₋₁₅₀₀

the source responsible. If the flux were equally split between two sources, we would expect them to have SMA flux densities of ~ 12 – 19 and 7 – 12 mJy, respectively; low, but, given the detection limits in Bootes15 and HeLMS493, still detectable. This is not a detailed analysis, but it suggests that in at least some of the maps where we do not detect any sources there may be ≥ 3 DSFGs separated by ~ 2 – 20 arcsec (the SMA and SPIRE beam sizes, respectively). This raises the intriguing possibility that such sources may be protocluster cores, with numerous far-IR luminous sources spread over a ~ 20 arcsec region, such as the Distant Red Core (Oteo et al. 2018) or SPT2349-56 (Miller et al. 2018). We have thus begun a programme to obtain deeper observations of these ‘blank field’ sources both at high resolution with the SMA, and with the JCMT to determine integrated fluxes (Cairns et al., in preparation).

5 DISCUSSION

5.1 The physical properties of 500-risers

To determine the physical parameters of our sources, we take the measured SPIRE and SMA flux densities, and, following literature convention (Blain 2002; Magnelli et al. 2012; Bianchi 2013; Casey, Narayanan & Cooray 2014), fit them to a single temperature modified blackbody function of the form

$$S_{\nu} \propto (1 - \exp(-\tau_{\nu}))B_{\nu}(T), \quad (1)$$

where S_{ν} is the observed flux density at frequency ν , $\tau_{\nu} = (\frac{\nu}{\nu_0})^{\beta}$ gives the optical depth at frequency ν , ν_0 is the frequency at which the optical depth equals unity, and $B_{\nu}(T) = B_{\nu}(\nu, T_{\text{dust}})$ is the Planck function at temperature T . β , the spectral index, is usually assumed to be $\beta = 1.5 - 2$ for SMGs (Blain 2002; Casey et al. 2014). Since it is unclear how to distribute SPIRE flux between multiple SMA sources, we do not include these sources in this analysis.

We use the affine invariant Markov Chain Monte Carlo (Goodman & Weare 2010) ensemble sampler PYTHON package EMCEE

(Foreman-Mackey et al. 2013) to fit this model to our data, with the following uninformative priors for our parameters: $0 < z \leq 12$, $T_{\text{CMB}}(z) \leq T_{\text{dust}} \leq 80$, $1 \leq \beta \leq 3$, $1 \mu\text{m} \leq c/\nu_0 \leq 1 \text{mm}$, and $-2 \leq \log_{10}(a) \leq 2$, where T_{CMB} gives the CMB temperature at redshift z , and c gives the speed of light. For numerical stability, at each sample we first normalize to the $500 \mu\text{m}$ observation, and allow the normalization a to vary from there. In this model, there are therefore five parameters to fit: the redshift z , the average dust temperature T_{dust} , the spectral index β , ν_0 , and an overall normalization parameter a . Where a redshift is known the value of z in the fit is fixed to this value. We run our sampler with 100 walkers and for 3000 steps and, after manually checking to ensure the model has burnt in, throw away the initial 300 steps as a burn in phase. The dust temperature and redshift that emerge from such fits are strongly degenerate, so these values do not provide any direct physical insight into our DSFGs. However, the far-IR luminosity, dust mass, and star formation rate that can be derived from these fits suffer much less from this degeneracy (see e.g. Greenslade et al. 2019). The far-IR luminosity (defined as the integrated luminosity from 42.5 to $122.5 \mu\text{m}$) is calculated by integrating the fitted SEDs between these two rest frame frequencies. This is done for all samples on each source to produce the posterior luminosity distribution of each source marginalized over all redshifts. The dust mass is then calculated following Riechers et al. (2013) using

$$M_{\text{dust}} = S_{\nu} D_L^2 [(1+z)\kappa_{\nu} B_{\nu}(T)]^{-1} \tau_{\nu} [1 - \exp(-\tau_{\nu})]^{-1}, \quad (2)$$

where S_{ν} gives the rest-frame flux density at $125 \mu\text{m}$, D_L is the luminosity distance, κ_{ν} is the mass absorption coefficient, and is assumed to be $\kappa_{\nu} = 2.64 \text{m}^2 \text{kg}^{-1}$ at $125 \mu\text{m}$ (Dunne, Eales & Edmunds 2003), and other symbols are defined as before. Finally, the FIR luminosity is converted to an SFR following Riechers et al. (2013), using a Chabrier stellar IMF (Chabrier 2003). These results can be found in Table 4 where we show the median, 14th and 86th percentiles of these distributions for each of our sources. The resulting SED fits compared to the observational data are shown in Fig. 1.

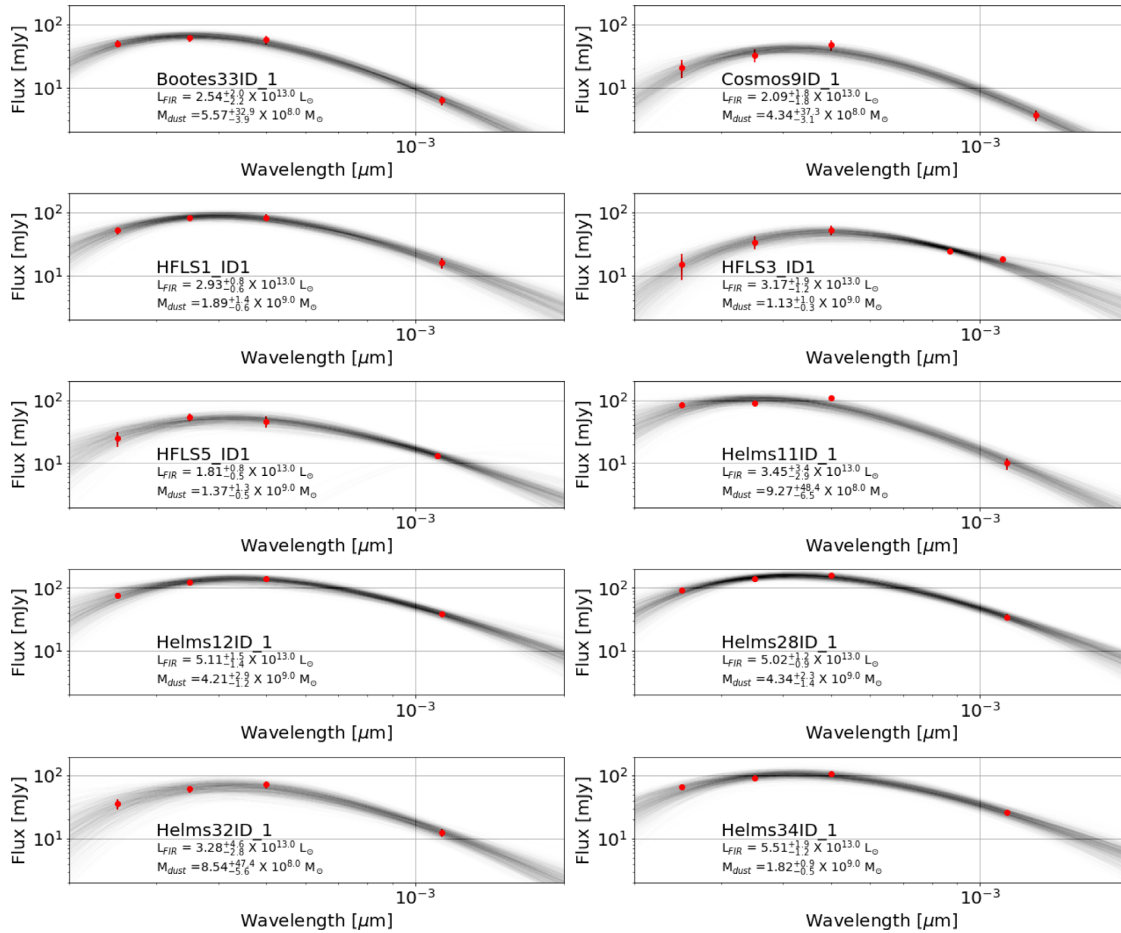
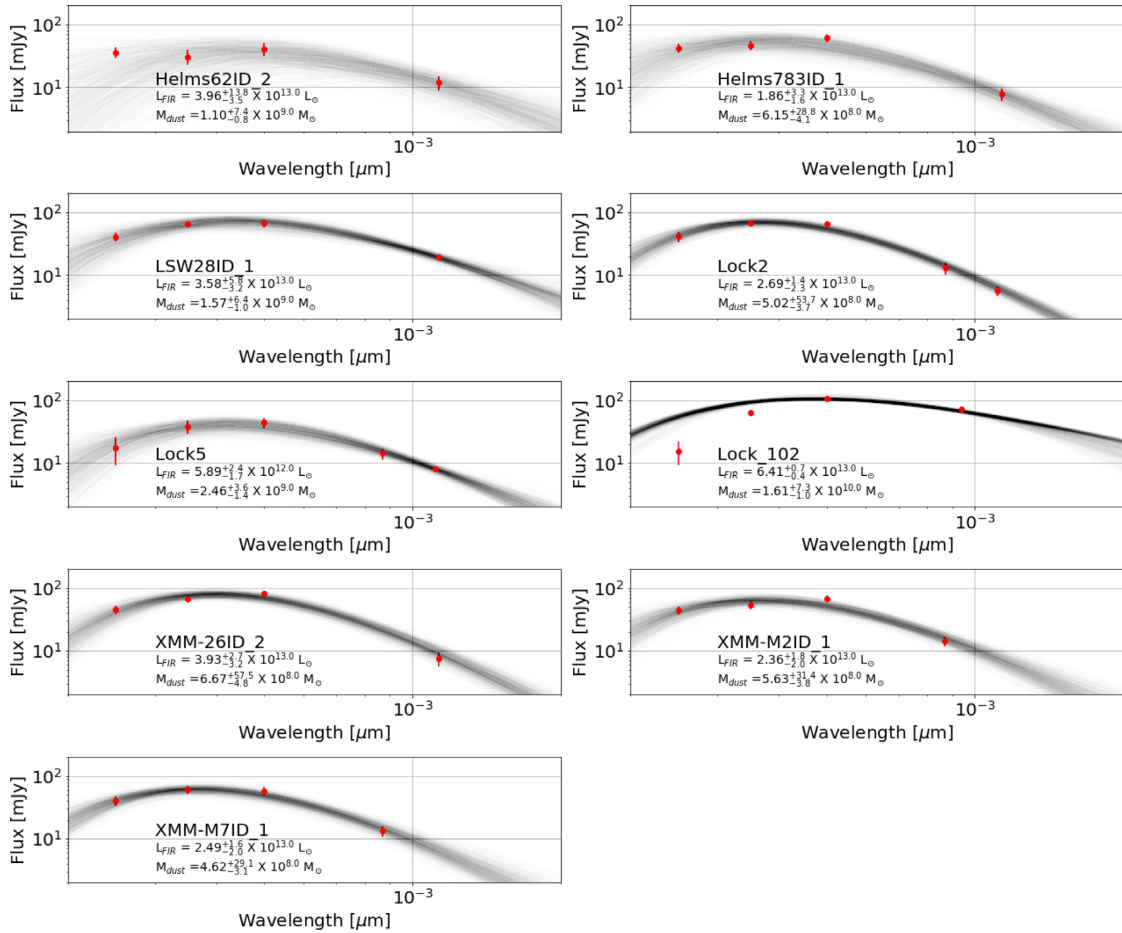


Figure 1. Sampled SED fits to the SPIRE and SMA flux densities measured for sources where we only detect a single counterpart to the SPIRE source. The black lines show 1000 samples after fitting, while the red points give the flux densities and error bars to the observed photometry of each source. Names, estimated FIR luminosities, and estimated dust masses are provided as text on each of the plots. The poor fit for Lock_102 results from the presence of the [C II] 157 μm line in the bandpass of these observations, contaminating and enhancing the submm flux derived from our SMA data (see Perez-Fournon et al., in preparation).

As can be seen from Table 4, those sources where we detect a single SMA counterpart all have luminosities $> 10^{13} M_{\odot}$, leading to their classification as HLIRGs. While lensing may have boosted the derived luminosity of some sources, a modest lensing amplification at the level of 2–3, as is the case for HFLS3 (Cooray et al. 2014) would still leave most of our sources as HLIRGs. In the absence of any AGN contribution to their high far-IR luminosity, the derived SFR for these objects is in excess of $1000 M_{\odot} \text{yr}^{-1}$, with some sources potentially having SFRs in excess of $5000 M_{\odot} \text{yr}^{-1}$, leading to their classification as Extreme Starbursts (Rowan-Robinson et al. 2018), a class of object that cannot be accounted for by any current model of galaxy formation. While the error bars on these derived properties are large, with many lying within 1σ of HLIRG luminosity, a number, such as HeLMS34, are well above this value. The simulations used by Béthermin et al. (2017) to determine the effects of source clustering on single dish instruments limited the maximal starburst to $1000 M_{\odot} \text{yr}^{-1}$. Similar to the higher resolution ALMA observations of Oteo et al. (2017b), we here show that sources beyond this limit exist in the 500-riser population, so it is likely that the results of the clustering simulations are in need of revision at the brightest fluxes. Further observations at complementary wavelengths, including the optical, X-ray, and far-IR spectroscopy, will be needed to determine what role AGN play in powering these sources.

5.2 The role of lensing

One of our sources is known to be strongly lensed (Lock_102), and another (HFLS3) is weakly lensed with a magnification of 1.3. Seventeen of our sources have $S_{500} > 60 \text{ mJy}$, which Béthermin et al. (2017) suggests should all be lensed galaxies. However, we only have extended array observations for five of our sources (HFLS1, HFLS3, HFLS5, Lock5, and Lock_102), and the 2–4 arcsec beamsizes in the compact and subcompact configurations lack the resolution necessary to provide a definitive test of lensing. In Fig. 2, we show those sources where we do have extended or very extended configuration observations. The synthesized beams for these range from 1 to 0.35 arcsec across. Only one of these sources, Lock_102, clearly shows the arc-like structure indicating strong lensing, which is confirmed by observations at other telescopes. These observations were at a resolution of $\sim 0.35 \text{ arcsec}$. The other known lens, HFLS3 with a magnification $\mu = 2.2 \pm 0.3$ (Cooray et al. 2014), appears reasonably compact, with some sign of extension to the north, but no obvious indications of lensing. These observations were also at a resolution of $\sim 0.35 \text{ arcsec}$. Similar suggestions of extension are seen in HFLS1 and Lock5 (resolutions of 1 and 0.8 arcsec, respectively) whilst HFLS5 (resolution 1 arcsec) appears compact, indicating a diameter $< 6.9 \text{ kpc}$.

Figure 1. *continued*

Oteo et al. (2017b) in their ALMA examination of 500 μm risers find 40 per cent of their sources display some lensing signature (arcs, elongation, or rings). This is a significantly higher fraction than we find here, but the superior resolution and sensitivity offered by ALMA, with an average beam FWHM of only 0.12 arcsec, makes this kind of analysis much easier. Nevertheless, both the current work and Oteo et al. (2017b) find that the lensing fraction for 500-risers with flux densities ≥ 60 mJy seems to be much lower than predicted by Béthermin et al. (2017).

5.3 The role of blended sources

Several authors have claimed, mostly through simulations, that most Herschel sources are multiple galaxies blended together in the large Herschel beam (Cowley et al. 2014; Scudder et al. 2016, 2018; Béthermin et al. 2017). Cowley et al. (2014) focused on 850 μm selected sources and Scudder et al. (2016, 2018) focused on 250 μm selected sources, so their results are not well matched to the current results on 500-risers. In contrast, Béthermin et al. (2017) did look at 500 μm selected sources. They suggest that on average 60 per cent of a $S_{500} < 60$ mJy Herschel flux density at 500 μm comes from a single source, while sources brighter than 60 mJy should largely be strongly lensed single sources. Our observations show a more complicated picture. Of the 17 sources in our sample with $F_{500} > 60$ mJy, three are found to have two counterparts and three are found to have no counterparts at all, suggesting that they are blends of three or more sources. This means that 35 per cent of

the sources in our sample with flux densities > 60 mJy at 500 μm do not match the expectation that they should be single, strongly lensed sources. Instead they are blends of multiple components. At brighter fluxes, $S_{500} > 100$ mJy, they do appear to be single sources, with one, Lock_102, being a confirmed lens. At somewhat fainter fluxes, where lensing is still expected to dominate, we have counterexamples such as Bootes13, which resolves into two sources ($S_{500} = 82$ mJy) and Bootes15, where we detect no counterparts despite a flux density of $S_{500} = 66$ mJy. At fainter fluxes, $S_{500} < 60$ mJy, where 40 per cent of sources are expected to be multiple, we find that of 17 sources, seven are single sources, one resolves into two components, and nine have no detection, suggesting that they are made up of multiple components, suggesting a multiplicity fraction of 60 per cent.

In Fig. 3, we plot a histogram of the 500 μm flux density of our sources split by whether we do or do not detect counterparts in our SMA maps. The maps where we do not detect any sources are generally fainter at 500 μm . This tends to support the conclusion of Béthermin et al. (2017) that sources below 60 mJy are more likely to be blends. However, with a likely multiplicity fraction of 60 per cent our data suggests that the blended fraction is likely to be higher than the predicted 40 per cent.

Comparing our results to those of Oteo et al. (2017b) is complicated by the fact that our observations were taken with a range of sensitivities, frequencies, and configurations. Our best sensitivities would be able to detect at least two separate sources in all of their multiple sources, while our best resolution would be able to resolve

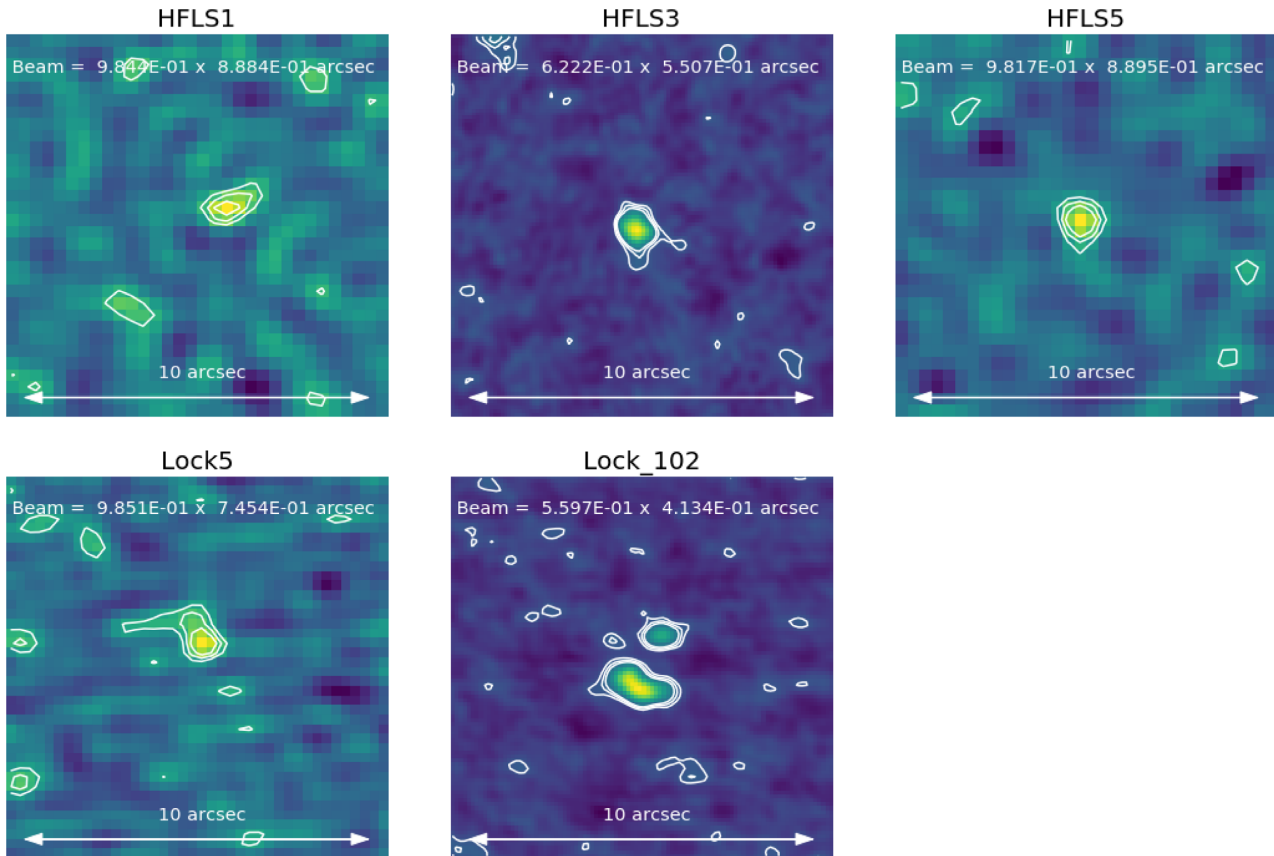


Figure 2. Extended (all maps) and very extended (HFLS3 and Lock_102) only observations on maps where these configurations are available. All maps are 10 arcsec \times 10 arcsec, and orientated so that north is up and east is to the left. Contours show 2, 3, 4, and 5 σ levels. The beam is approximated as a Gaussian, with semimajor and semiminor axis given in the text at the bottom of each image.

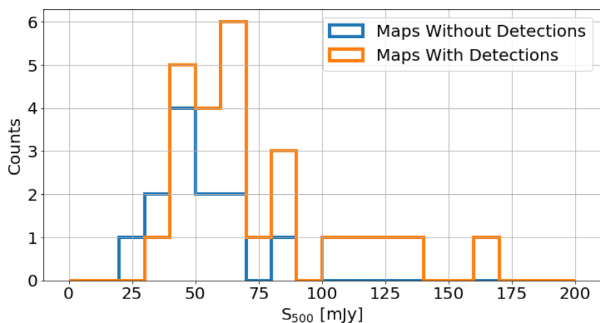


Figure 3. The Herschel 500 μ m flux density of a source, as a function of whether a source is detected in the SMA map (orange histogram) or not (blue histogram). Bins are of width 20 mJy.

all the multiple components in all of their targets. More typical observations, however, with a resolution of 2 arcsec and a 345 GHz equivalent sensitivity of 2 mJy, would miss many of the multiple components found in Oteo et al. (2017b), with most being classified as single sources and three being classified as blank fields. This suggests that the true multiplicity fraction may be even higher than we report above, implying a greater disagreement with the models. Differences in target selection for the two samples, however, mean that this is not yet a solid conclusion. Our comparison with Oteo et al. (2017b) also suggests that deeper observations of our blank field sources will likely uncover further detections, something that

our continuing SMA observations have already demonstrated for some sources (Cairns et al., in preparation).

6 CONCLUSIONS

We have conducted SMA interferometric observations of 34 candidate high redshift ($z > 4$) DSFGs selected from *Herschel* survey data. The sources were selected to be 500-risers, i.e. sources whose redshifted far-IR dust SED is still rising in the SPIRE 500 μ m band. This has been found to be a good indicator of high redshift. Theoretical models (B  thermin et al. 2017) have ascribed the unexpectedly high number counts of such 500-riser sources to the effects of gravitational lensing for bright sources (flux densities >60 mJy) and blending of multiple sources in the large *Herschel* beams at lower flux densities. Our observations, at resolutions from 0.35 to 3 arcsec, detect counterparts to 22 of our 34 sources, with those sources lacking counterparts likely to be made up of three or more blended sources. We examine the rates of detection of counterparts above and below the 60 mJy 500 μ m flux density break above which B  thermin et al. (2017) suggests all sources should be strongly lensed, and below which they suggest 40 per cent should be blends. We find that bright (>60 mJy) 500 μ m sources are a more diverse population with ~ 35 per cent being blends rather than single sources. Fainter (< 60 mJy) sources are found to be blends rather more often than predicted. We use the *Herschel* fluxes and our new SMA fluxes to determine the SED properties of the single sources we have detected, and then to measure their far-IR luminosity and

SFR. We find that all our targets are HLIRGs, i.e. have $L_{\text{FIR}} > 10^{13} L_{\odot}$, and with SFRs in excess of $> 1000 M_{\odot} \text{ yr}^{-1}$ if the effects of gravitational lensing are ignored. The 500-riser population thus seems likely to represent galaxies undergoing major bursts of star formation at very high redshifts. Such high luminosity objects were not considered in the Béthermin et al. (2017) models. Determining the broader role of this population in galaxy formation, and whether, as is likely for any reasonable luminosity function, there is an underlying much larger population of lower luminosity high redshift DSFGs, remain important tasks for far-IR and submm astronomy.

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APPENDIX A: SMA IMAGES

We here present the dirty maps, beam shapes, and cleaned maps for all at targets at all their observed frequencies. All images are 20 arcsec on a side. North is up and east is to the left.

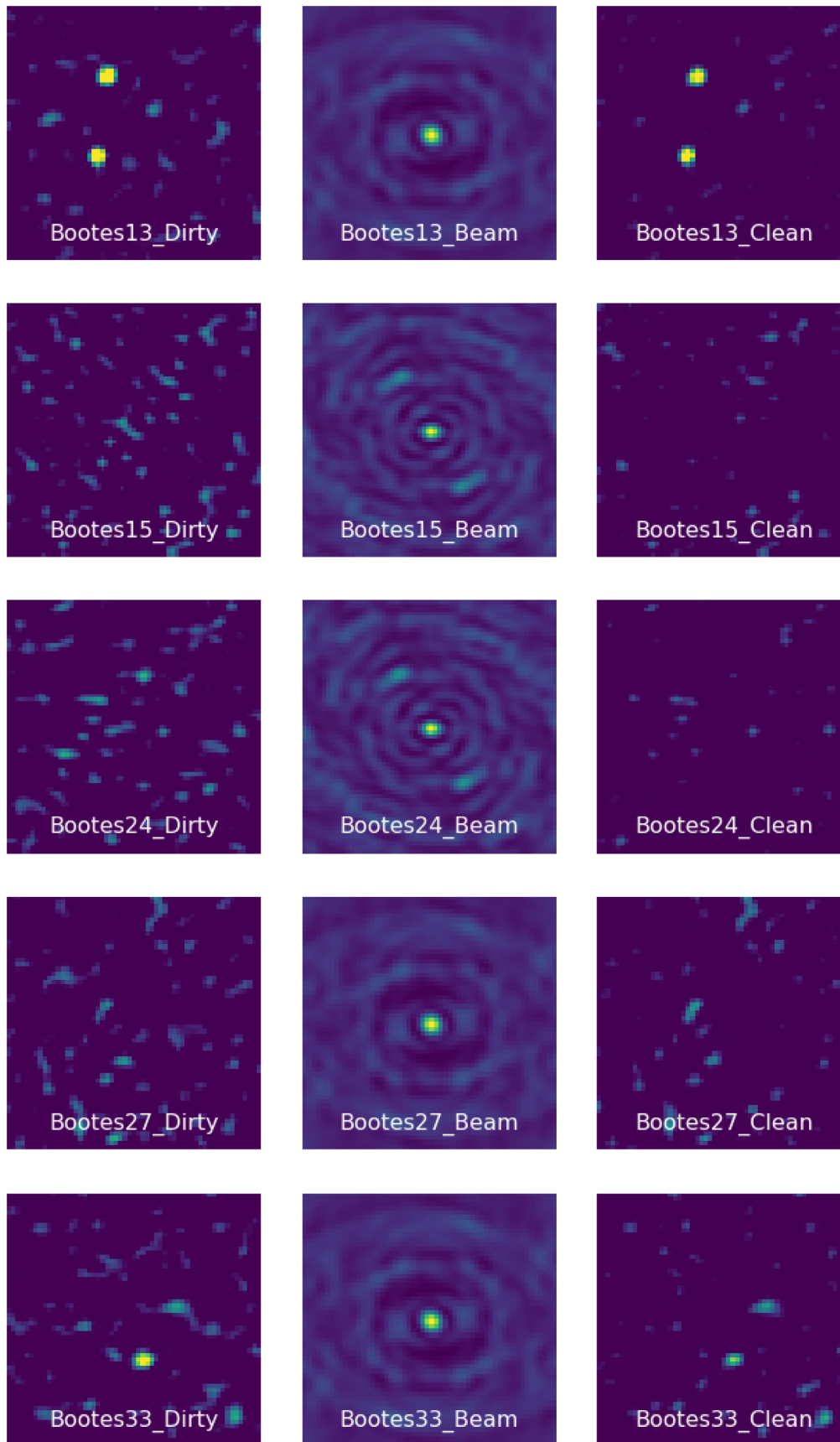
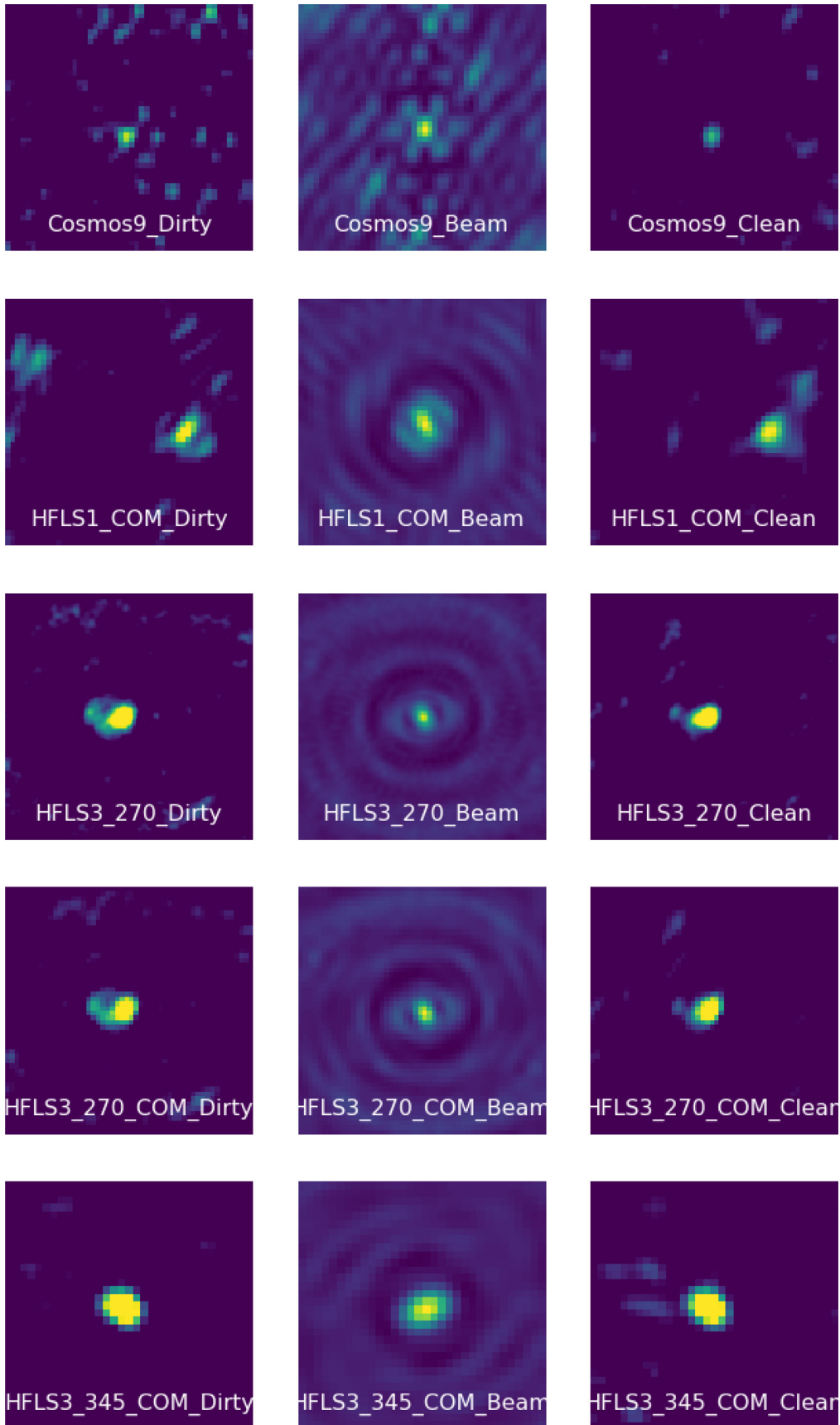


Figure A1. Images of the dirty (left-hand column), beam (middle), and CLEANed (right-hand column) maps of our 34 sources. The colour scheme in the dirty and clean maps ranges between 1 and 5σ , using the rms on the dirty map. Maps which include the `.COM_` label indicate where we have only included the visibilities from the compact or subcompact array configurations (see the text for more details).

Figure A1. *continued.*

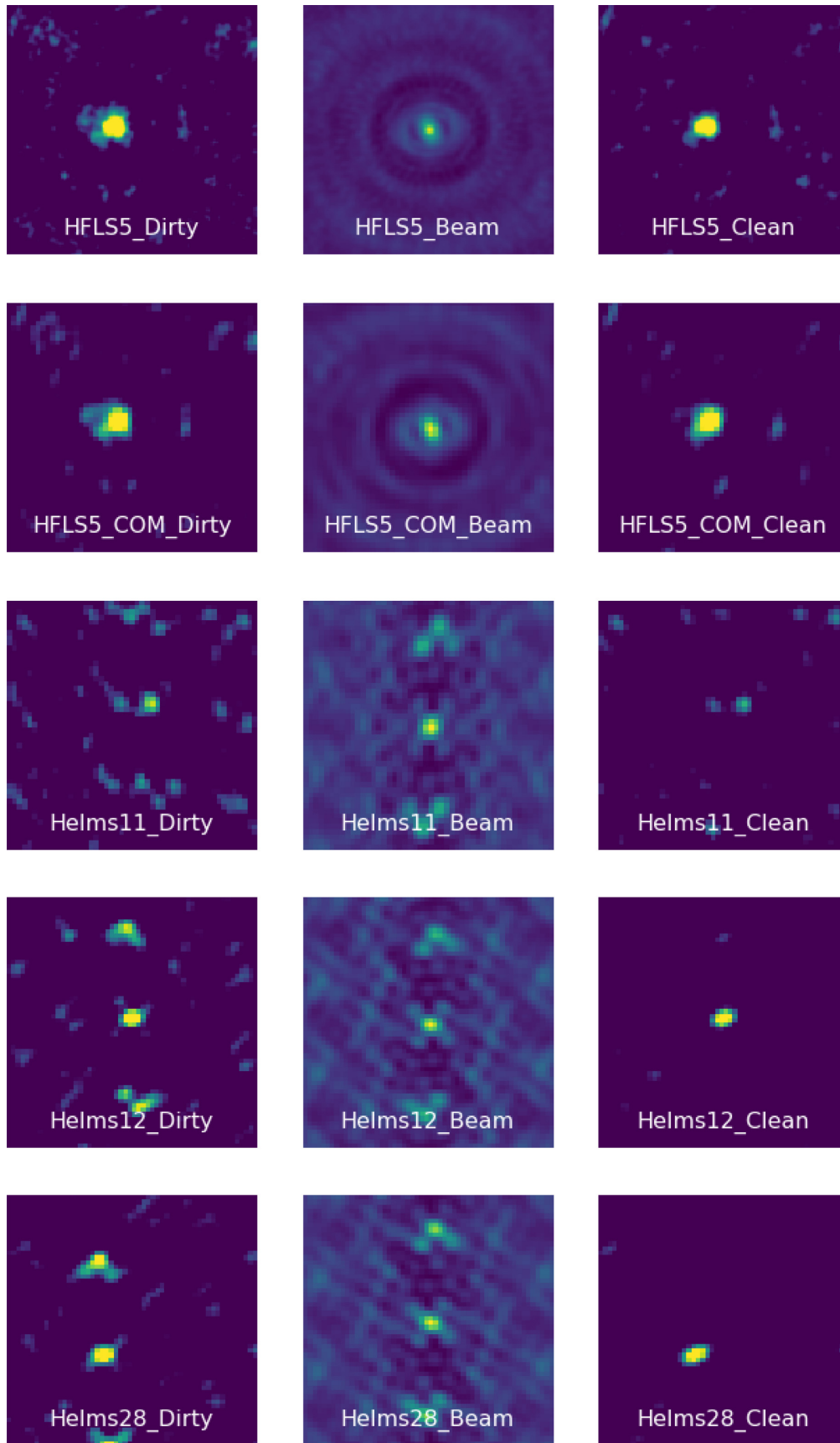
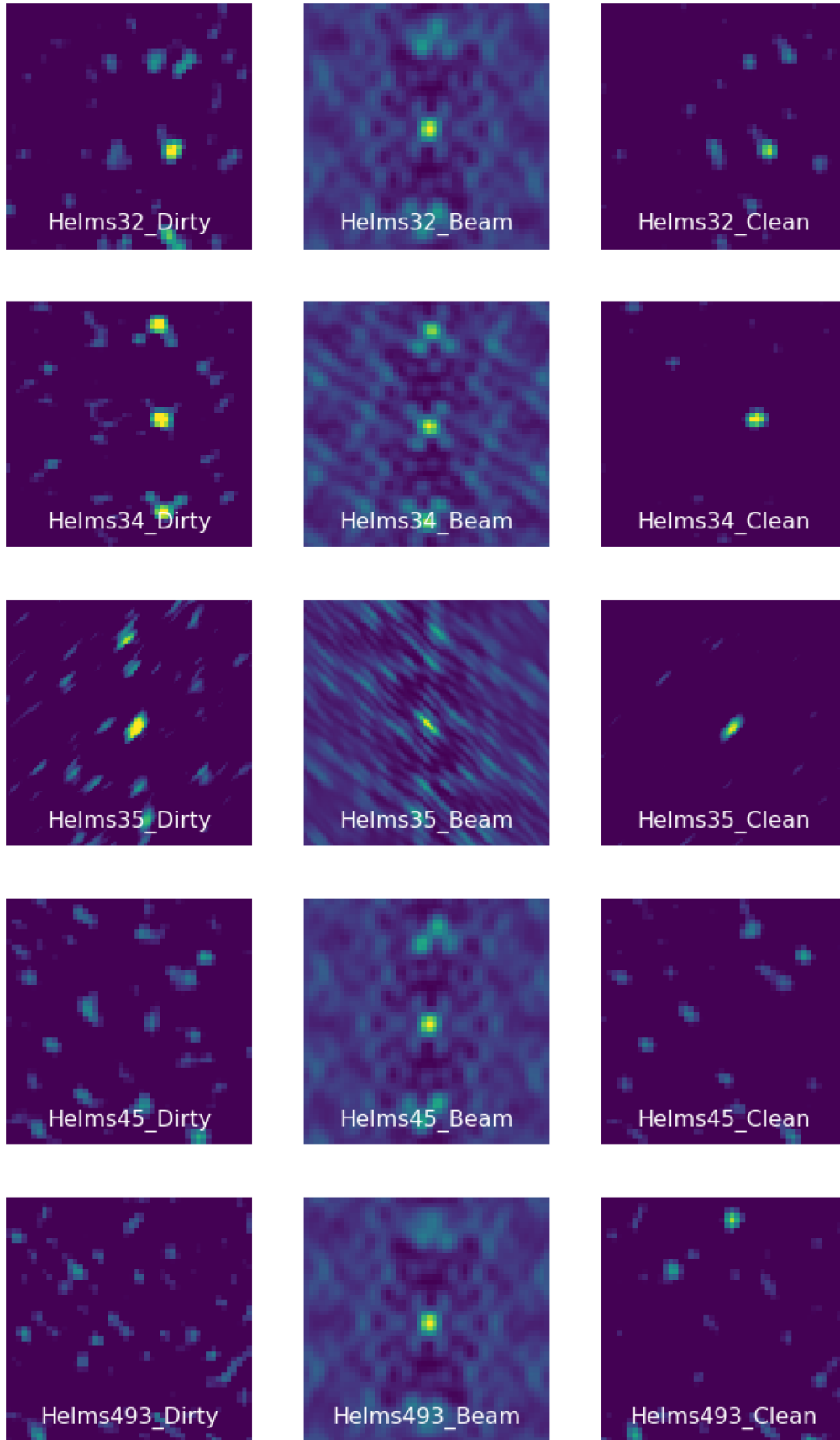


Figure A1. *continued.*

Figure A1. *continued.*

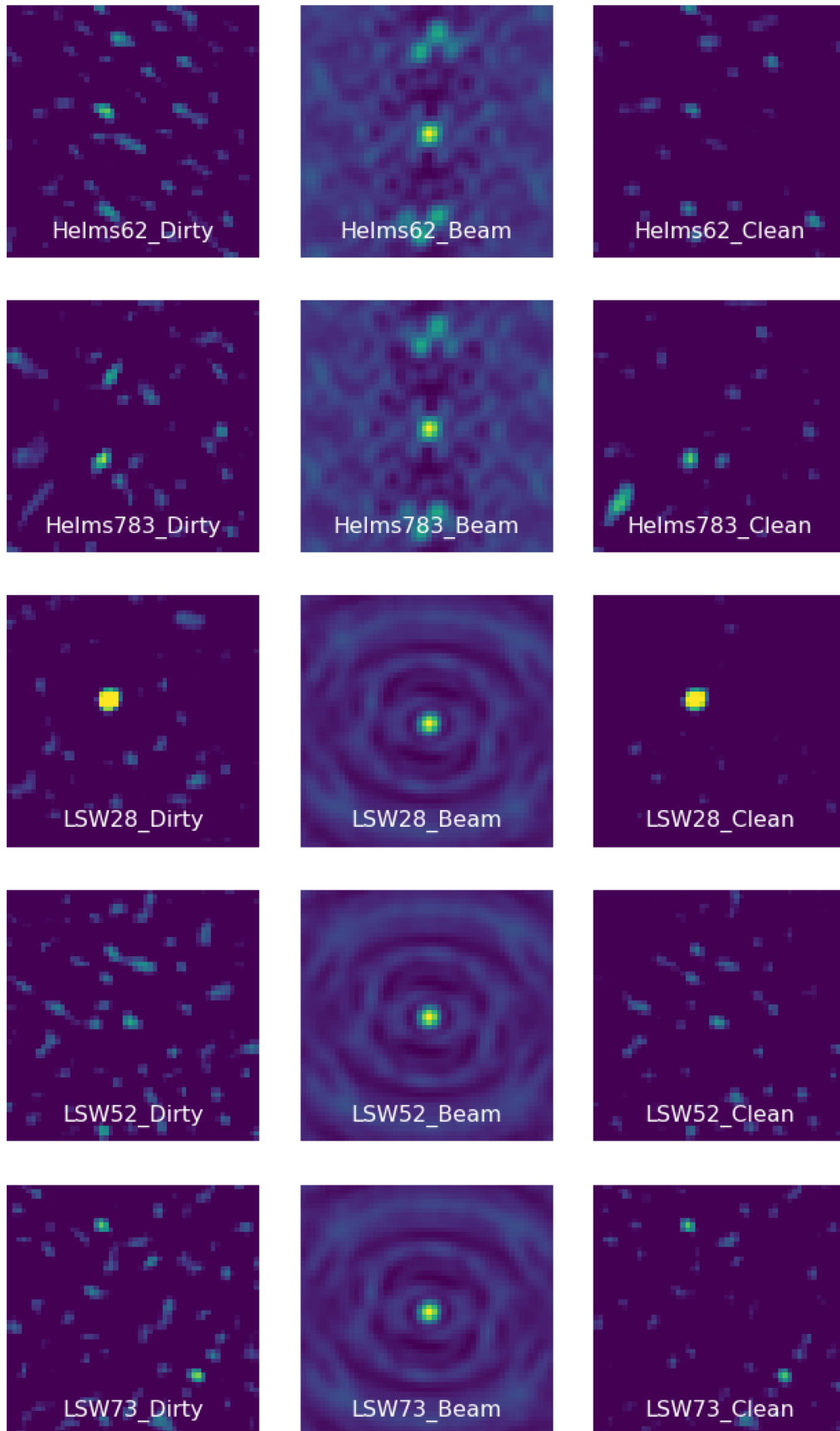
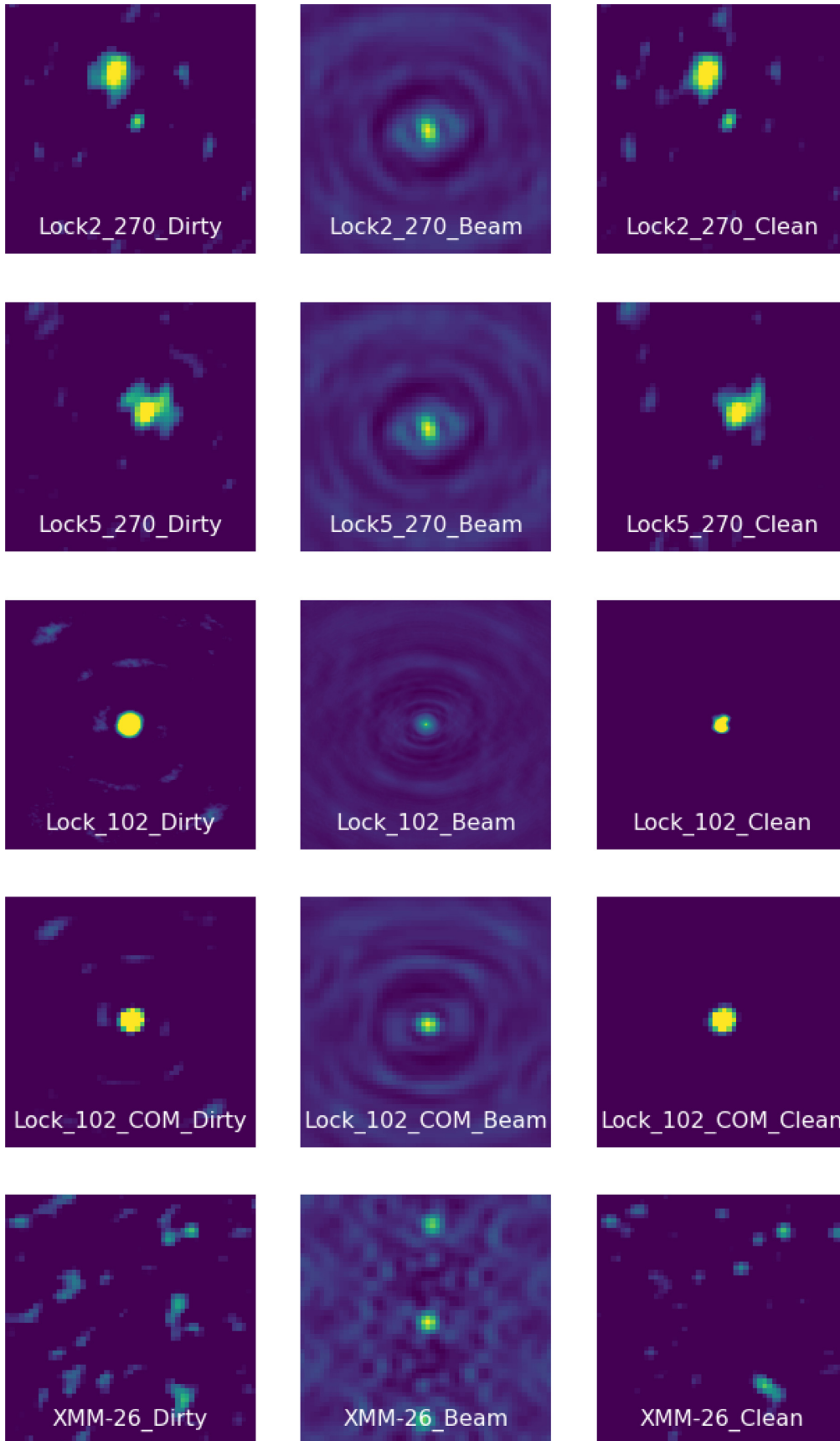


Figure A1. *continued.*

Figure A1. *continued.*

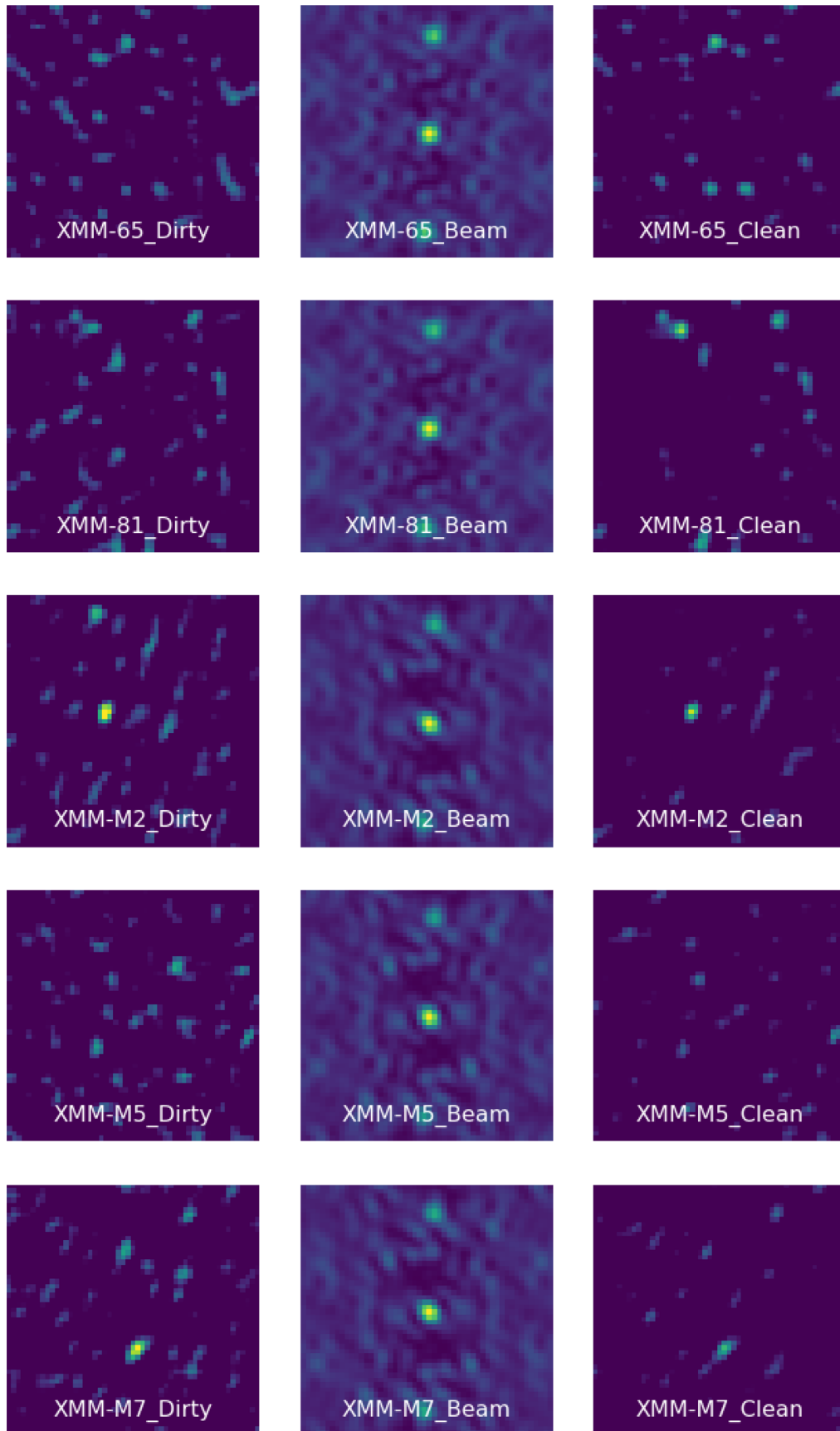


Figure A1. *continued.*

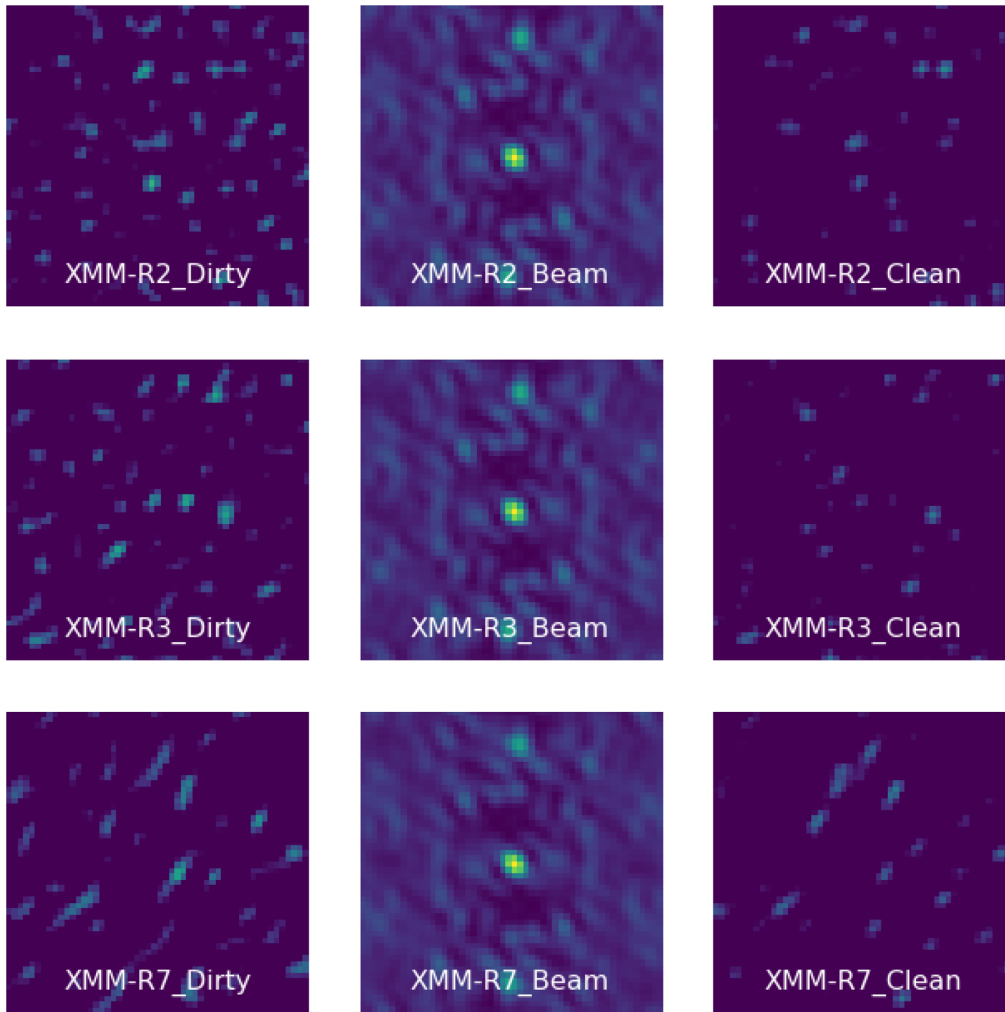


Figure A1. *continued.*

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